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GAUSSIAN JITTER OF A FOCUSED BEAM OF LIGHT. (U)

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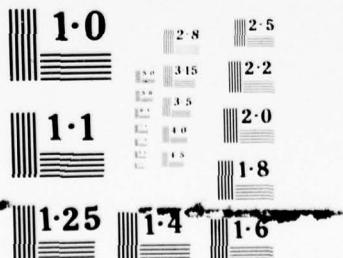
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LEVEL III

## GAUSSIAN JITTER OF A FOCUSED BEAM OF LIGHT

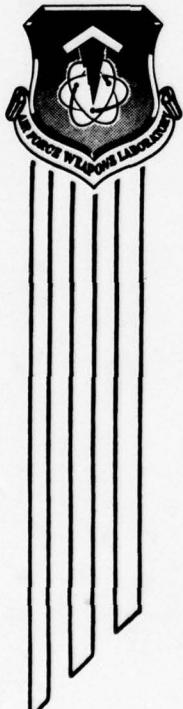
Warren T. White  
John R. Baumgardner  
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LEVEL III

April 1976

Final Report

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AIR FORCE WEAPONS LABORATORY  
Air Force Systems Command  
Kirtland Air Force Base, NM 87117

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The research was prepared by the Air Force Weapons Laboratory, Kirtland Air Force Base, New Mexico, under Job Order 317J1897. Capt. Warren T. White, III (ALO) was the Laboratory Project Officer-in-Charge.

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This technical note has been reviewed and is approved.

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effective peak irradiance as a function of jitter. Annuli with area obscuration ratios of 0.0, 0.1, 0.2, and 0.3 are considered.

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SECTION I  
INTRODUCTION

Negro (1), Esposito (2), and others have proposed expressions to describe effects of jitter upon propagating, fundamental-mode, gaussian laser beams. No one has yet considered the effect of jitter on annular laser beams.

In many unstable resonator lasers, the output of the laser is approximately a beam of uniform phase and gaussian irradiance truncated by an annular aperture. As a first approximation to such a beam, this article considers a monochromatic beam of uniform phase and irradiance apertured by an annular slit.

Some effects of two-dimensional, random gaussian jitter on the Fraunhoffer pattern are examined. By two-dimensional, random gaussian jitter we mean a process which causes the center of the Fraunhoffer pattern to wander in such a way that the time-averaged probability density function for finding the pattern centered at any given location in the Fraunhoffer plane is gaussian. The main effect of jitter is to smear the irradiance distribution. The time-averaged Fraunhoffer pattern loses its characteristic fringes, the time-averaged peak irradiance diminishes, and in general the amount of power transmitted by a circular hole that is concentric with the center of the time-averaged Fraunhoffer irradiance pattern diminishes as jitter increases in amplitude.

Section II of this article presents a simple jitter model. Section III looks at the effect of jitter upon time-averaged power transmitted by a

circular aperture centered in the Fraunhofer plane. Section IV looks at the effect of jitter upon peak irradiance of the time-averaged Fraunhofer pattern. Section V presents graphs of the effective, time-averaged Fraunhofer irradiance distribution,  $\langle I \rangle$ , as a function of radial coordinate and of jitter. It concludes the paper by presenting an approximation to  $\langle I \rangle$ .

## SECTION II

## THE MODEL

The fundamental equation of this paper is the following two-dimensional convolution of the circularly symmetric functions  $I(\rho; \alpha)$  and  $p(\rho; \sigma)$ :

$$\begin{aligned} \langle I \rangle &= I(\rho; \alpha) * p(\rho; \sigma) \\ &= \int_0^\infty \int_0^{2\pi} I(\sqrt{\rho^2 + \rho_0^2 - 2\rho\rho_0 \cos\theta_0}; \alpha) p(\rho_0; \sigma) \rho_0 d\theta_0 d\rho_0 \end{aligned} \quad (1)$$

where

$$I(\rho; \alpha) = \frac{4I_0}{(1 - \alpha^2)^2} \left\{ \frac{J_1(\pi\rho) - \alpha J_1(\pi\alpha\rho)}{\pi\rho} \right\}^2 \quad (2)$$

and

$$p(\rho; \sigma) = \frac{\exp(-\rho^2/\sigma^2)}{\pi\sigma^2} \quad (3)$$

$I(\rho; \alpha)$  is the unjittered Fraunhofer irradiance pattern caused by focusing a plane wave of uniform irradiance through an annular aperture having a ratio of inner diameter to outer diameter of  $\alpha$  (figure 1). We shall refer to  $\alpha^2$  as the obscuration ratio of the aperture.

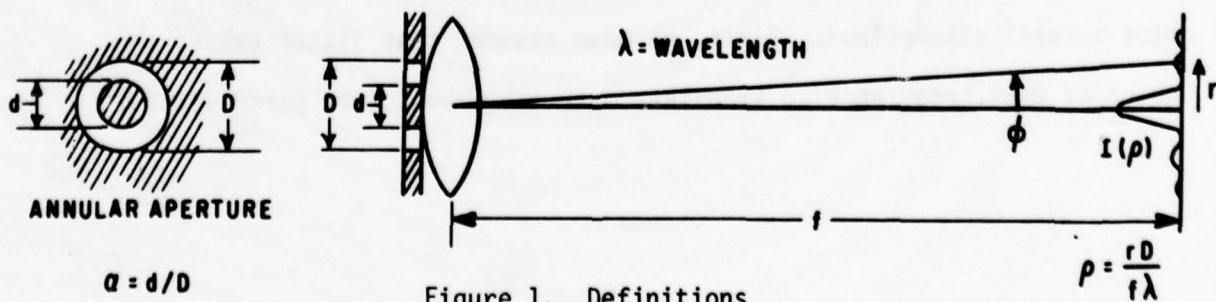


Figure 1. Definitions

$I_0$  is the peak Fraunhoffer irradiance.  $\rho$  is a dimensionless far-field coordinate given by

$$\rho = \frac{rD}{f\lambda} \quad (4)$$

where  $D$  is the outer diameter of the aperture,  $f$  is the distance from the aperture to the Fraunhoffer plane,  $\lambda$  is the wavelength of radiation, and  $r$  is the distance from the origin of the Fraunhoffer plane to any arbitrary point in the Fraunhoffer plane. The function  $p(\rho; \sigma)$  is the time-averaged probability density function that the location of the peak value of the jittering irradiance distribution is inside of differential area  $\rho d\theta d\rho$  which is located at  $(\rho, \theta)$ . In this article the parameter  $\sigma$  is defined as the rms jitter. It is dimensionless. Mathematically it is the square root of  $\sigma^2$ , where

$$\begin{aligned} \sigma^2 &= \langle \rho^2 \rangle \\ &= 2\pi \int_0^\infty \rho^2 p(\rho, \sigma) \rho d\rho \end{aligned} \quad (5)$$

In the expression for  $I(\rho; \alpha)$ , we have assumed that we have a diffraction-limited focal spot and that  $f \gg D \gg \lambda$  and  $f \gg r$ .

By approximating the effective irradiance,  $\langle I \rangle$ , as a convolution, we have made several assumptions. First, we have assumed that jitter causes the point of peak irradiance to translate over the Fraunhoffer plane without

distorting the irradiance distribution relative to that point. Secondly, we have assumed that the focal surface is truly a plane. Thirdly, we have assumed that infinitely large jitter displacements might occur. For a brief discussion of these last three assumptions, see Appendix A.

SECTION III  
ENCIRCLED POWER AS A FUNCTION  
OF CIRCLE SIZE

Encircled power is the total radiant power contained within an illuminated circle. Let  $\langle P(a) \rangle$  denote time-averaged encircled power over a circle of radius  $a$  that is centered at the origin of the Fraunhoffer plane.

$$\langle P(a) \rangle = 2\pi \left(\frac{f\lambda}{D}\right)^2 \int_0^{aD/f\lambda} \langle I \rangle \rho d\rho \quad (6)$$

The quantity  $\langle P(a) \rangle / P_T$ , where  $P_T$  is the total output power over the entire Fraunhoffer plane, is a measure of how well focused a beam of light is. The effect of jitter upon  $\langle P(a) \rangle / P_T$  is generally to decrease its value. An exception occurs when jitter is slight and the rim of the circle is located near a fringe. In such cases, a small amount of jitter tends to increase the encircled power.

The graphs in figures 2(a)-2(d) depict encircled power as a function of circle size. The ordinates are normalized so that an infinitely large circle ( $aD/f\lambda = \infty$ ) causes the function to equal unity. Each graph consists of a family of curves for a given obscuration ratio,  $\alpha^2$ . Each curve in a family represents  $\langle P(a) \rangle / P_T$  as a function of the dimensionless variable  $aD/f\lambda$  for a fixed amount of rms jitter.

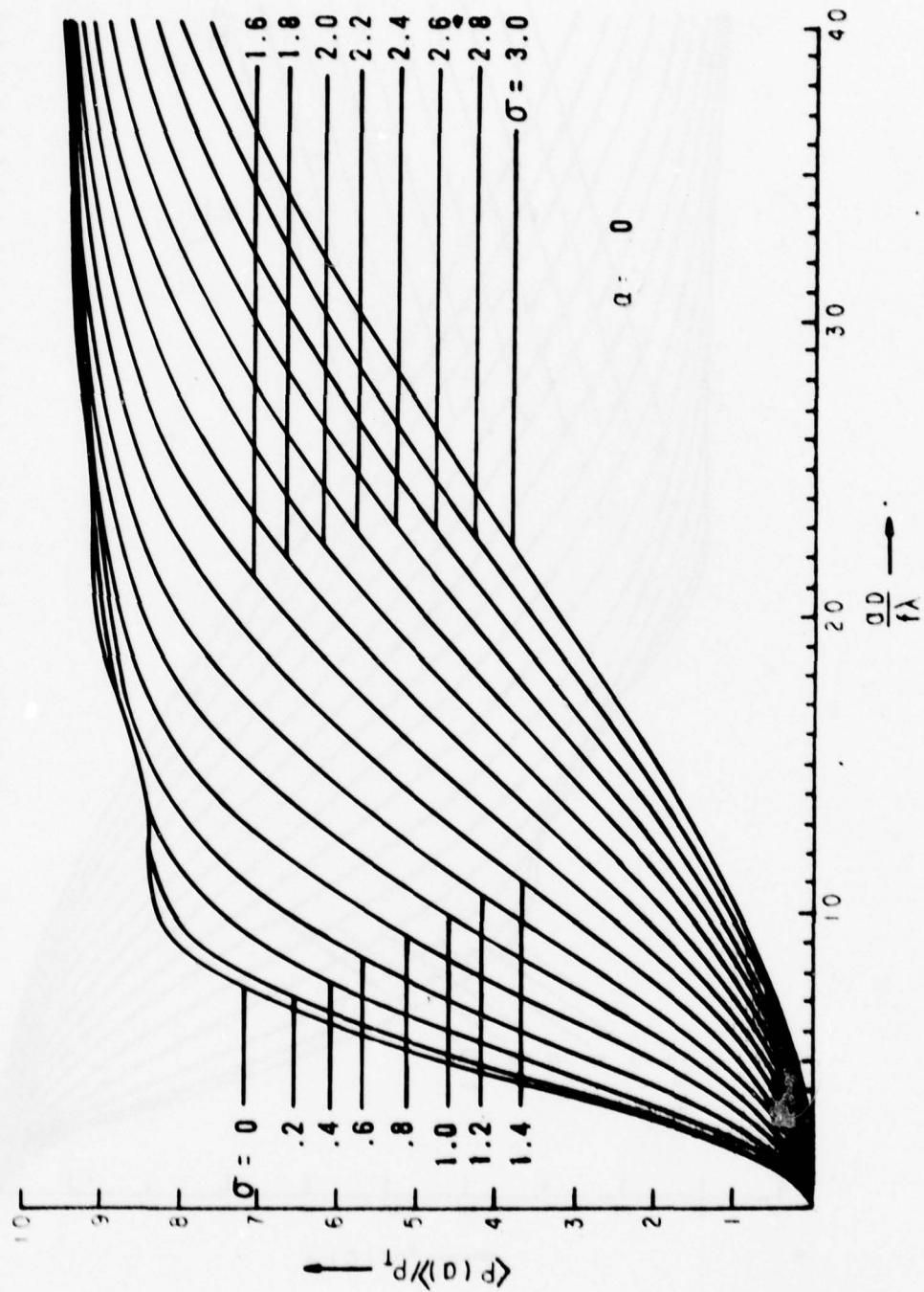


Figure 2(a). Encircled Far-Field Power as a Function of Circle Radius,  $a$ , and RMS Gaussian Jitter,  $\sigma$ . Near-field irradiance is a uniformly illuminated circle.

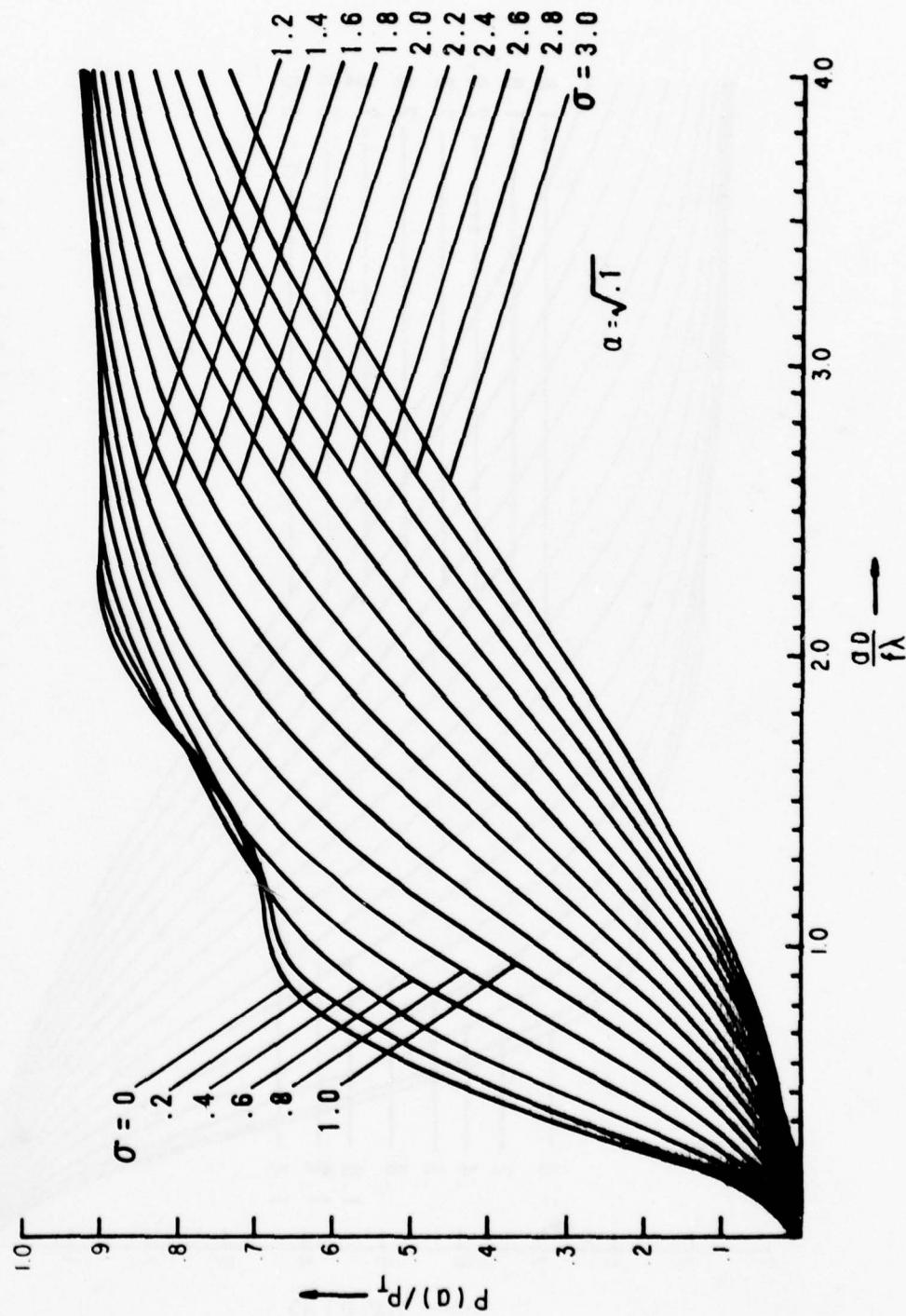


Figure 2(b). Encircled Far-Field Power as a Function of Circle Radius,  $a$ , and RMS Gaussian Jitter,  $\sigma$ . Near-field irradiance is a uniformly illuminated annulus with 10% obscuration.

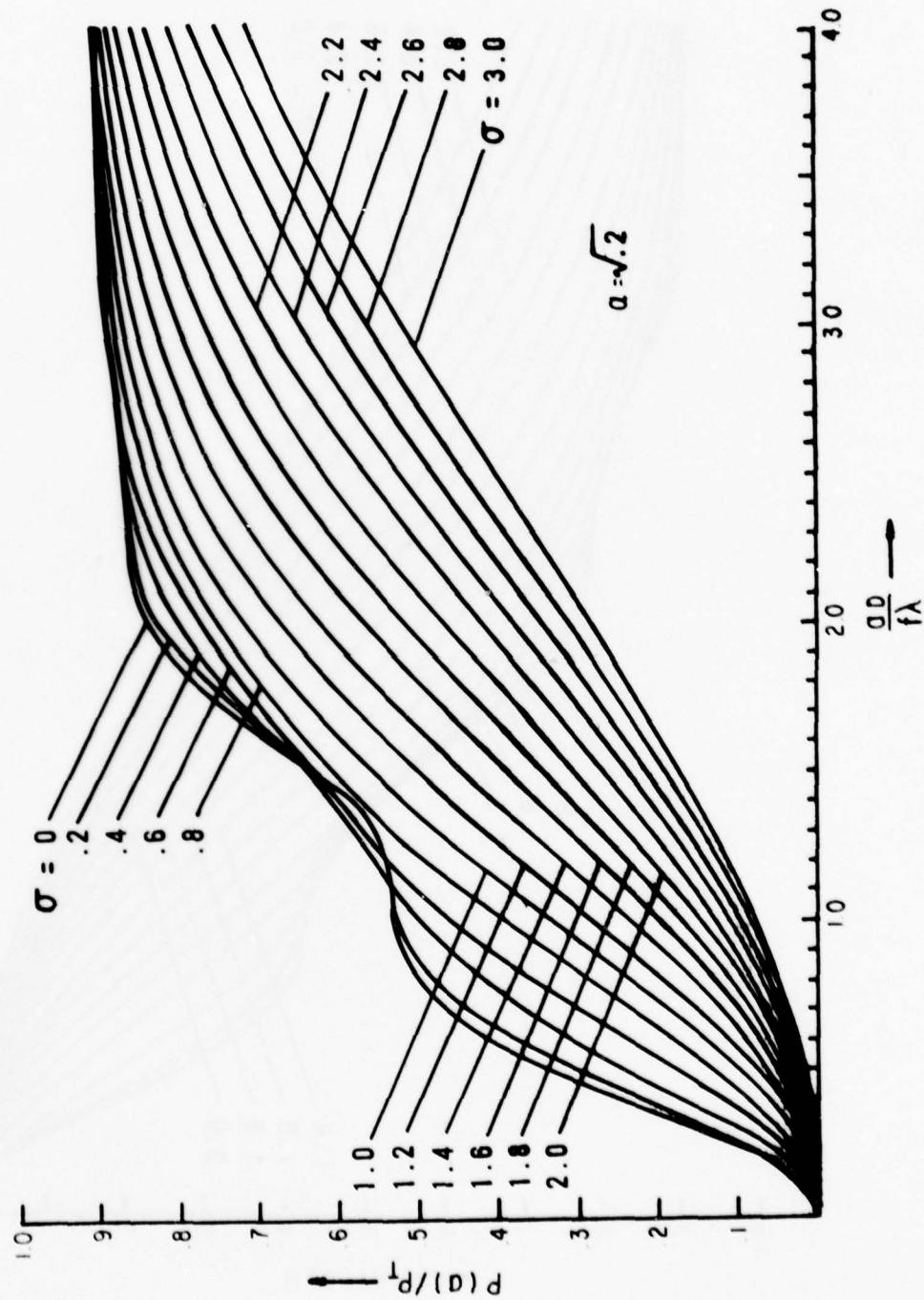


Figure 2(c). Encircled Far-Field Power as a Function of Circle Radius,  $a$ , and RMS Gaussian Jitter,  $\sigma$ . Near-field irradiance is a uniformly illuminated annulus with 20% obscuration.

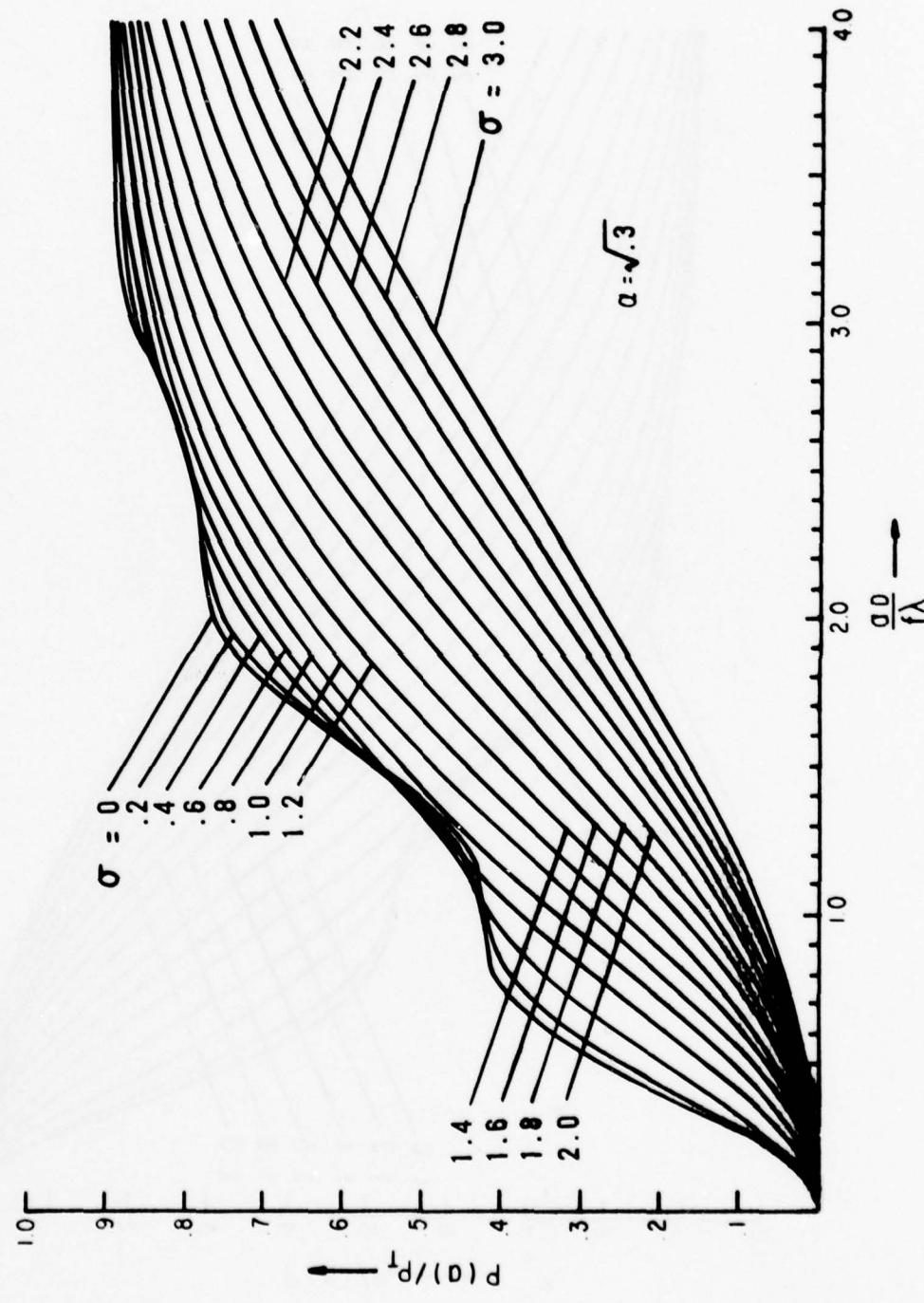


Figure 2(d). Encircled Far-Field Power as a Function of Circle Radius,  $a$ , and RMS Gaussian Jitter,  $\sigma$ . Near-field irradiance is a uniformly illuminated annulus with 30% obscuration.

Some general comments apply to all of these graphs. First, on each graph the curve corresponding to no jitter tends to have points of zero slope, corresponding to fringes. As jitter increases these points of zero slope disappear, corresponding to fringe wash-out. Secondly, given a value of  $aD/f\lambda$ , an increase in rms jitter,  $\sigma$ , tends to decrease encircled power except as previously noted. However, given a large enough circle, the effect of jitter, for small enough jitter, is unnoticeable. For example, figure 2(a) shows that for any circle whose radius,  $a$ , is larger than  $a' = 2f\lambda/D$ ,  $\langle P(a) \rangle / P_T$  is independent of the rms jitter,  $\sigma$ , to within a tolerance of about two percent as long as  $\sigma \lesssim 0.5$ . For a larger value of  $a'$ , say  $a' = 3f\lambda/D$ ,  $\langle P(a) \rangle / P_T$  is approximately independent of  $\sigma$  for all  $a$  such that  $a \geq a'$  and for  $\sigma \lesssim 1.0$ . In less mathematical terms, since most of the power in the focused laser beam is near the point of peak irradiance, if the point of peak irradiance, although it may be moving from point to point, stays well within the rim of a circle whose radius equals at least  $a'$ , then the amount of power transmitted by that circle is approximately the same as it would be if the laser beam were not jittering.

## SECTION IV

TIME-AVERAGED AXIAL  
FRAUNHOFFER IRRADIANCE,  $\langle I(0) \rangle$ 1.  $\langle I(0) \rangle$  AS A FUNCTION OF  $\sigma$ 

Figure 3 shows how time-averaged axial Fraunhofer irradiance deteriorates if jitter exists. These curves are hand-smoothed lines fit through a discrete sequence of digital computer evaluations of the integral

$$\frac{\langle I(0) \rangle}{I_0} = \frac{2\pi}{I_0} \int_0^\infty I(\rho) e^{-\rho^2/\sigma^2} \rho \, d\rho \quad (7)$$

The curves are normalized to unity for no jitter. The normalizing factor  $I_0$  is the peak instantaneous Fraunhofer irradiance, as in equation (2).

As an illustration of the jitter-induced degradation of axial Fraunhofer irradiance, consider an unstable optical resonator with uniformly intense, collimated annular output which is subsequently focused at  $f = 2.5$  km and has

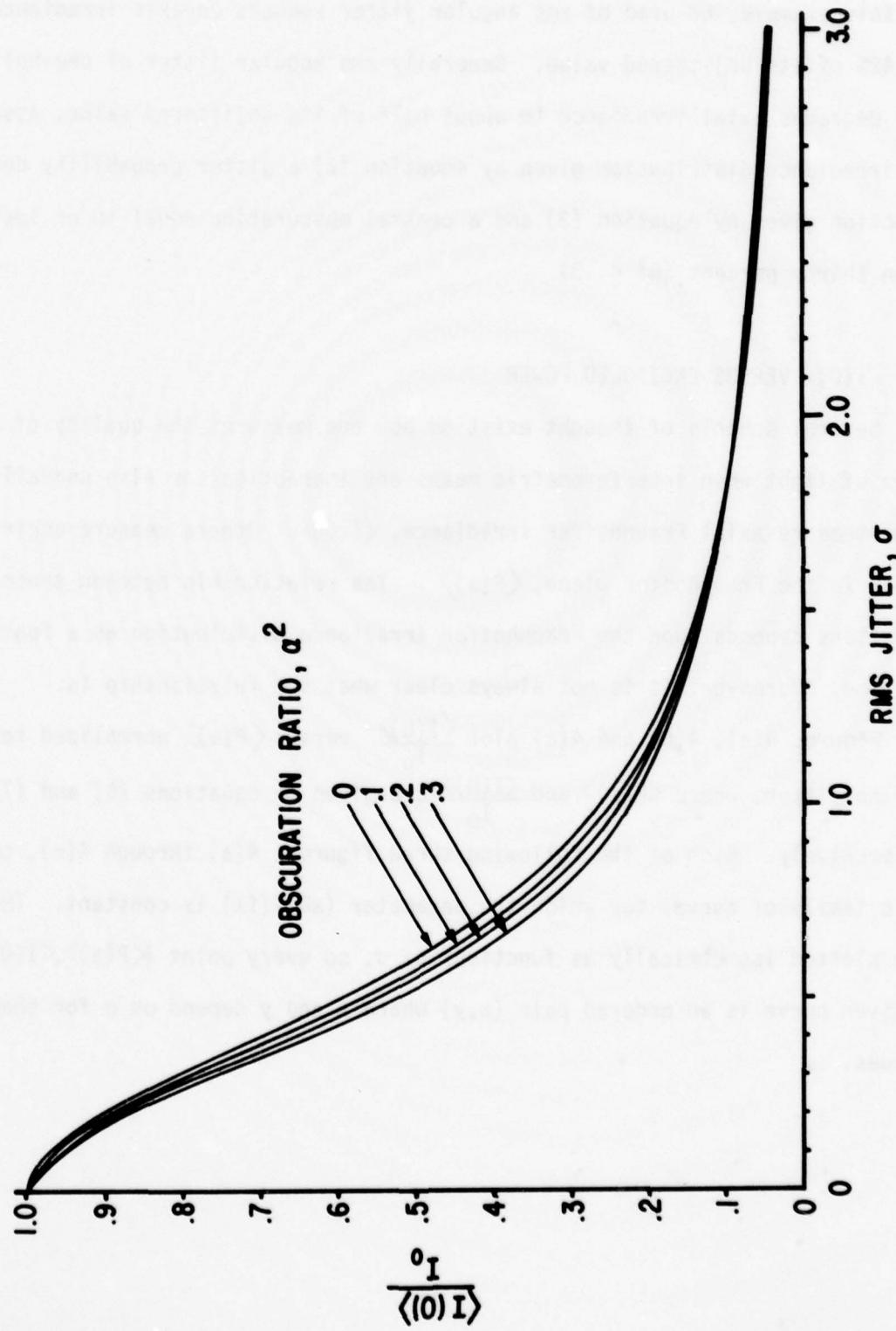
$$\lambda = 10 \mu\text{m}$$

$$D = 10 \text{ cm}$$

$$\alpha^2 = .3 \text{ (70\% geometric output coupling).}$$

Assume the beam vibrates in a random gaussian fashion with an rms angular jitter of  $60 \mu\text{rad}$ . This is equivalent to assuming  $\sigma = 0.6$ . From figure 3

$$\frac{\langle I(0) \rangle}{I_0} \Big|_{\substack{\sigma = .6 \\ \alpha^2 = .3}} = .42$$



**Figure 3.** Degradation of Far-Field On-Axis Irradiance Due to an RMS Line-of-Sight Jitter,  $\sigma$ . RMS jitter is dimensionless but may be thought of as being in units of  $\lambda/D$ .

In this example, 60  $\mu$ rad of rms angular jitter reduces on-axis irradiance to 42% of its unjittered value. Generally rms angular jitter of one-half  $\lambda/D$  degrades axial irradiance to about half of its unjittered value, assuming an irradiance distribution given by equation (2) a jitter probability density function given by equation (3) and a central obscuration equal to or less than thirty percent ( $\alpha^2 \leq .3$ ).

## 2. $\langle I(0) \rangle$ VERSUS ENCIRCLED POWER

Several schools of thought exist on how one measures the quality of a beam of light when interferometric means are impractical or else unavailable. Some measure axial Fraunhofer irradiance,  $\langle I(0) \rangle$ . Others measure encircled power in the Fraunhofer plane,  $\langle P(a) \rangle$ . The relationship between these two functions depends upon the Fraunhofer irradiance distribution as a function of time. Moreover, it is not always clear what the relationship is.

Figures 4(a), 4(b) and 4(c) plot  $\frac{\langle I(0) \rangle}{I_0}$  versus  $\langle P(a) \rangle$  normalized to unity for no jitter, where  $\langle P(a) \rangle$  and  $\frac{\langle I(0) \rangle}{I_0}$  are given by equations (6) and (7) respectively. Each of the following three figures, 4(a) through 4(c), consist of a family of curves for which the parameter  $(aD)/(f\lambda)$  is constant. The data are plotted isometrically as functions of  $\sigma$ , so every point  $(\langle P(a) \rangle, \langle I(0) \rangle)$  on a given curve is an ordered pair  $(x, y)$  where  $x$  and  $y$  depend on  $\sigma$  for their values.

The fact that all points corresponding to  $\sigma \neq 0$  in following graphs fall below the line  $y = x$  (where  $x$  equals  $\langle P(a) \rangle$  normalized to unity for  $\sigma = 0$  and where  $y$  equals  $\langle I(0) \rangle$ ) shows that jitter degrades  $\langle I(0) \rangle$  more than it does  $\langle P(a) \rangle$ .

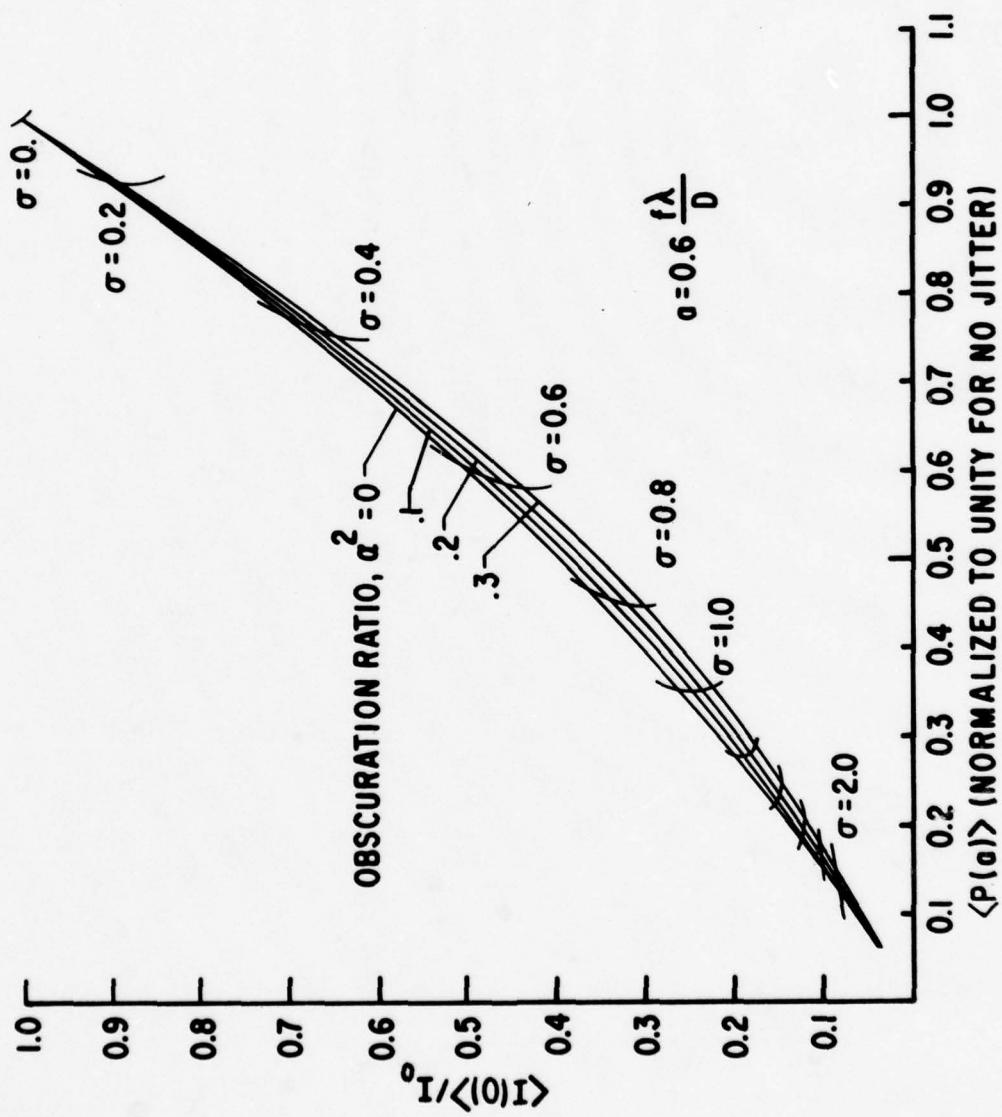


Figure 4(a). Degraded Far-Field On-Axis Irradiance Plotted Against Degraded Far-Field Encircled Power. The radius of the far-field circle is fixed at  $aD/f\lambda = 0.6$ .

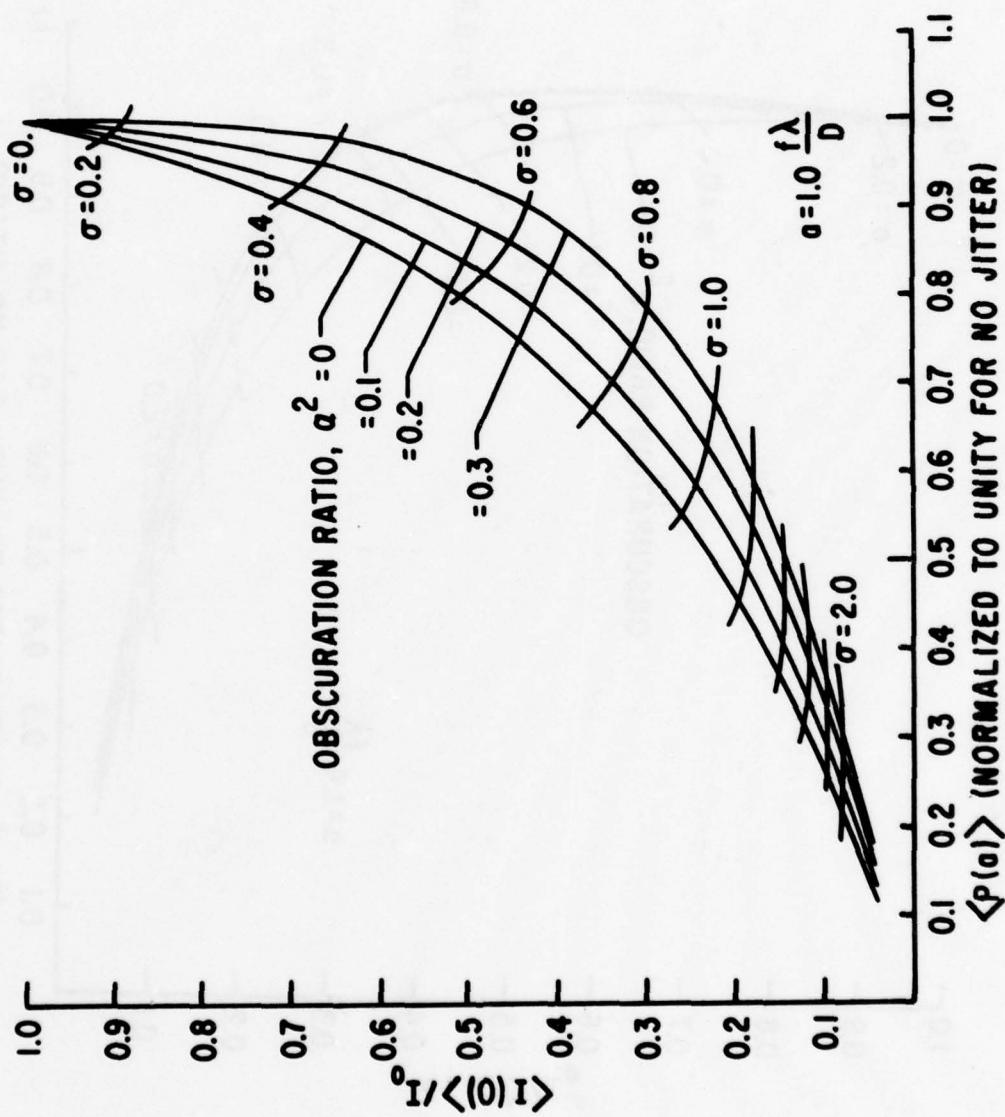
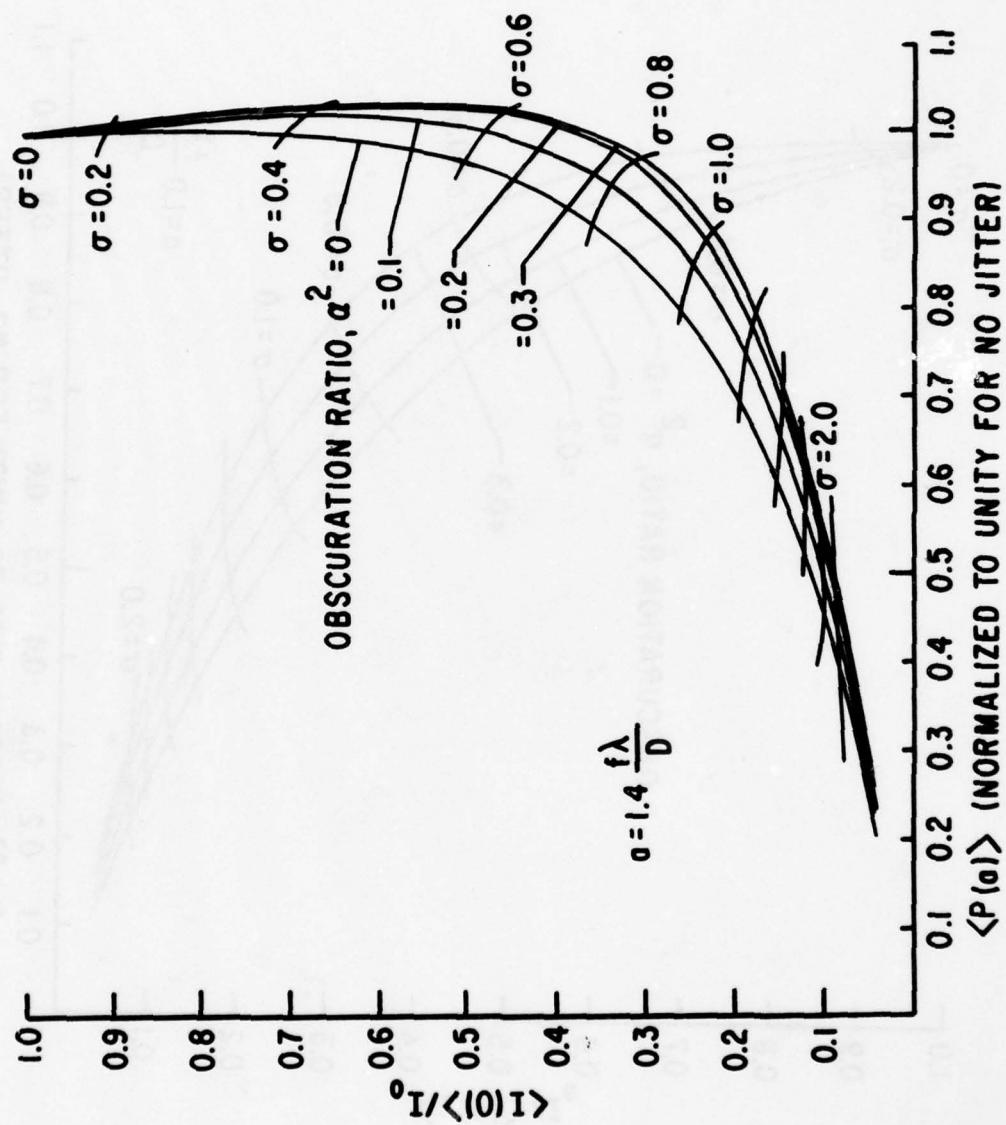


Figure 4(b). Degraded Far-Field On-Axis Irradiance Plotted Against Degraded Far-Field Encircled Power. The radius of the far-field circle is fixed at  $aD/f\lambda = 1.0$ .



**Figure 4(c).** Degraded Far-Field On-Axis Irradiance Plotted Against Degraded Far-Field Encircled Power. The radius of the far-field circle is fixed at  $aD/f\lambda = 1.4$ .

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SECTION V  
TIME-AVERAGED FRAUNHOFFER PATTERN

1.  $\langle I \rangle$  vs  $\rho$

As rms jitter increases, the effective far-field irradiance distribution tends to smooth out. The first part of this section presents computer calculated graphs of the effective irradiance,  $\langle I \rangle$ , versus the normalized far-field coordinate,  $\rho$ .

The computer calculated values of  $\langle I \rangle$  are plotted as a dashed line. Superimposed on each graph is a least-squares fit according to an approximation given in the last part of this section.

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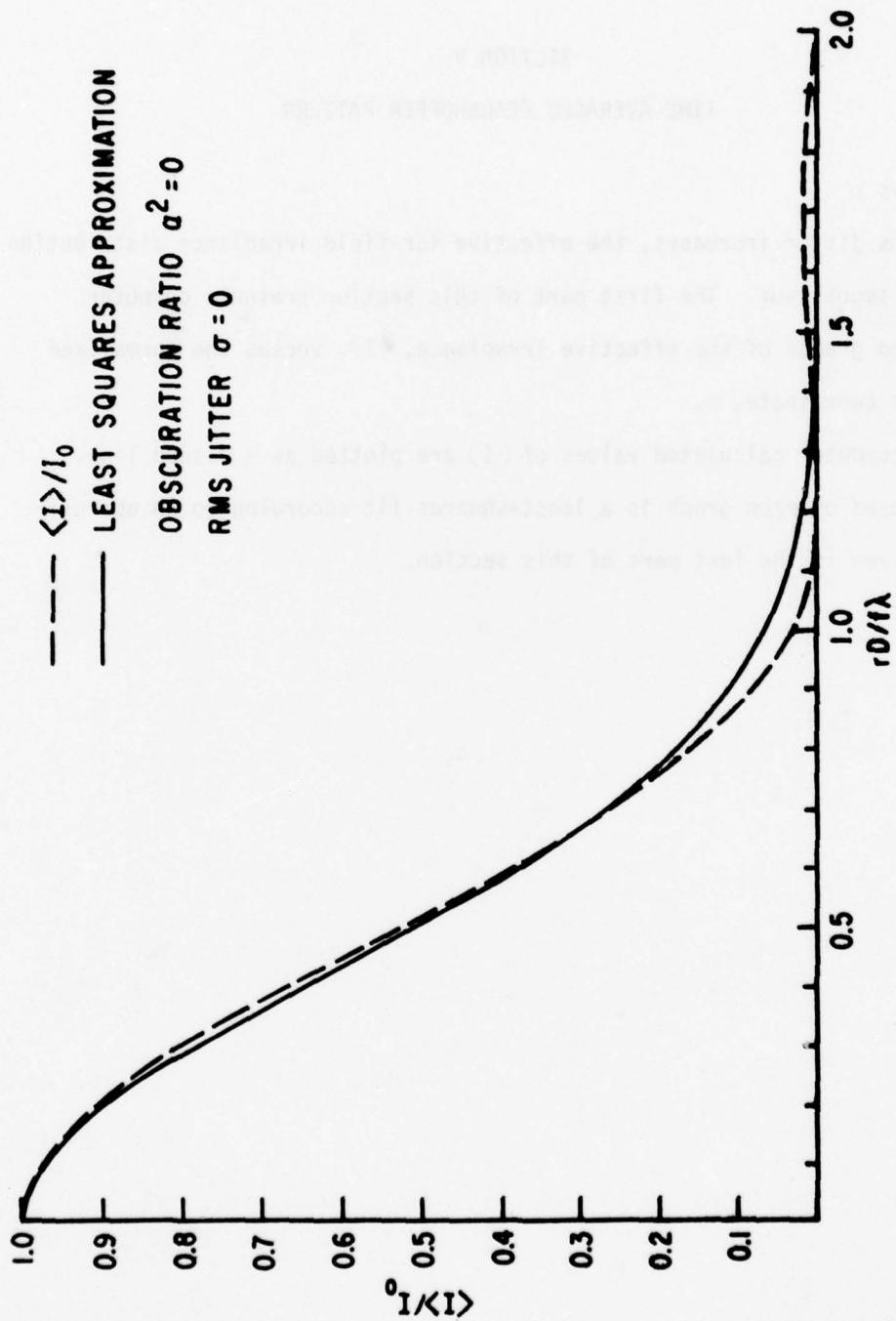


Figure 5(a)-1. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

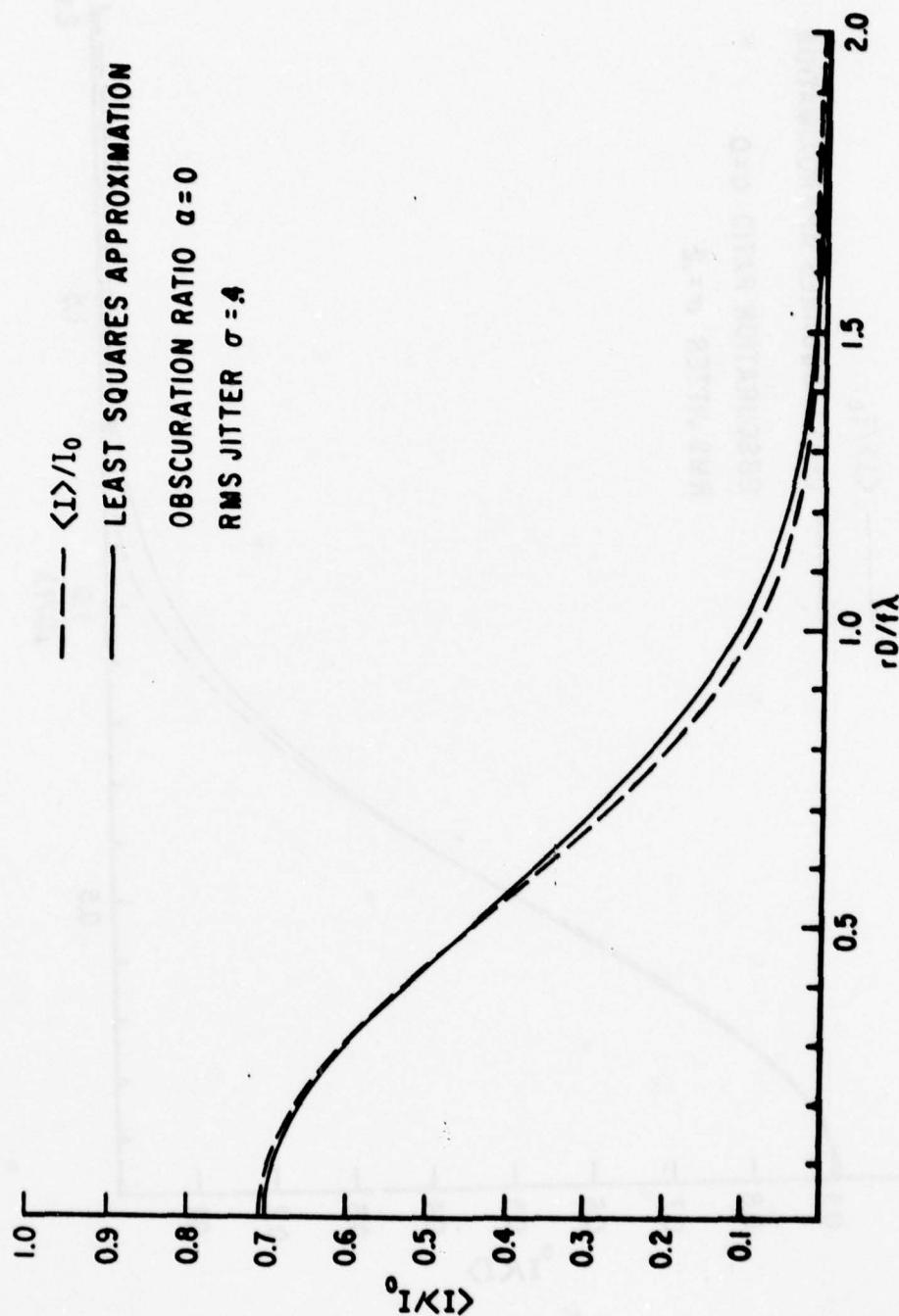


Figure 5(a)-2. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

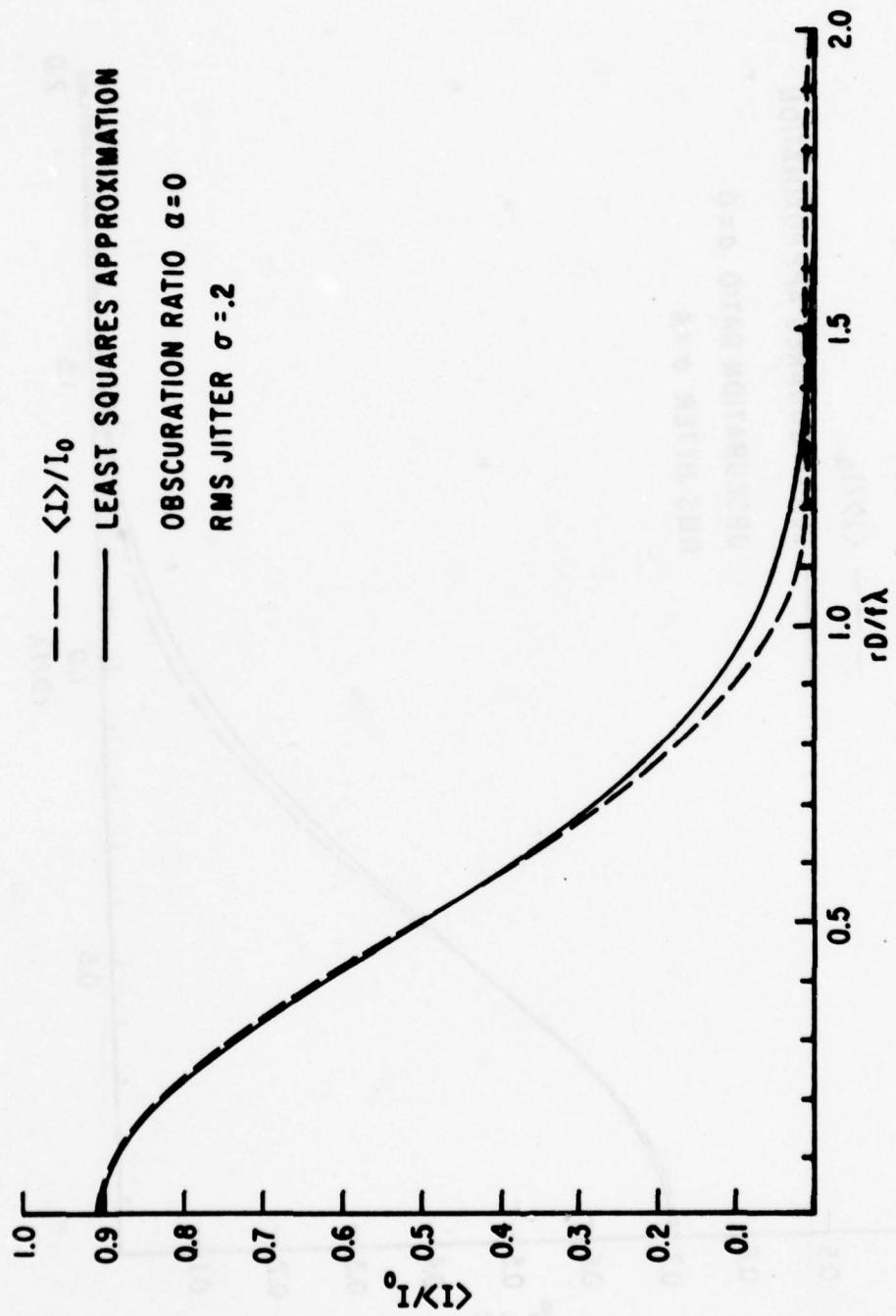


Figure 5(a)-3. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

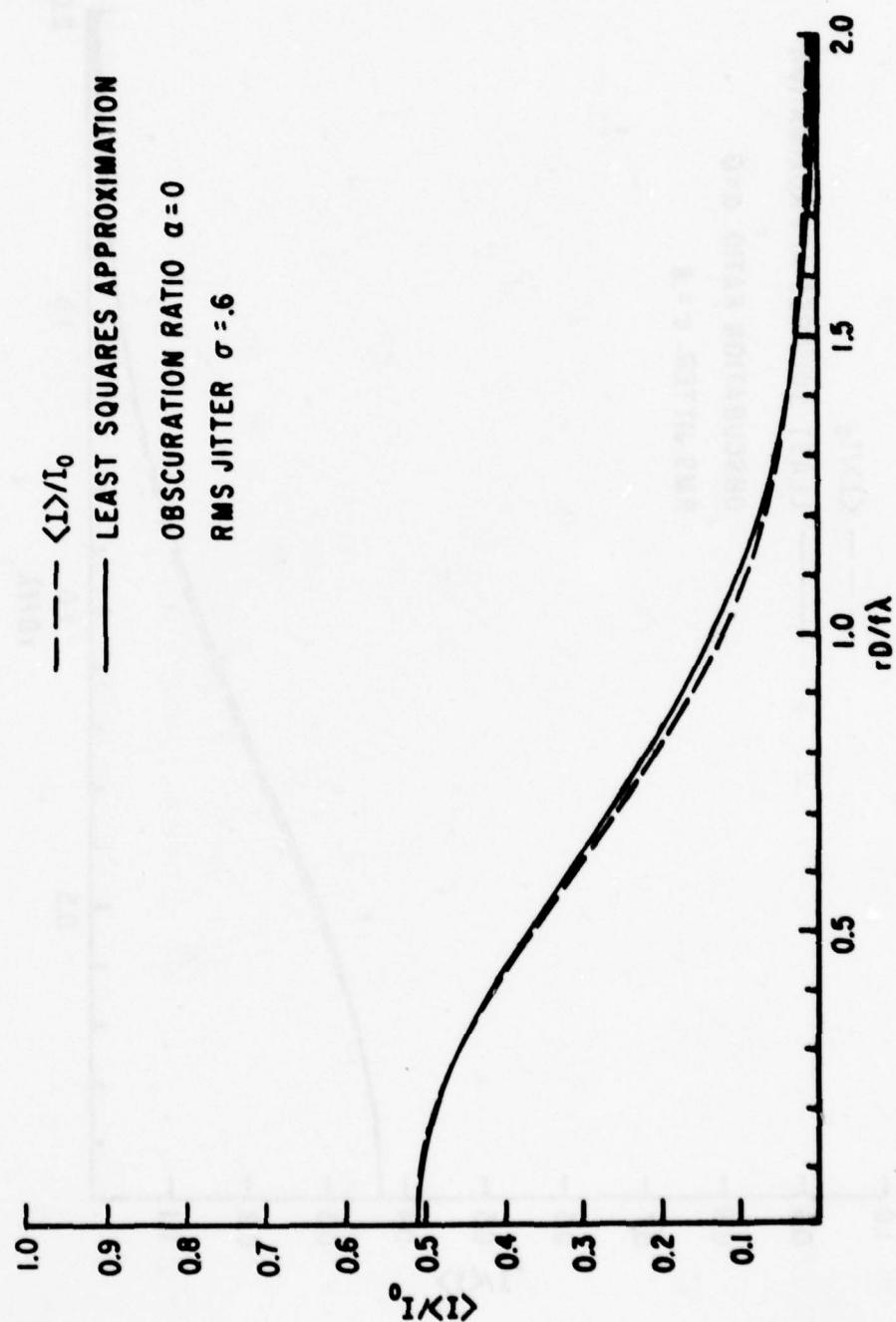


Figure 5(a)-4. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

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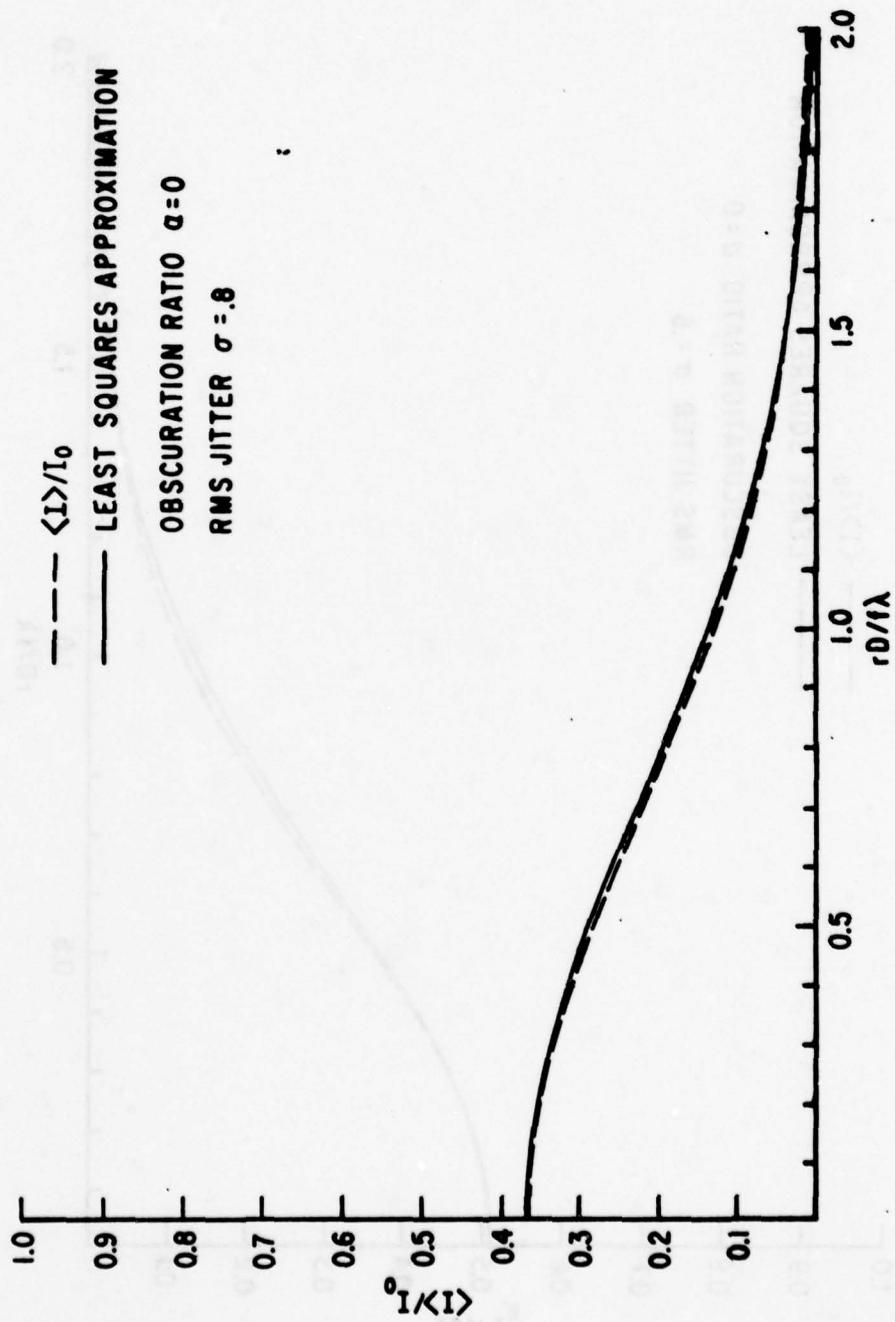


Figure 5(a)-5. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

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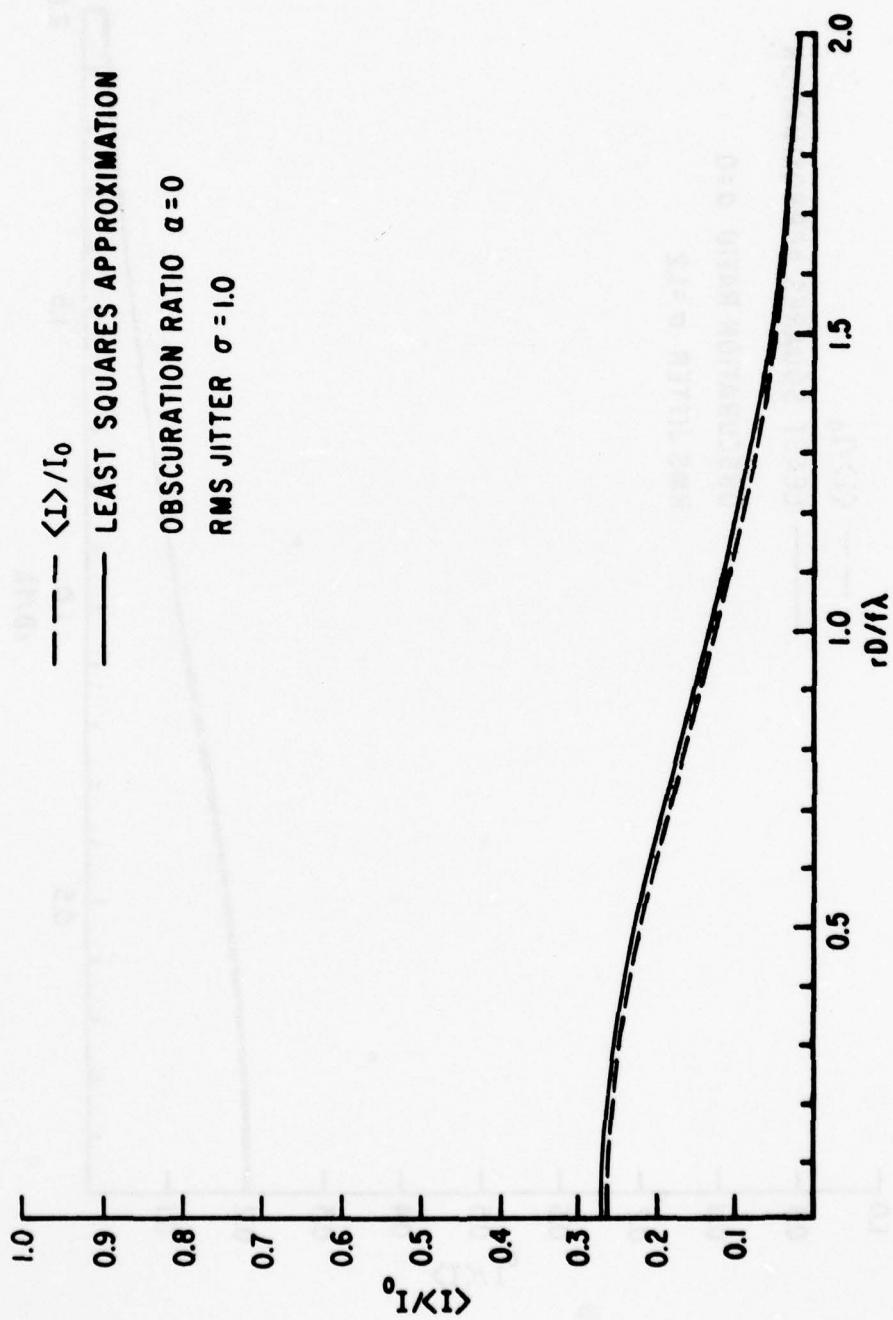


Figure 5(a)-6. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

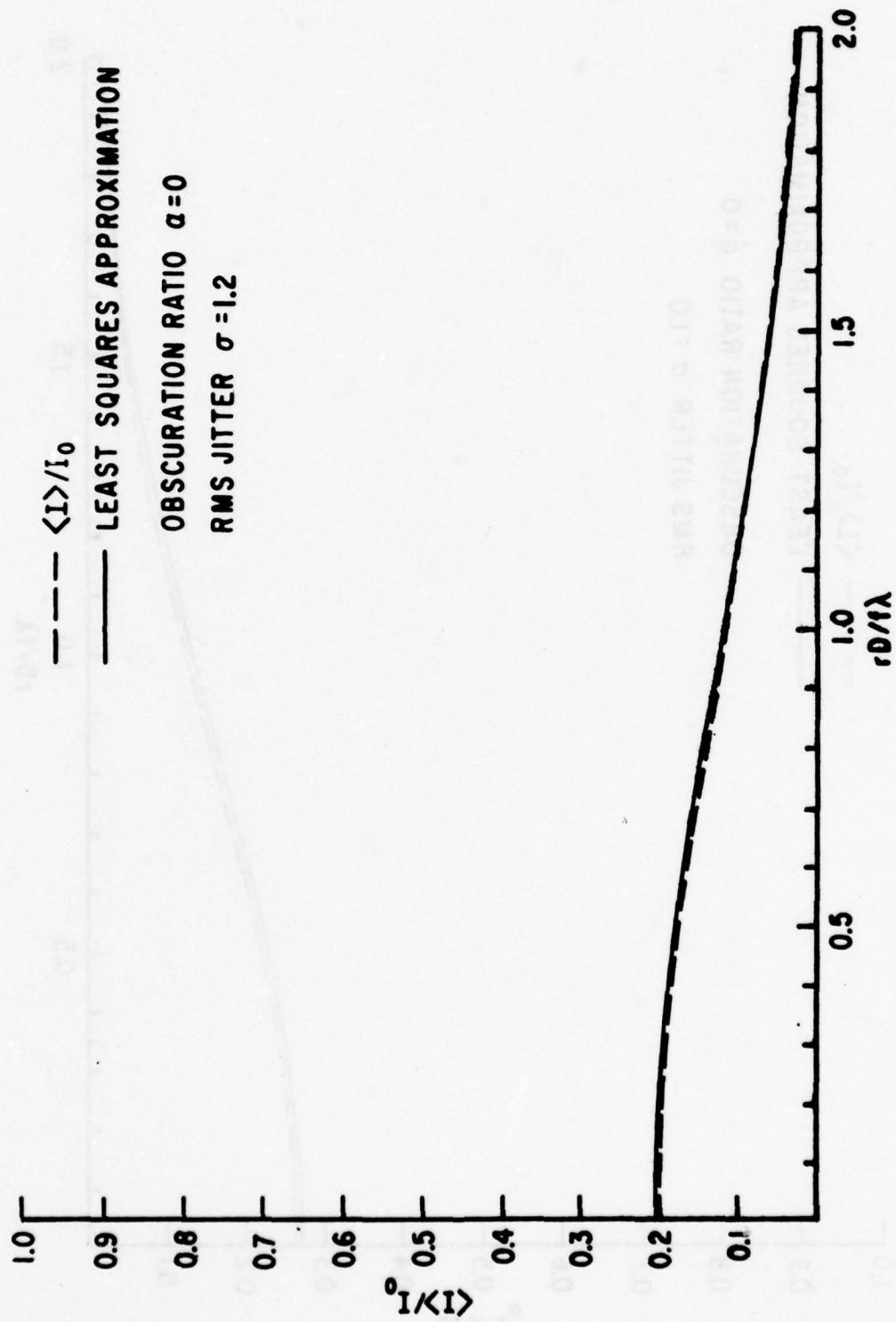


Figure 5(a)-7. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

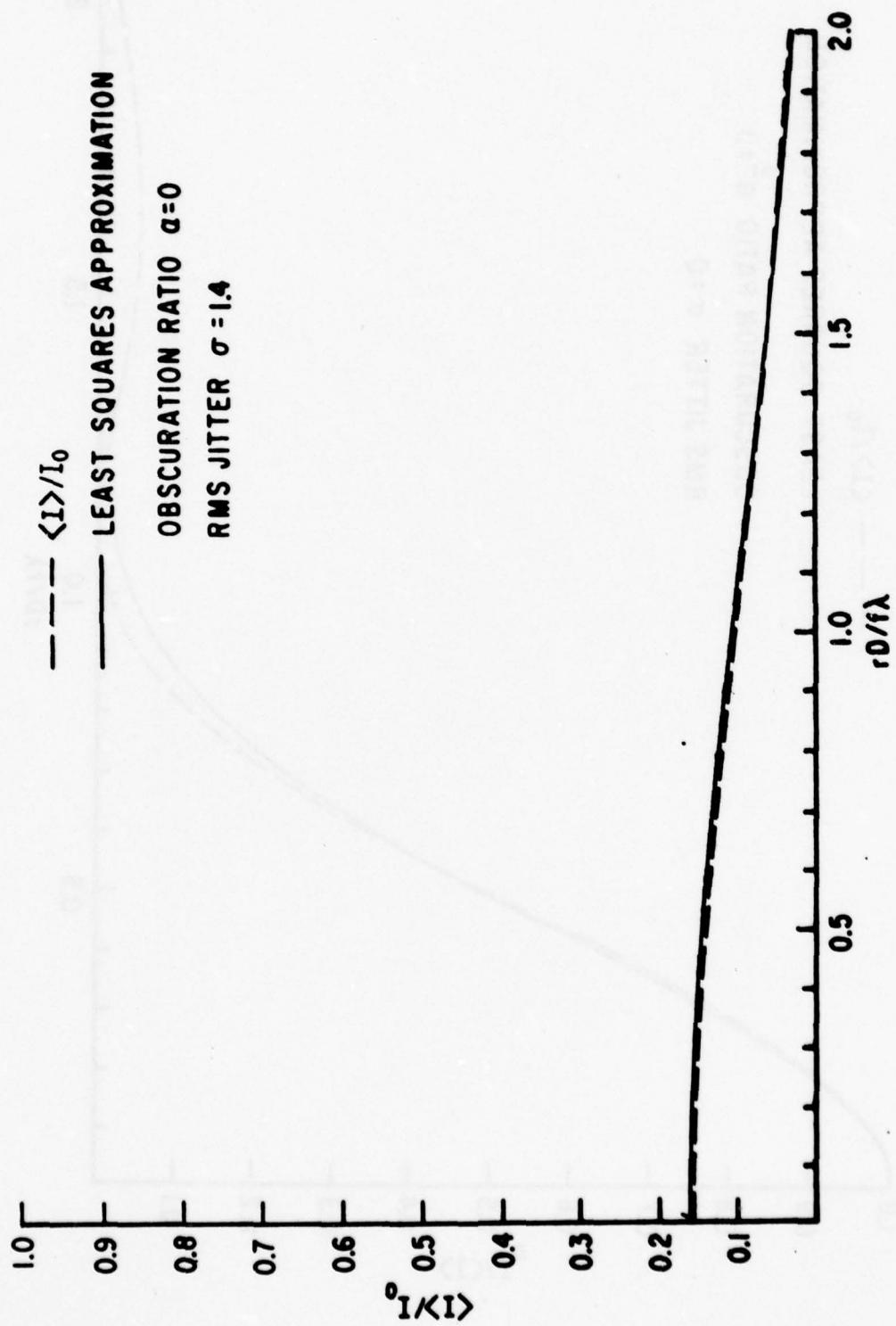


Figure 5(a)-8. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

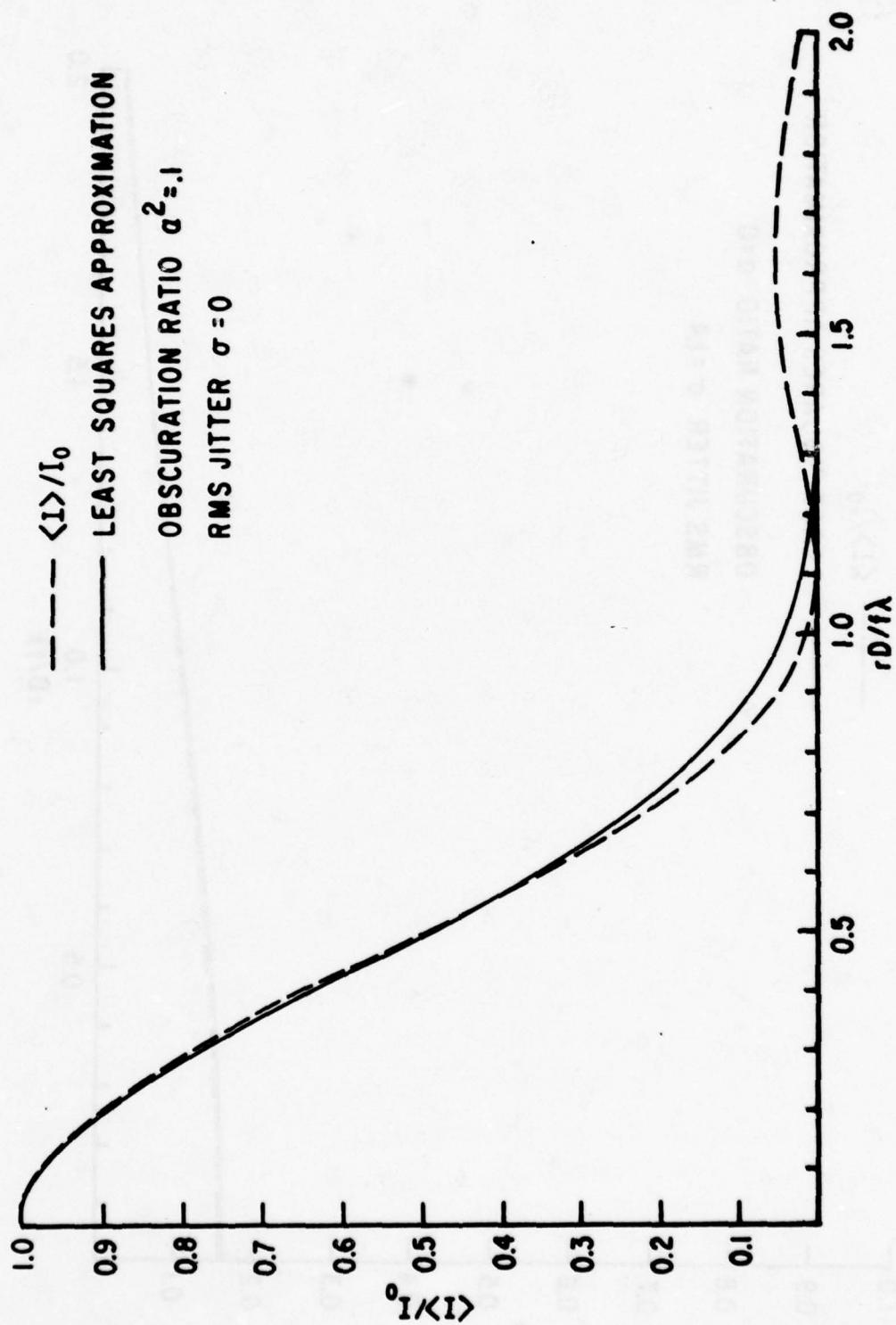
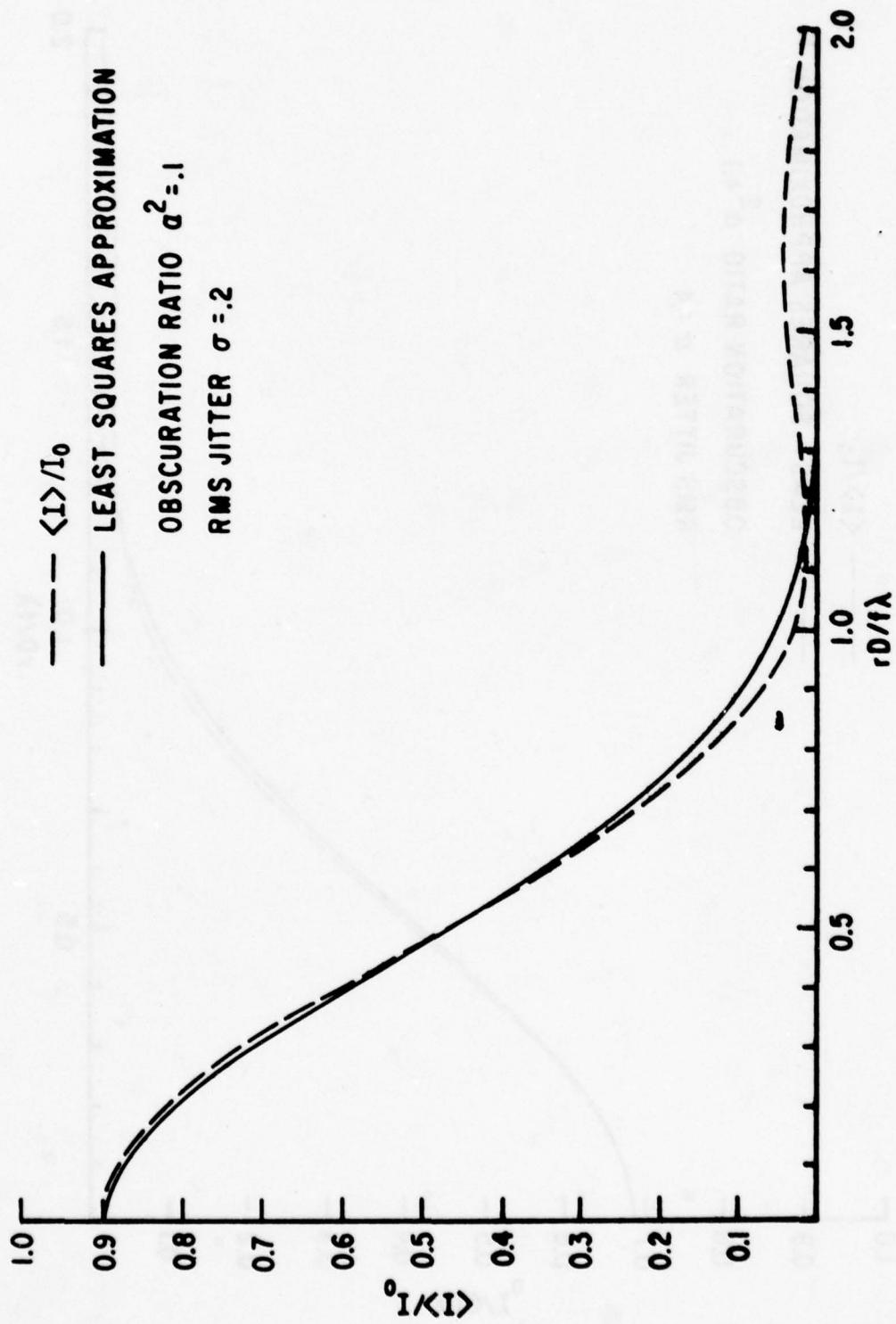


Figure 5(b)-1. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.



**Figure 5(b)-2.** Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

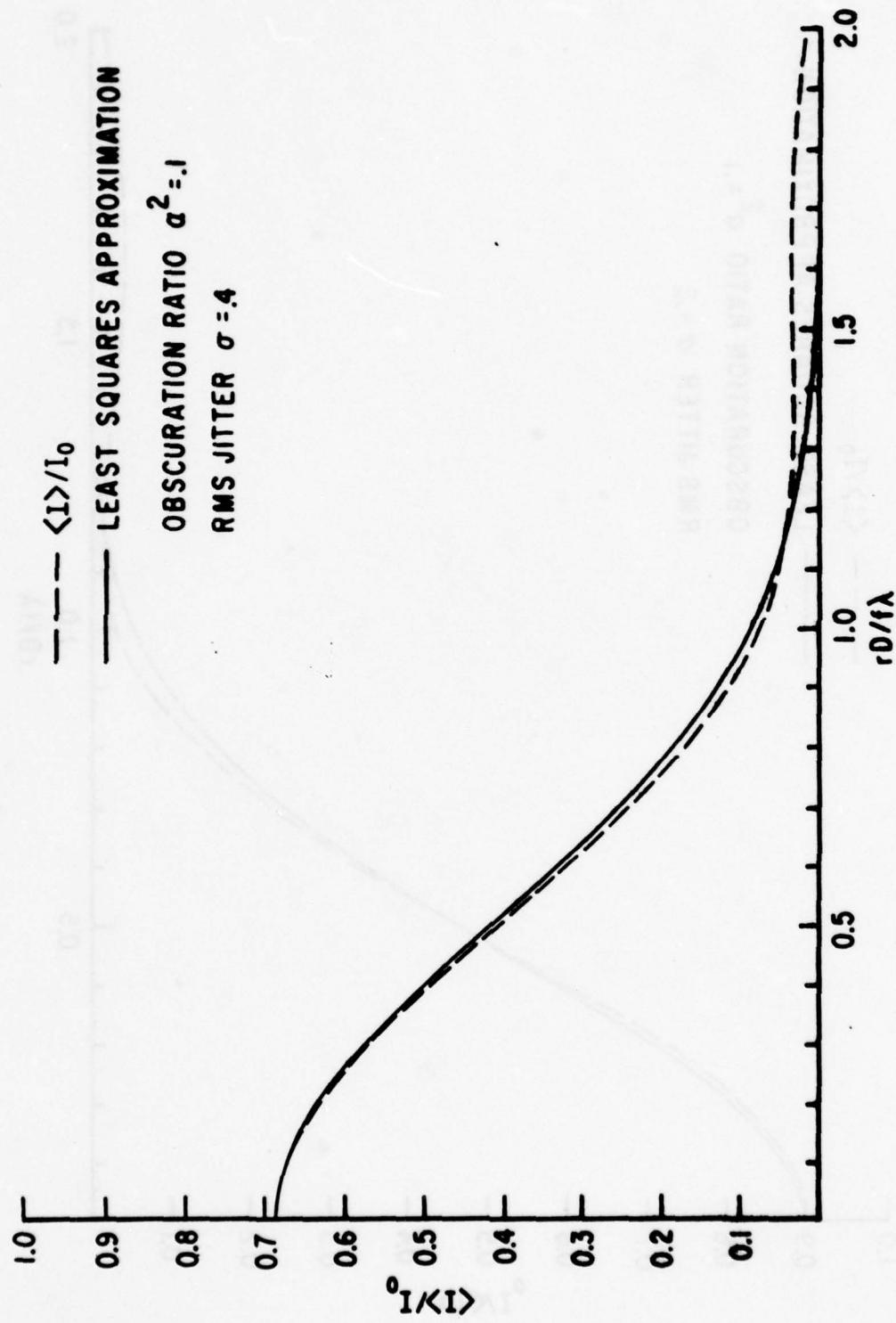


Figure 5(b)-3. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

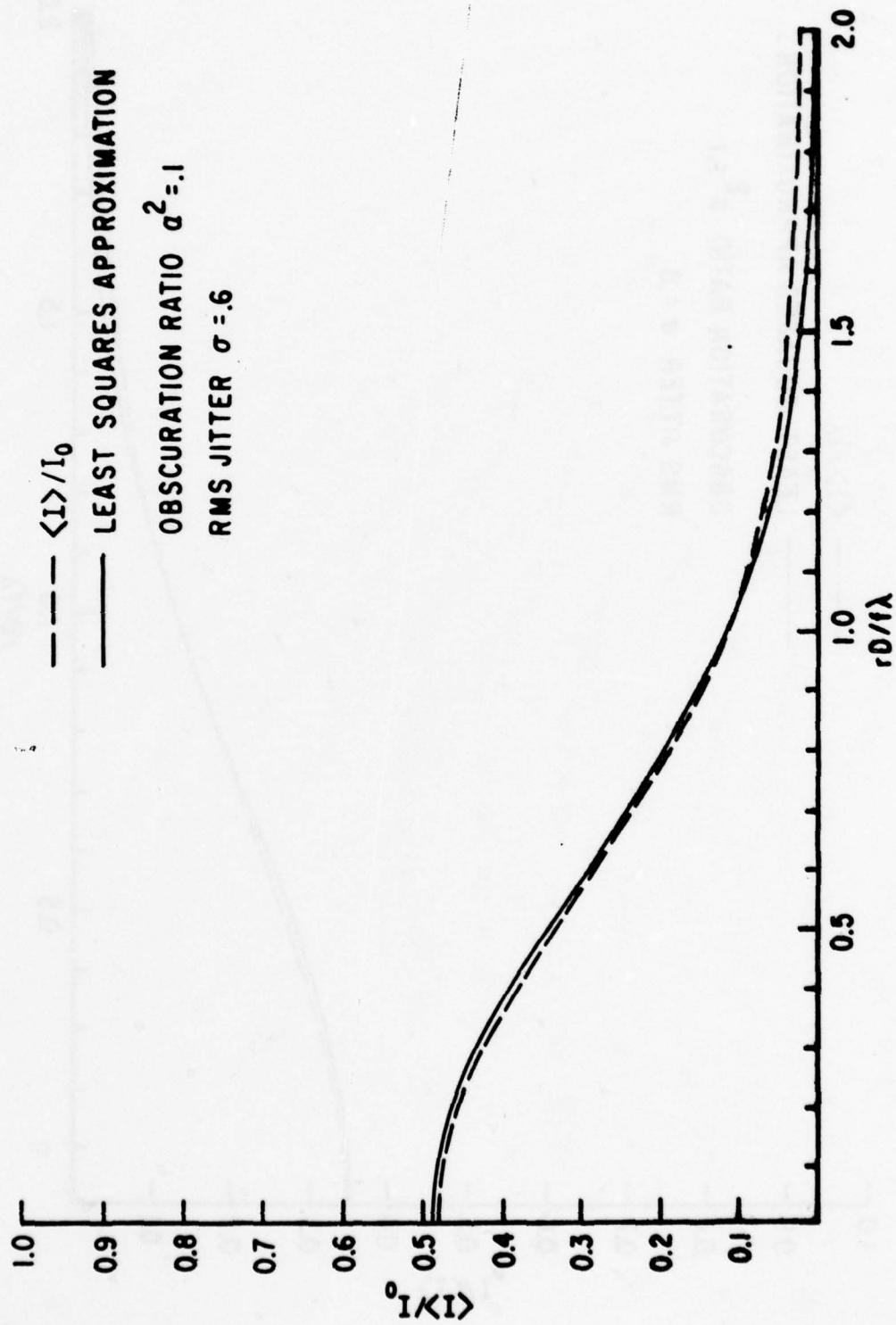
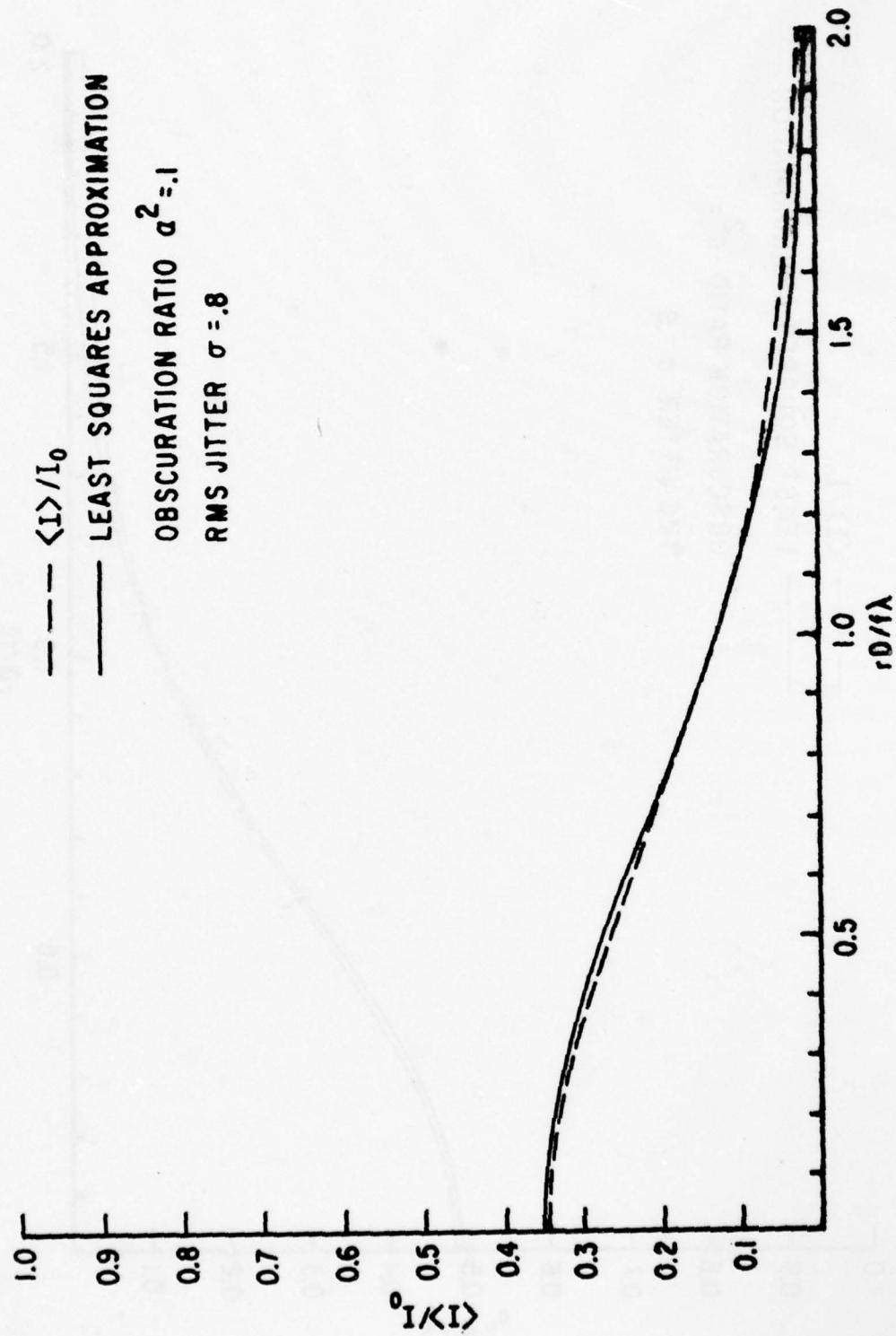


Figure 5(b)-4. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.



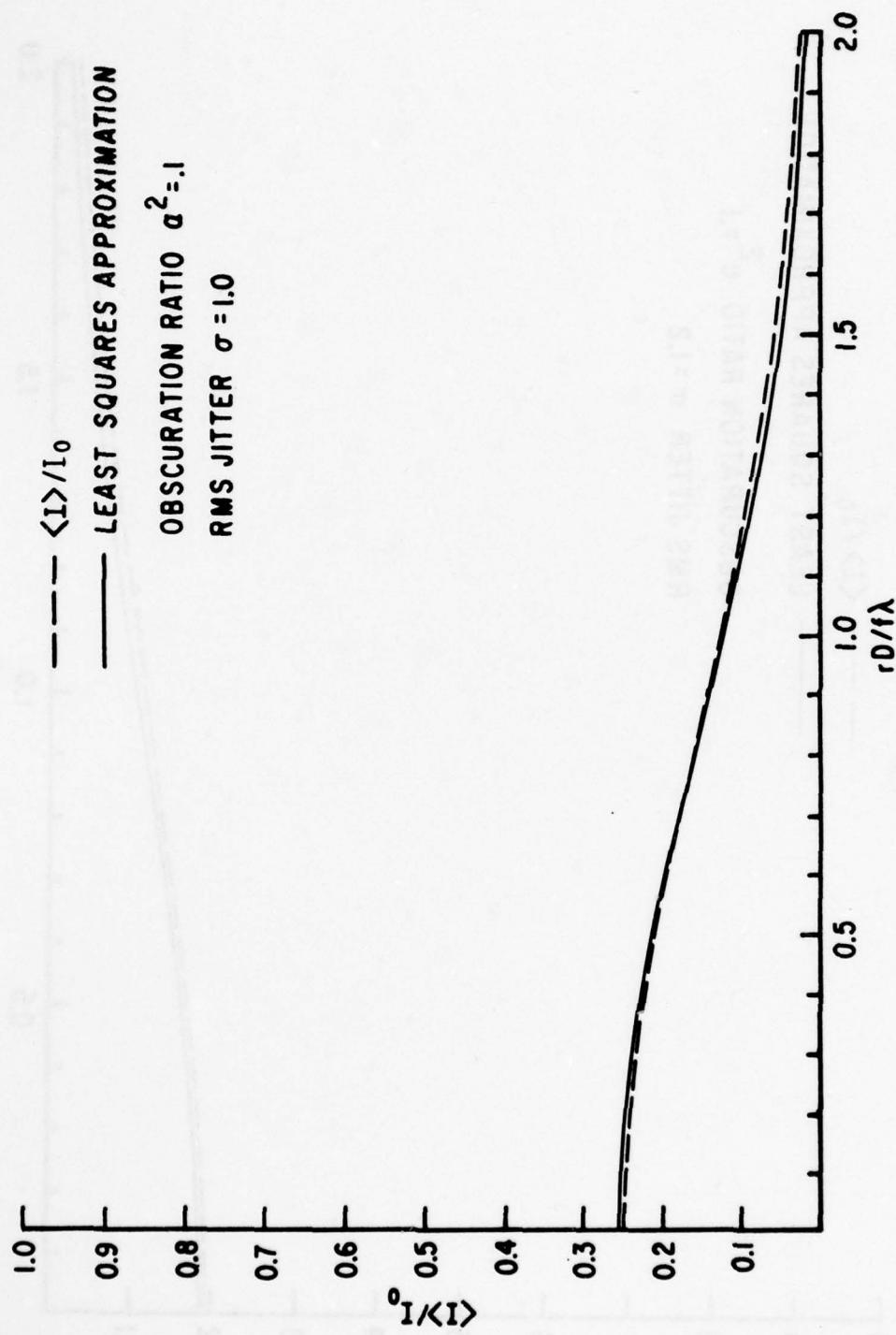


Figure 5(b)-6. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

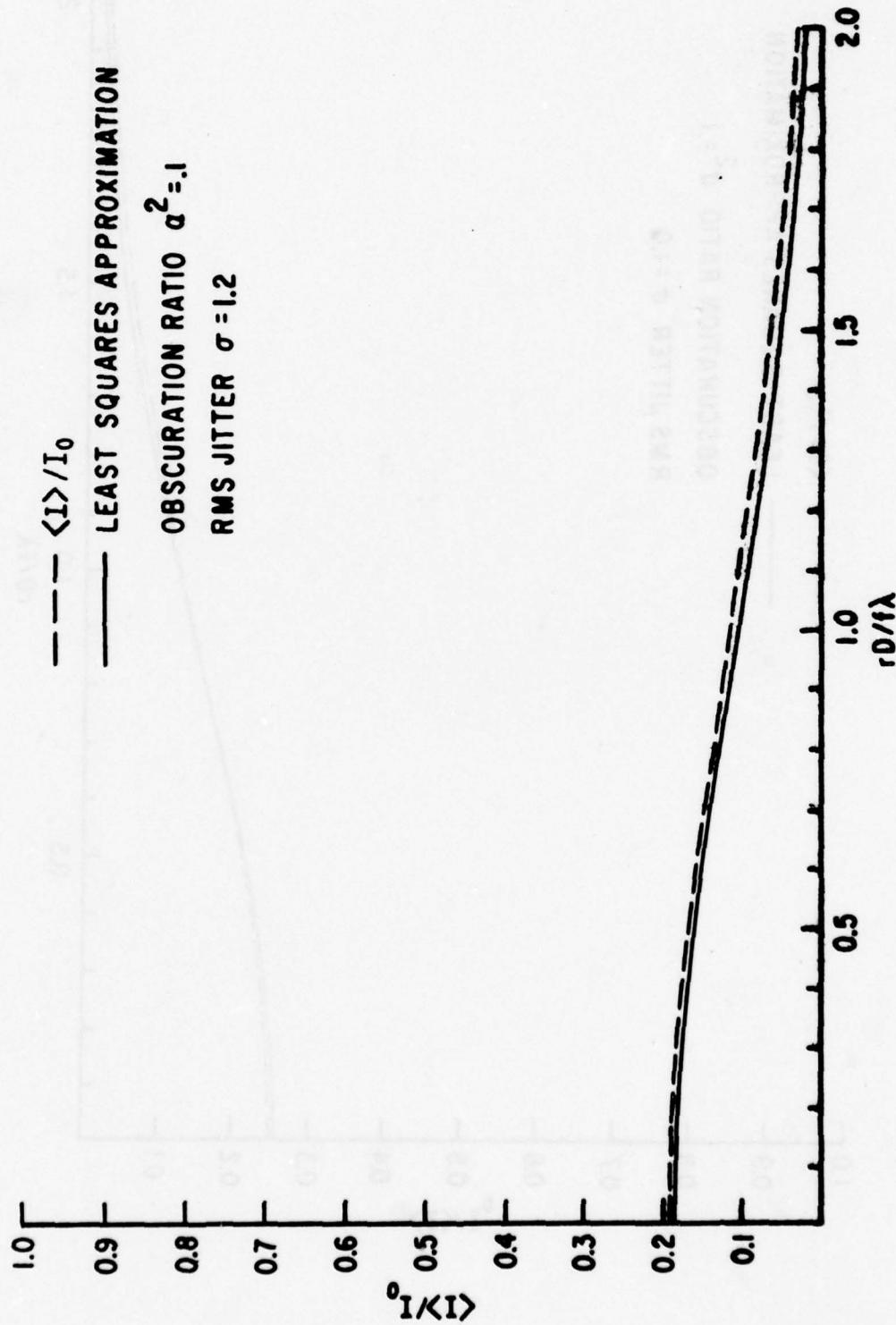


Figure 5(b)-7. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

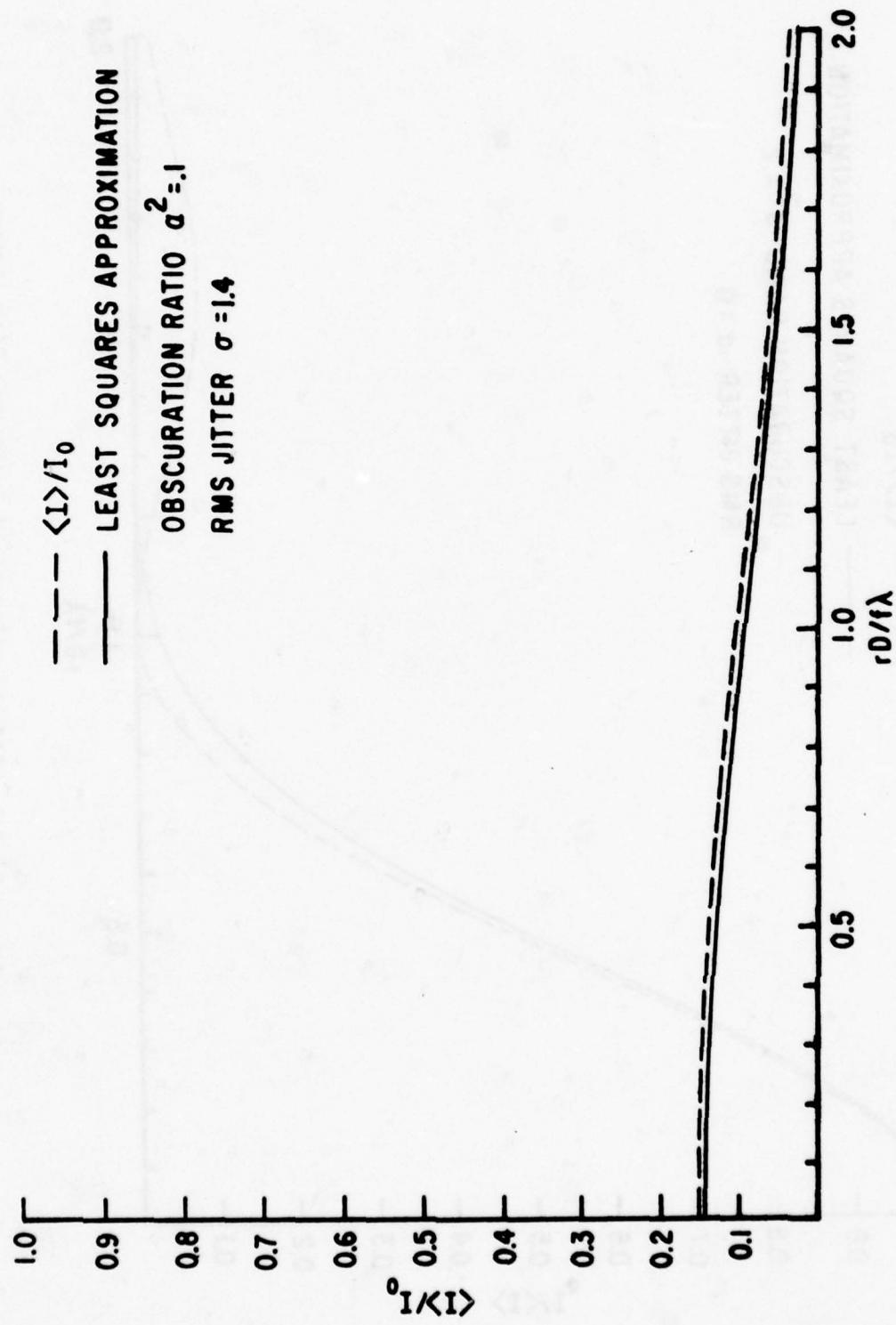


Figure 5(b)-8. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

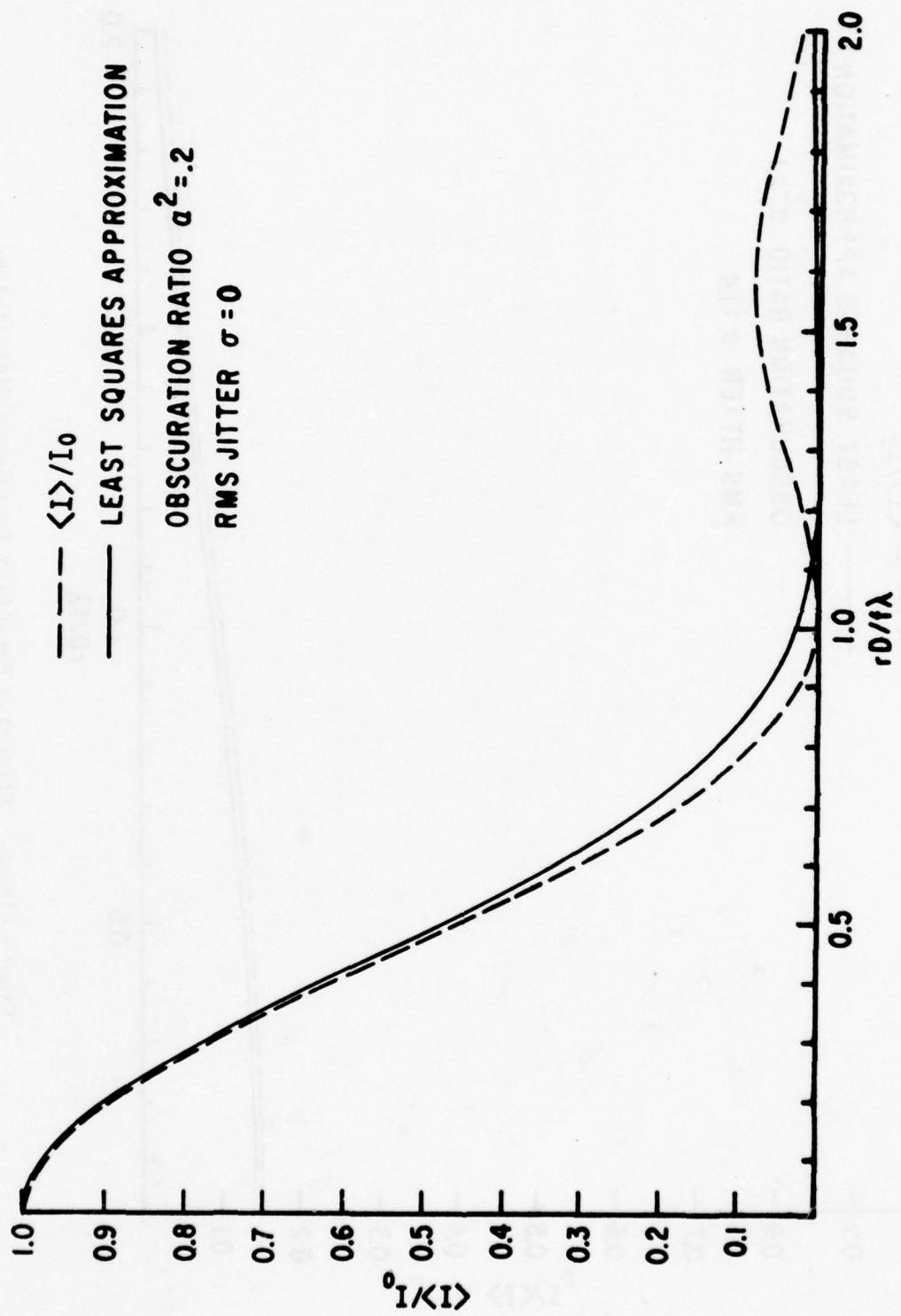


Figure 5(c)-1. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

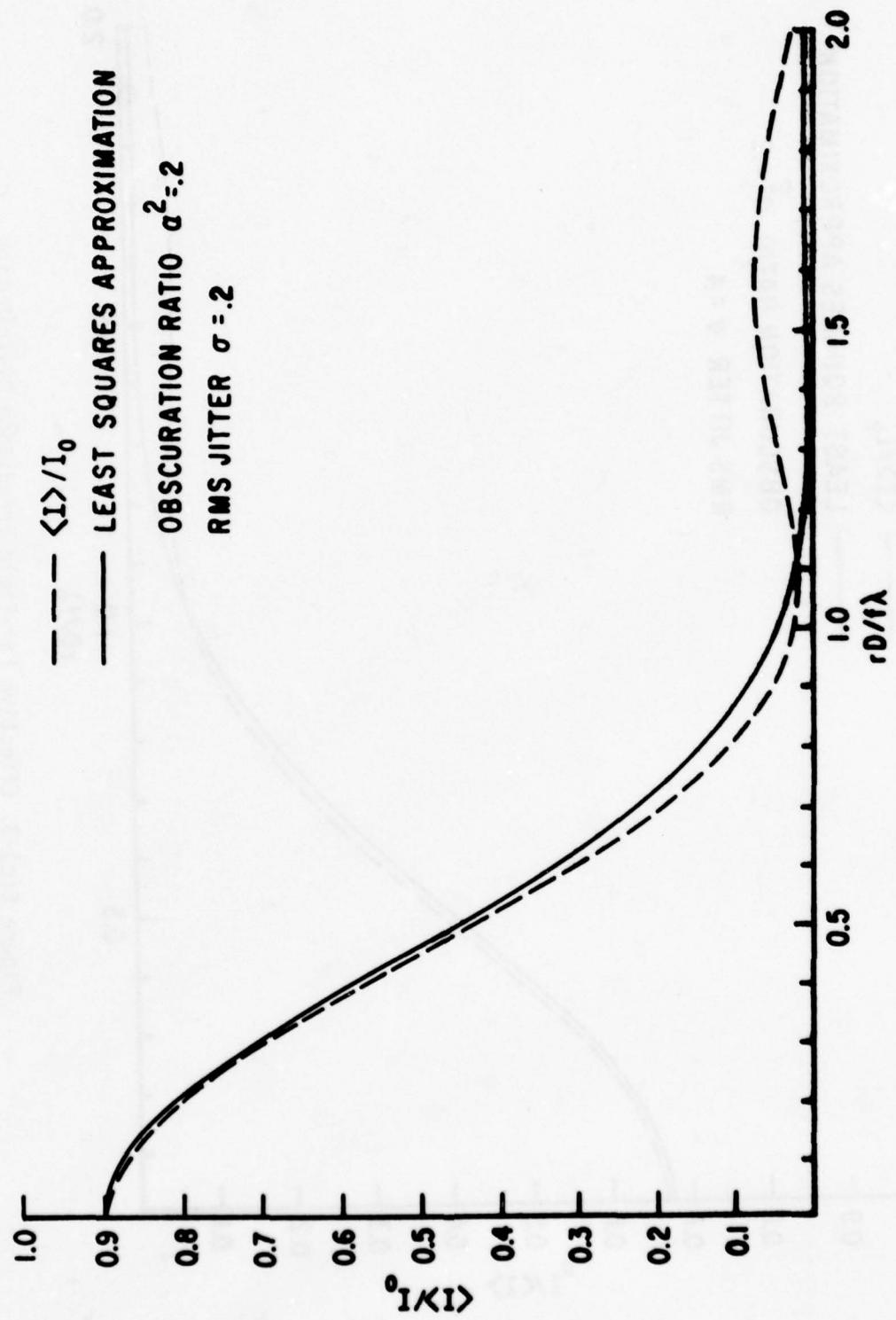


Figure 5(c)-2. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

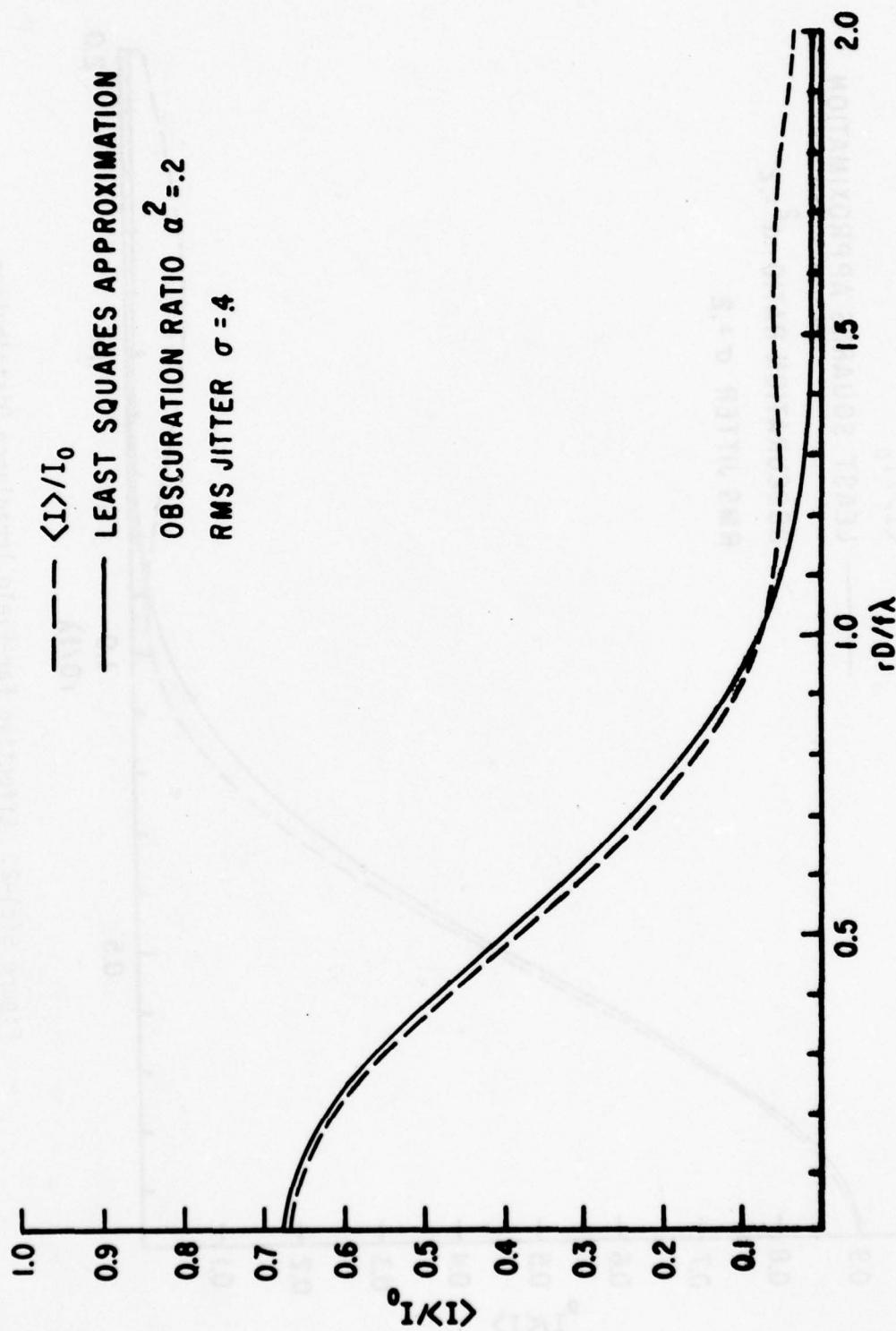


Figure 5(c)-3. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

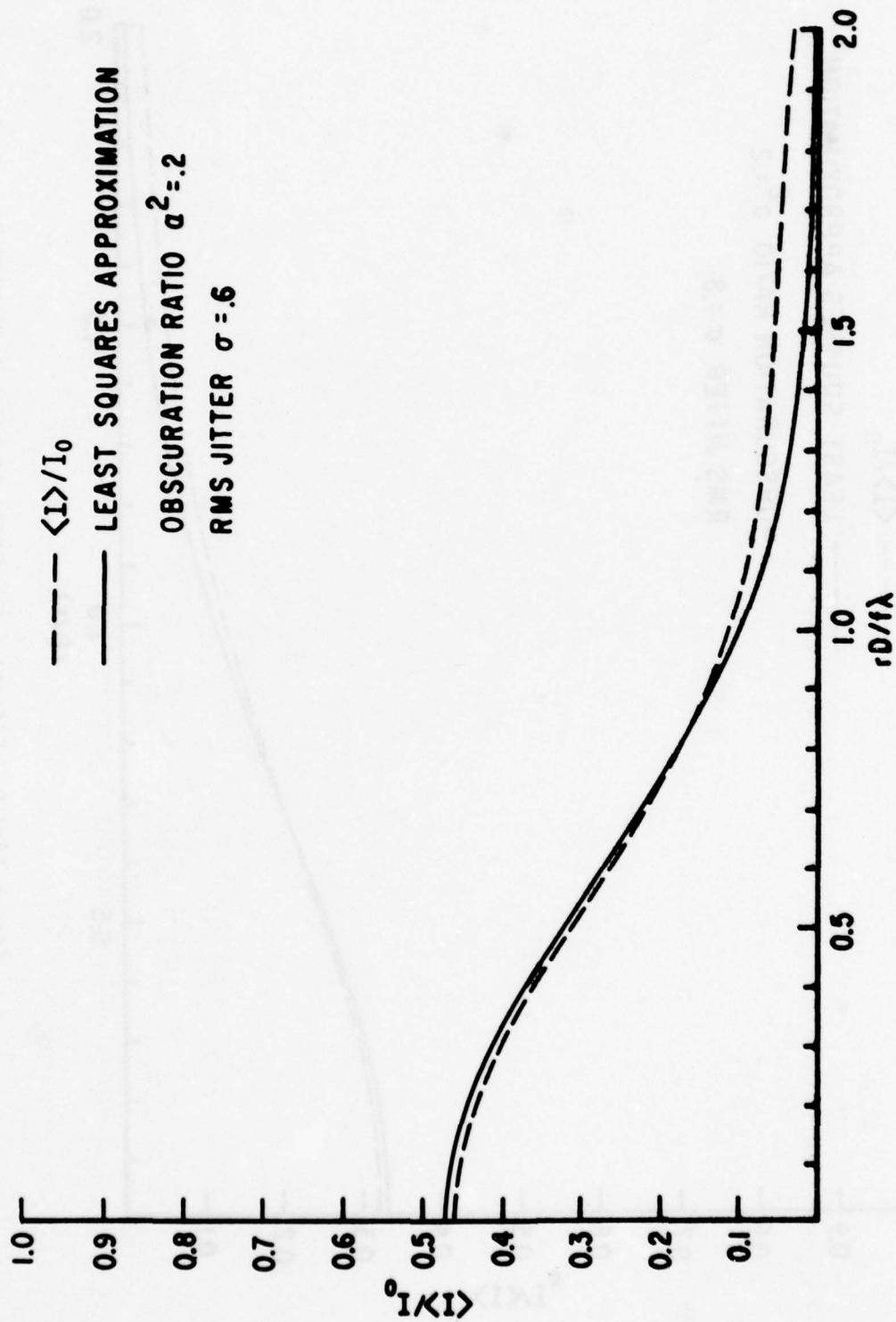


Figure 5(c)-4. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

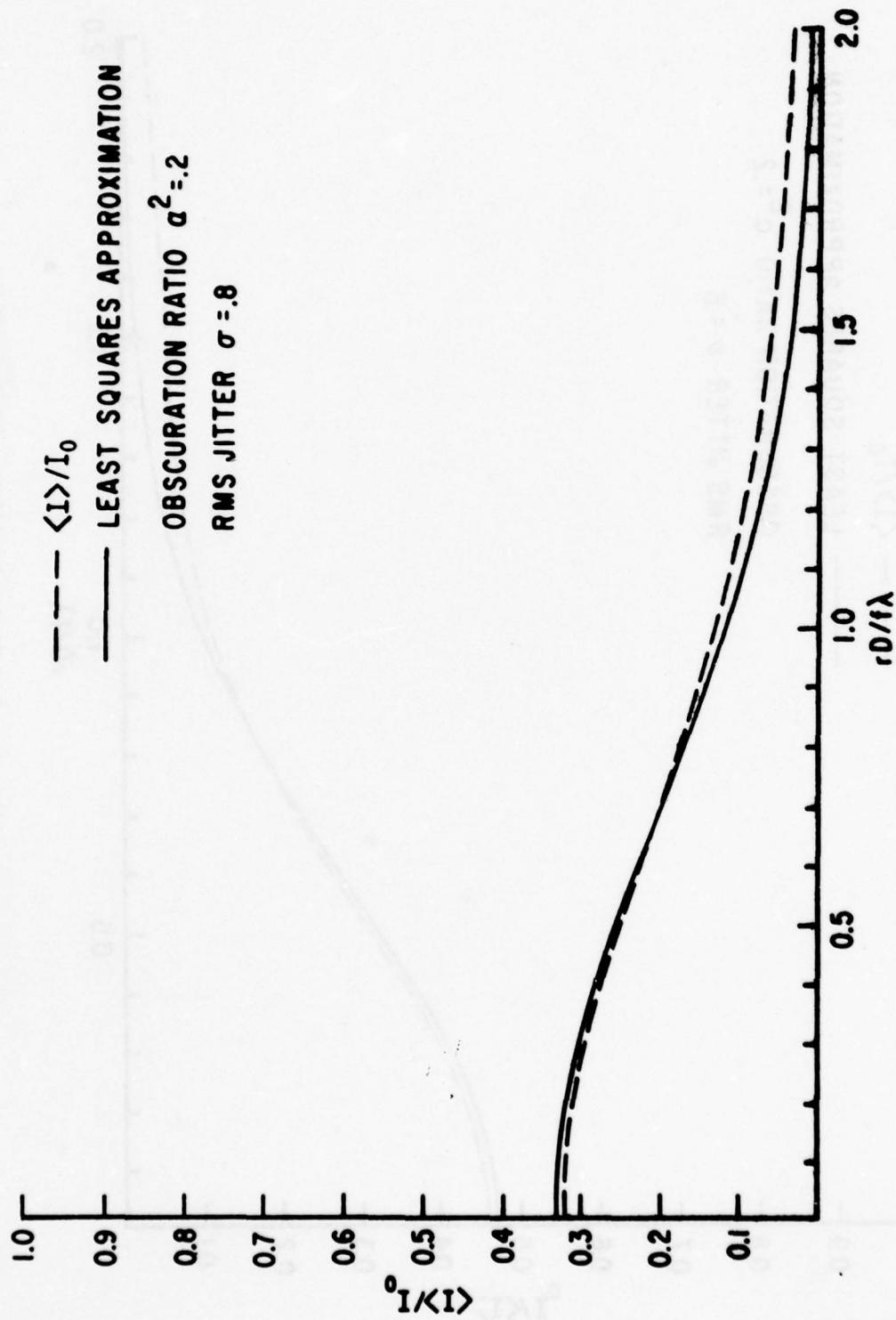


Figure 5(c)-5. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

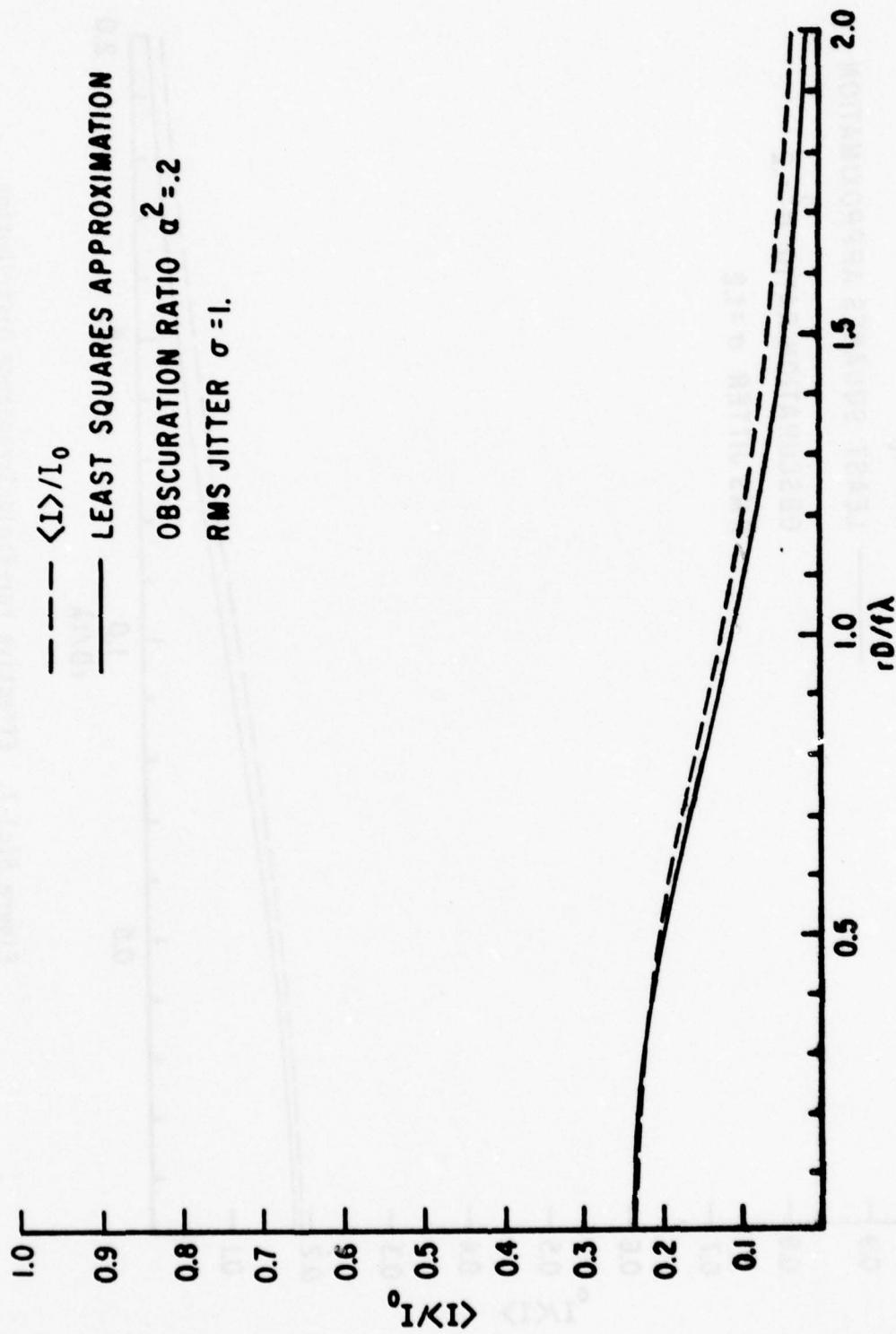


Figure 5(c)-6. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

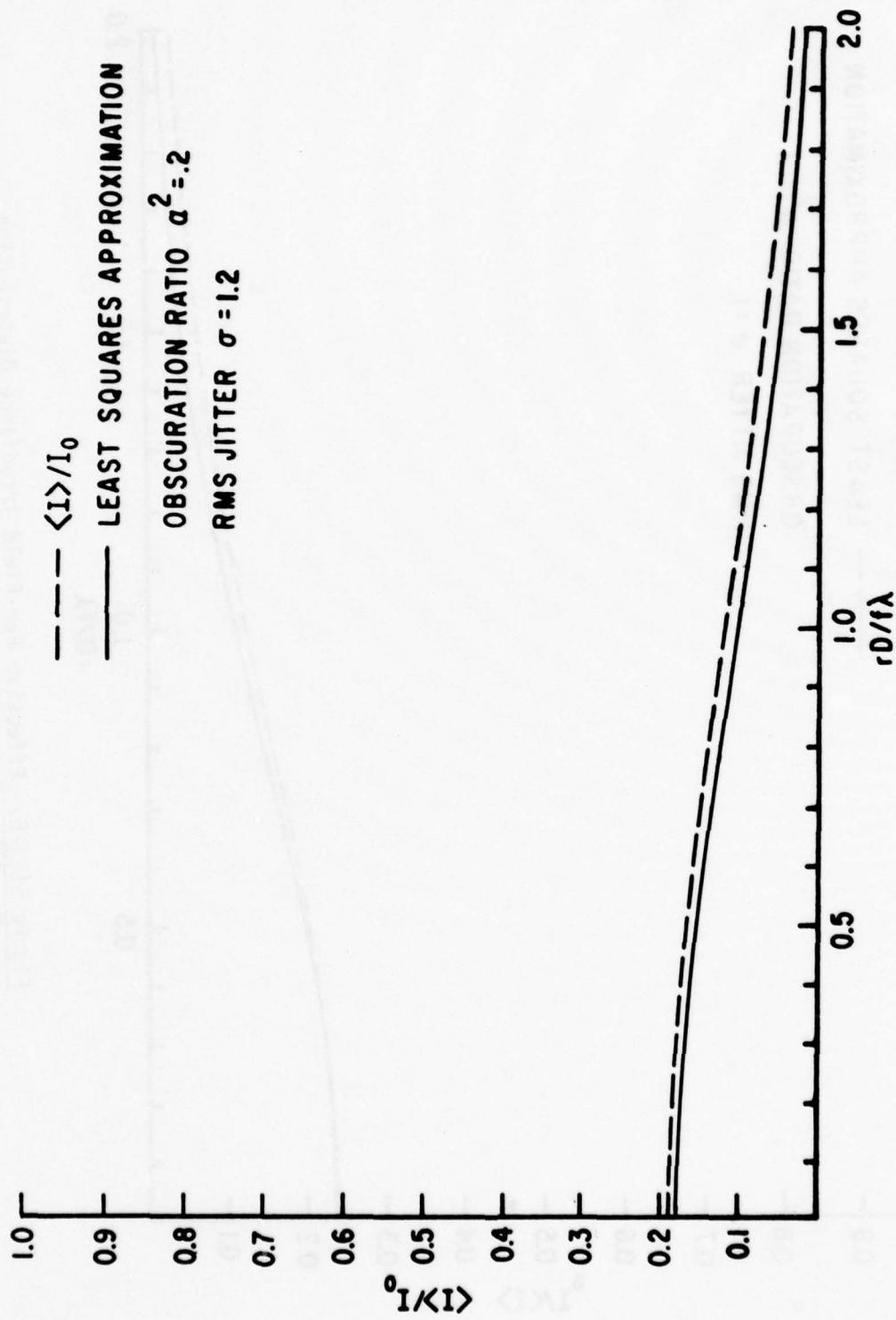


Figure 5(c)-7. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

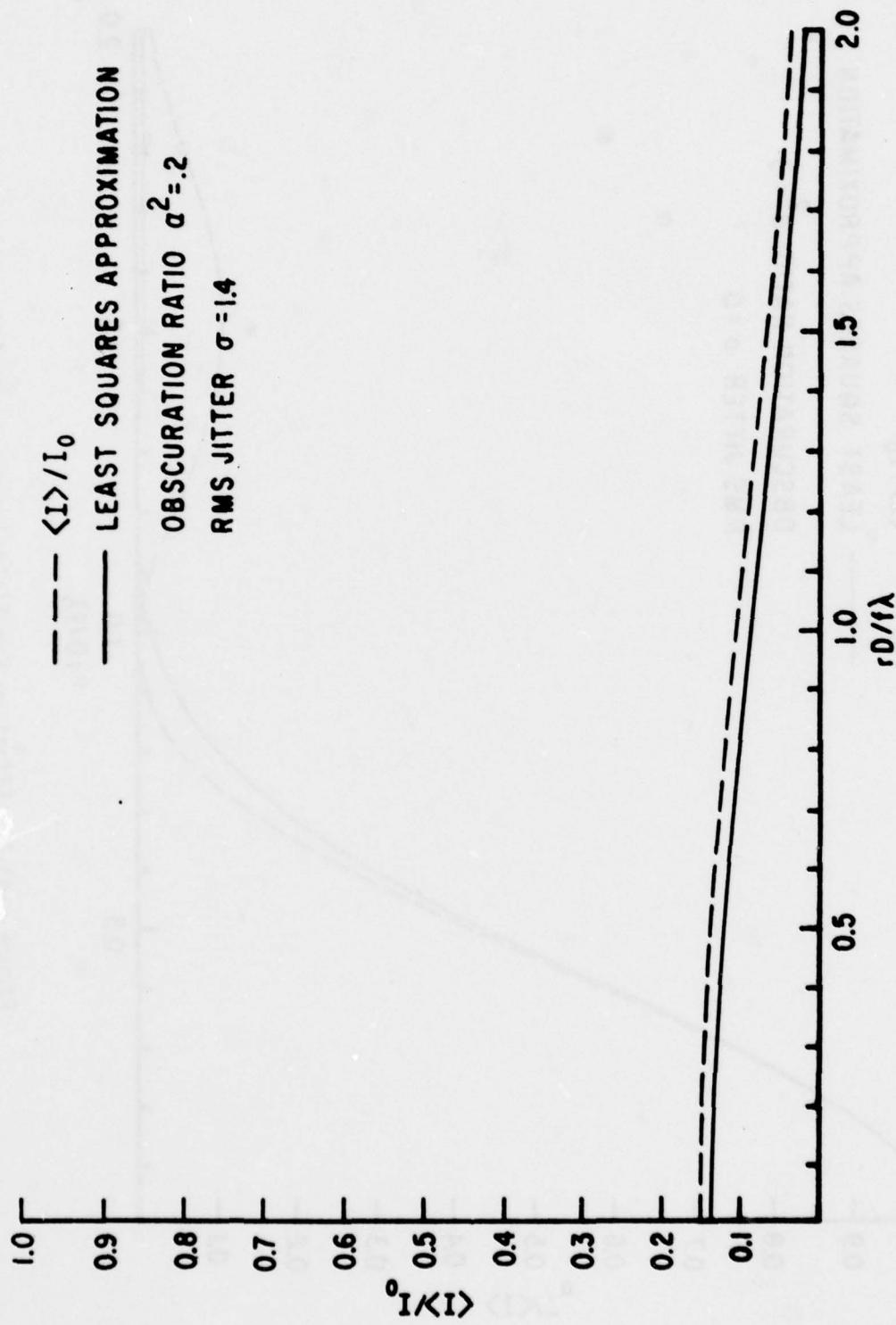


Figure 5(c)-2. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

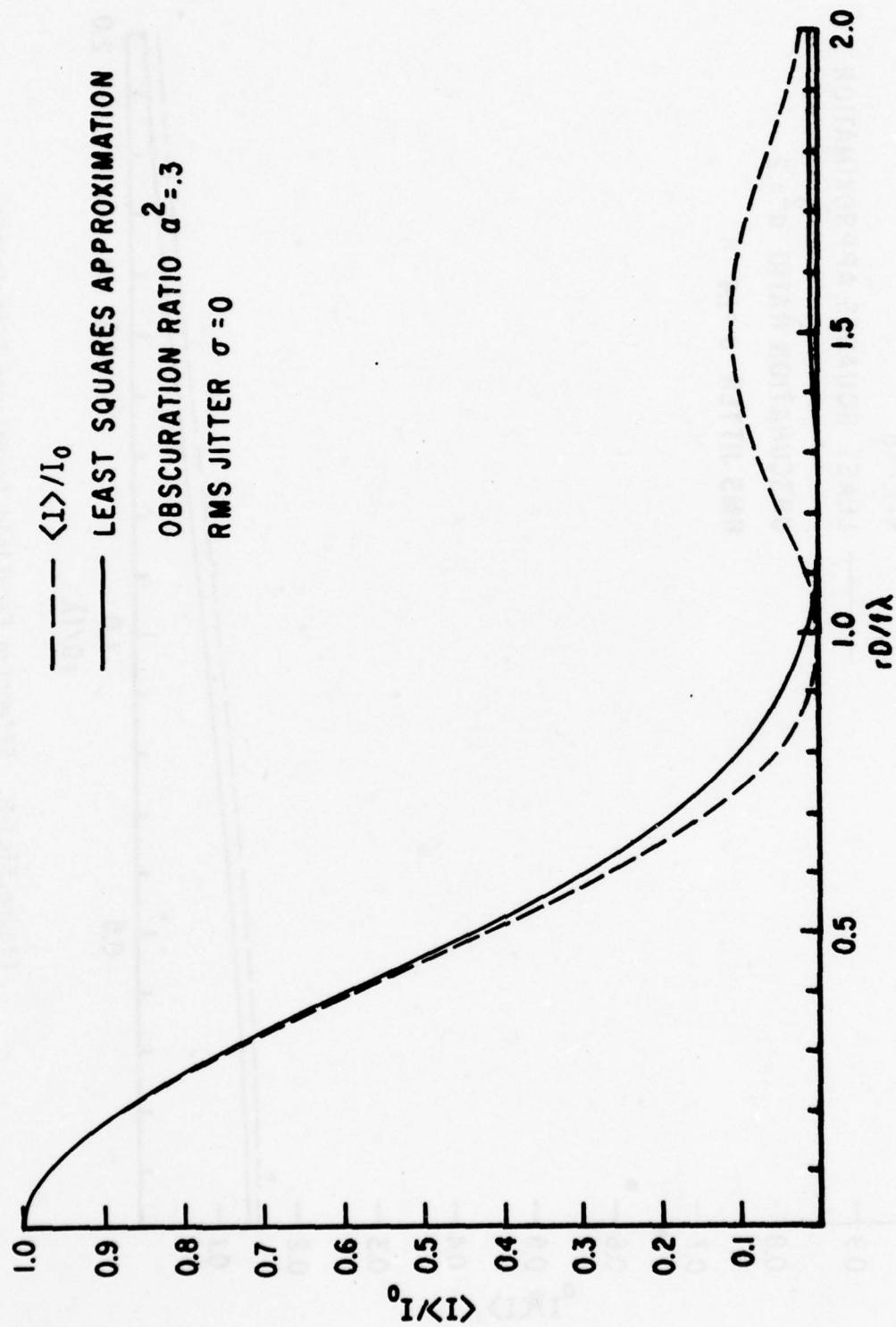


Figure 5(d)-1. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

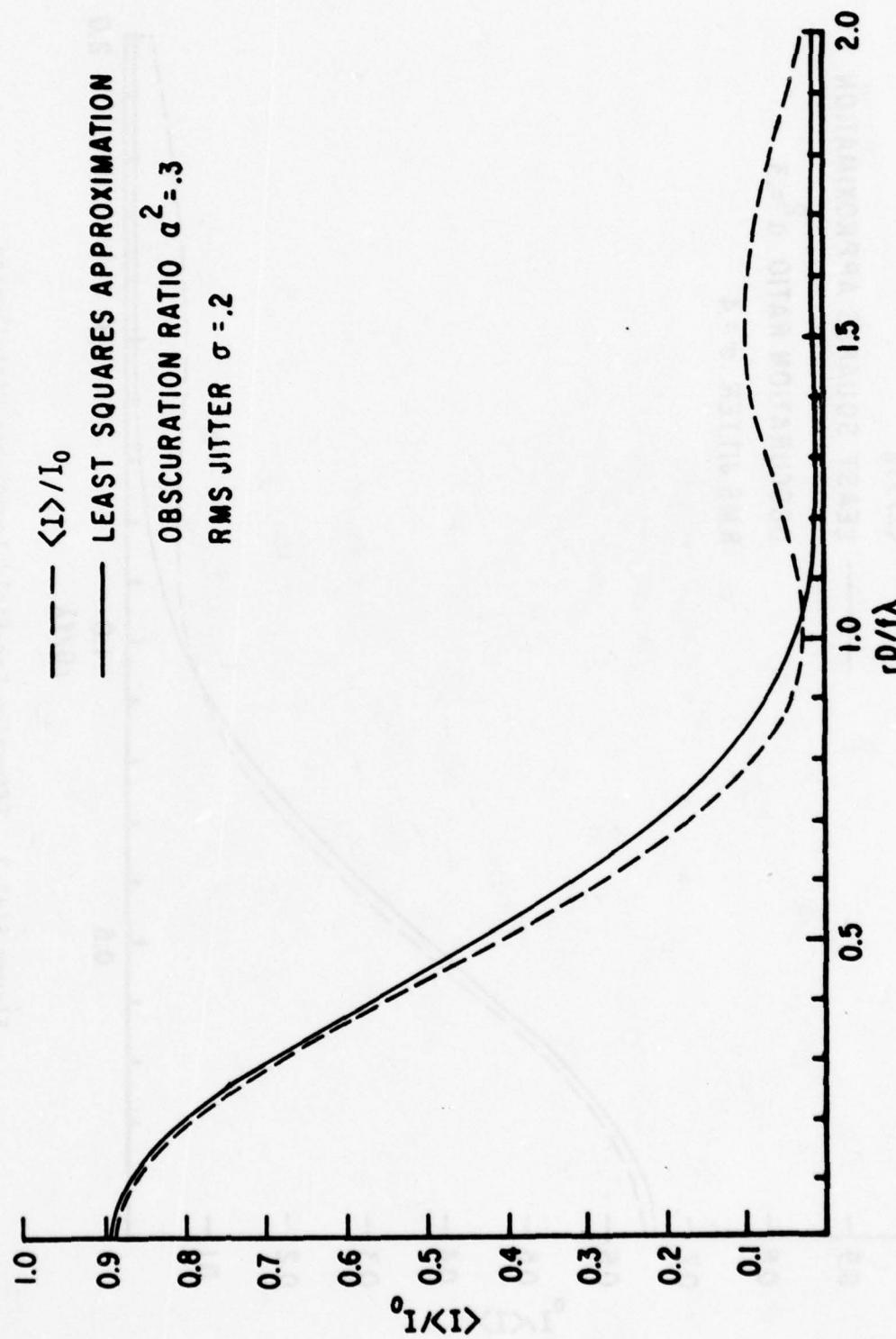


Figure 5(d)-2. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

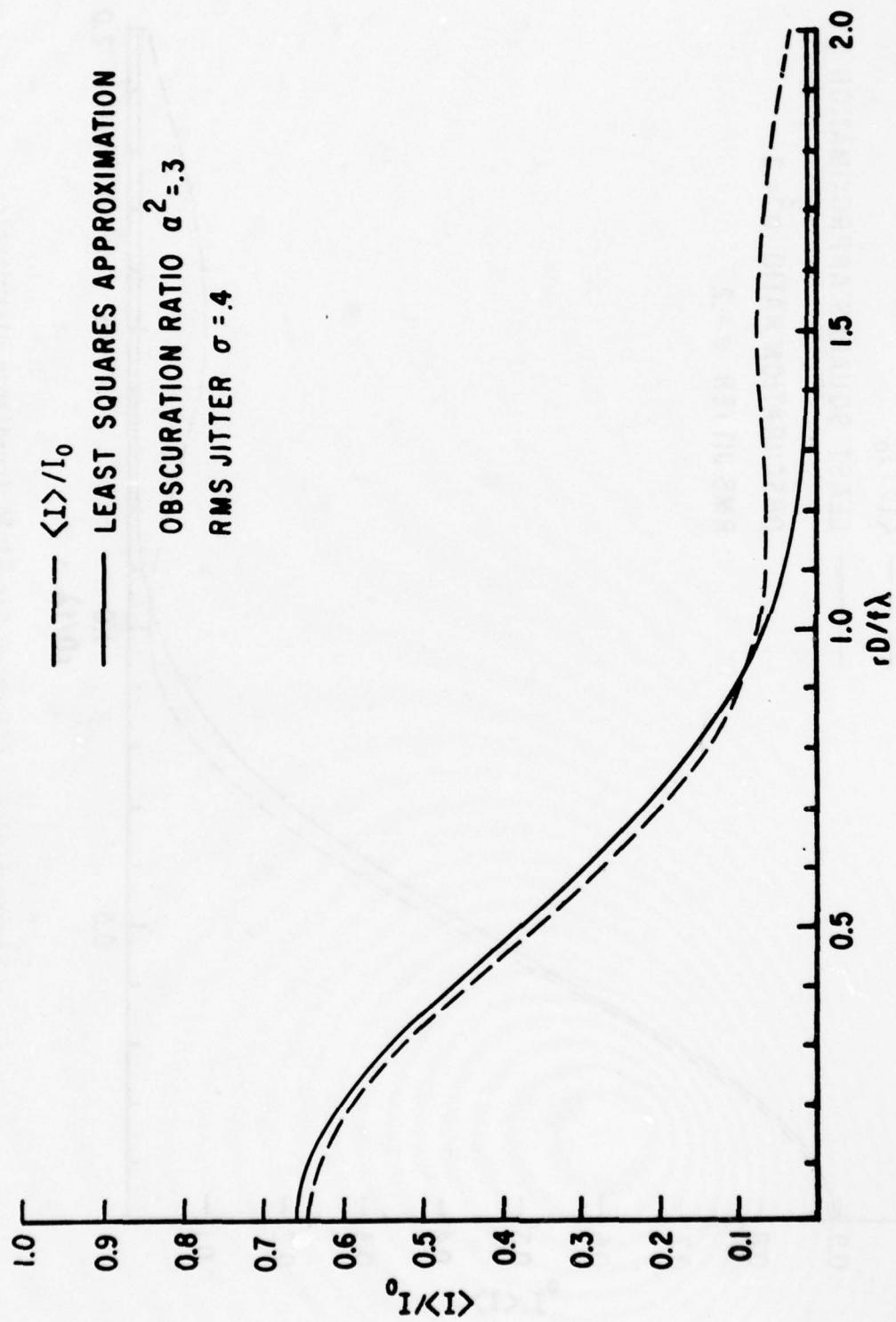


Figure 5(d)-3. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

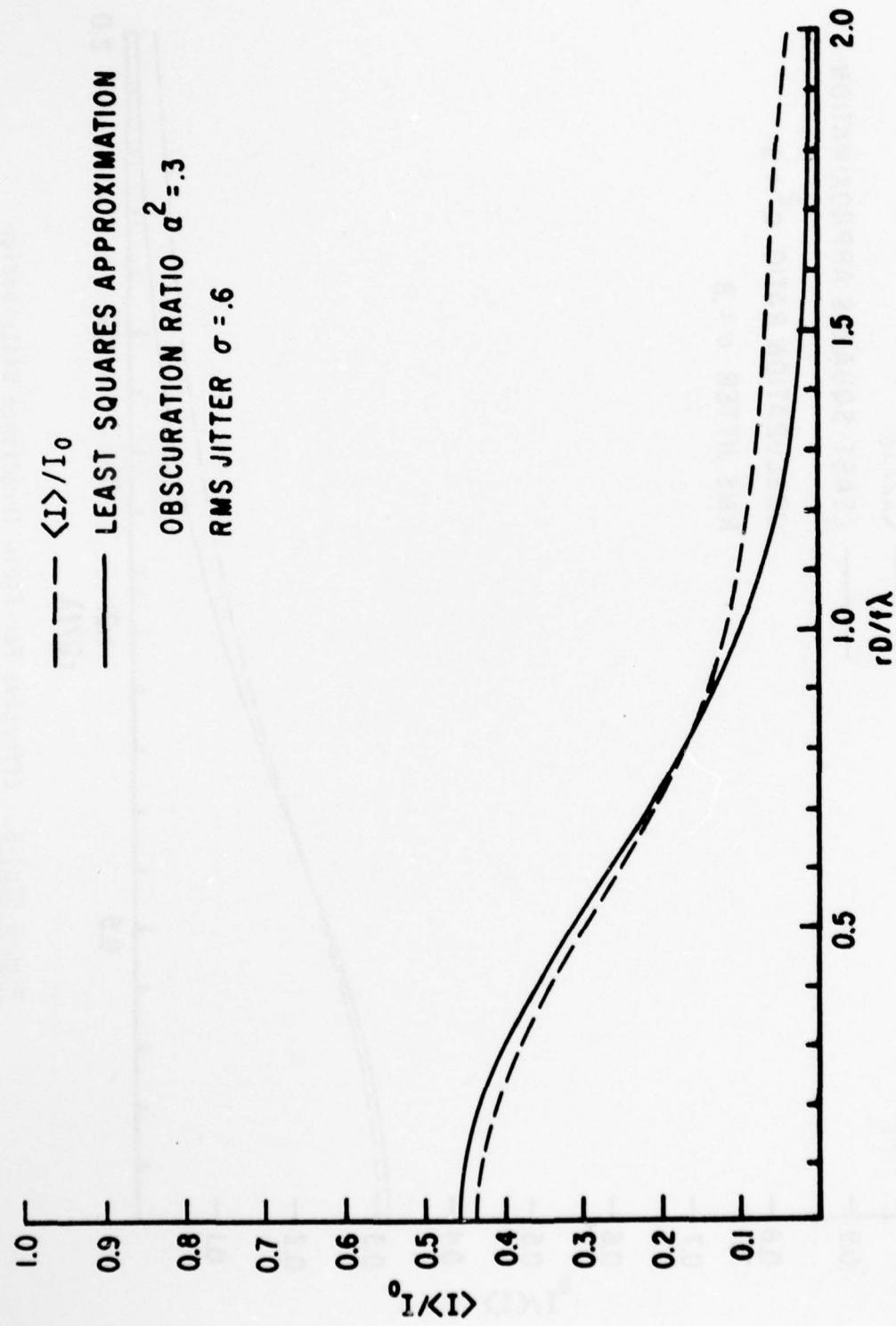


Figure 5(d)-4. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

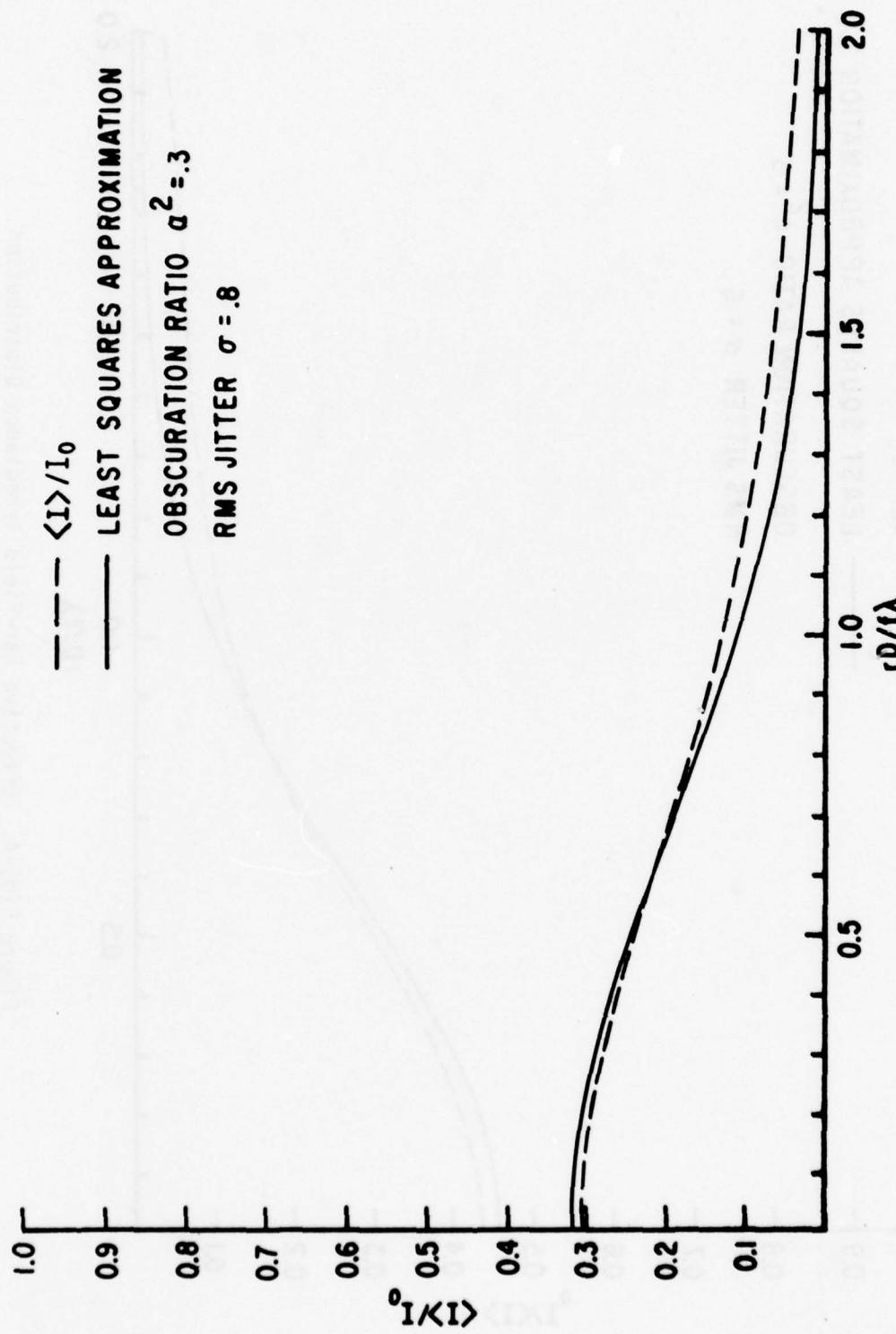


Figure 5(d)-5. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

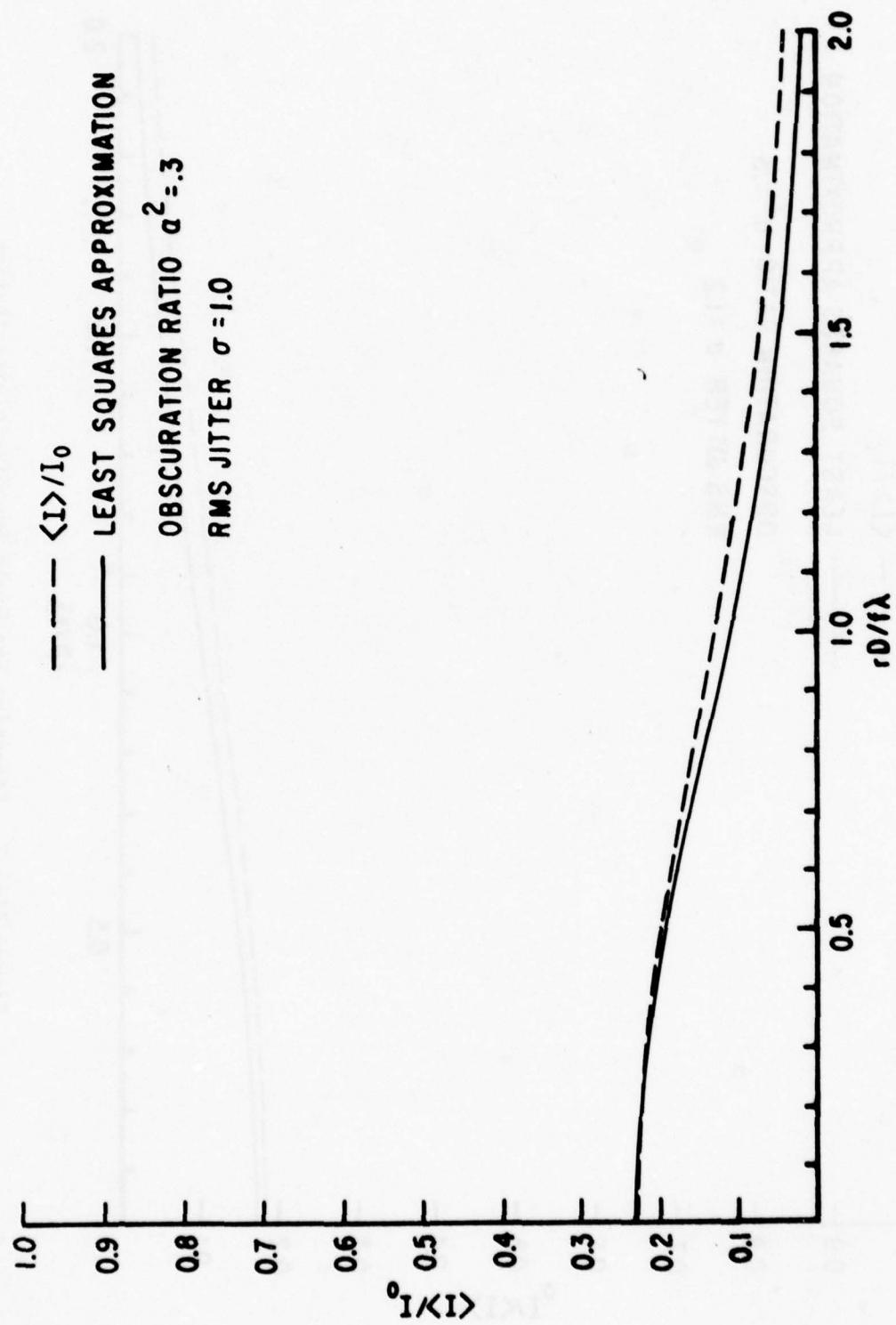


Figure 5(d)-6. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

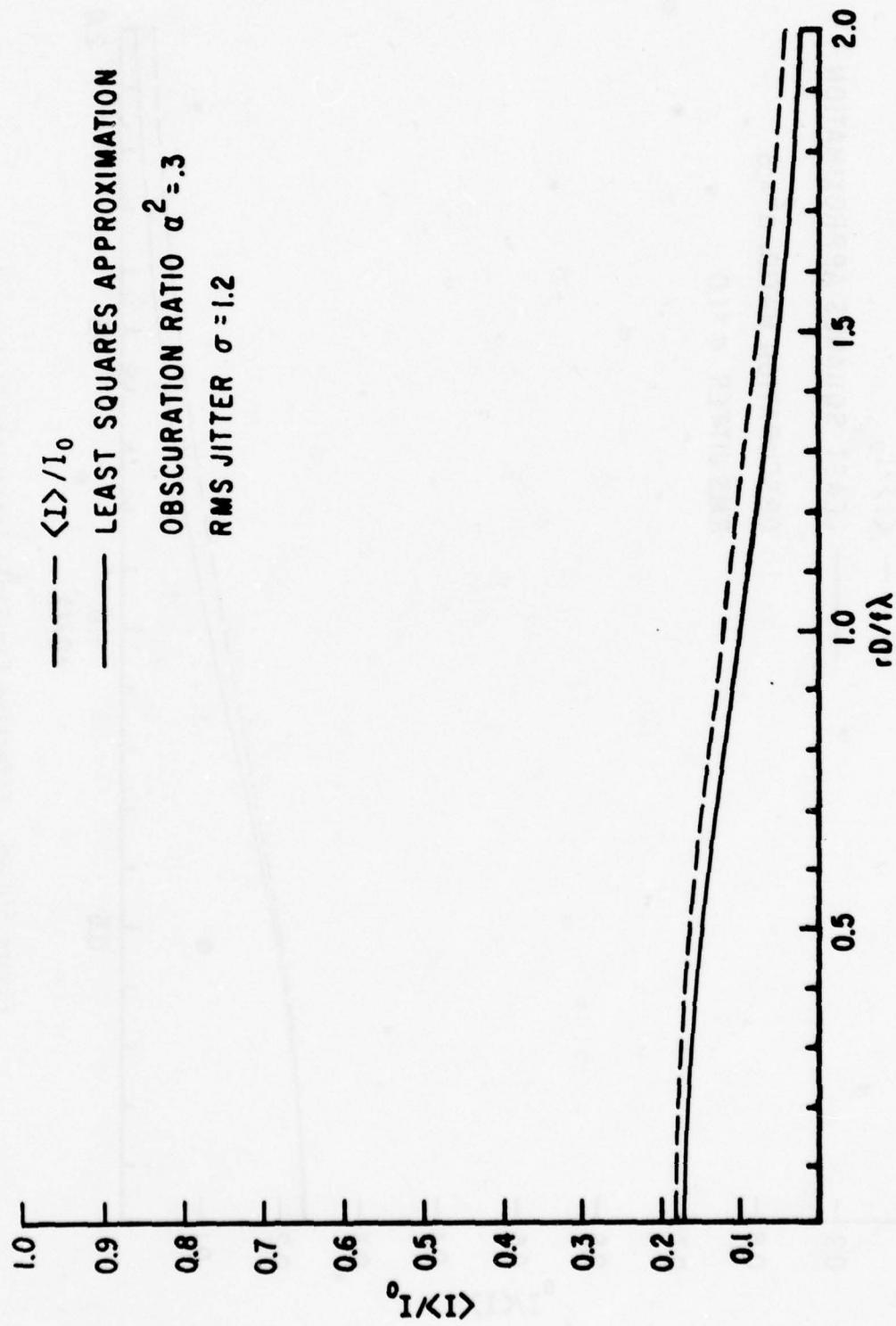


Figure 5(d)-7. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

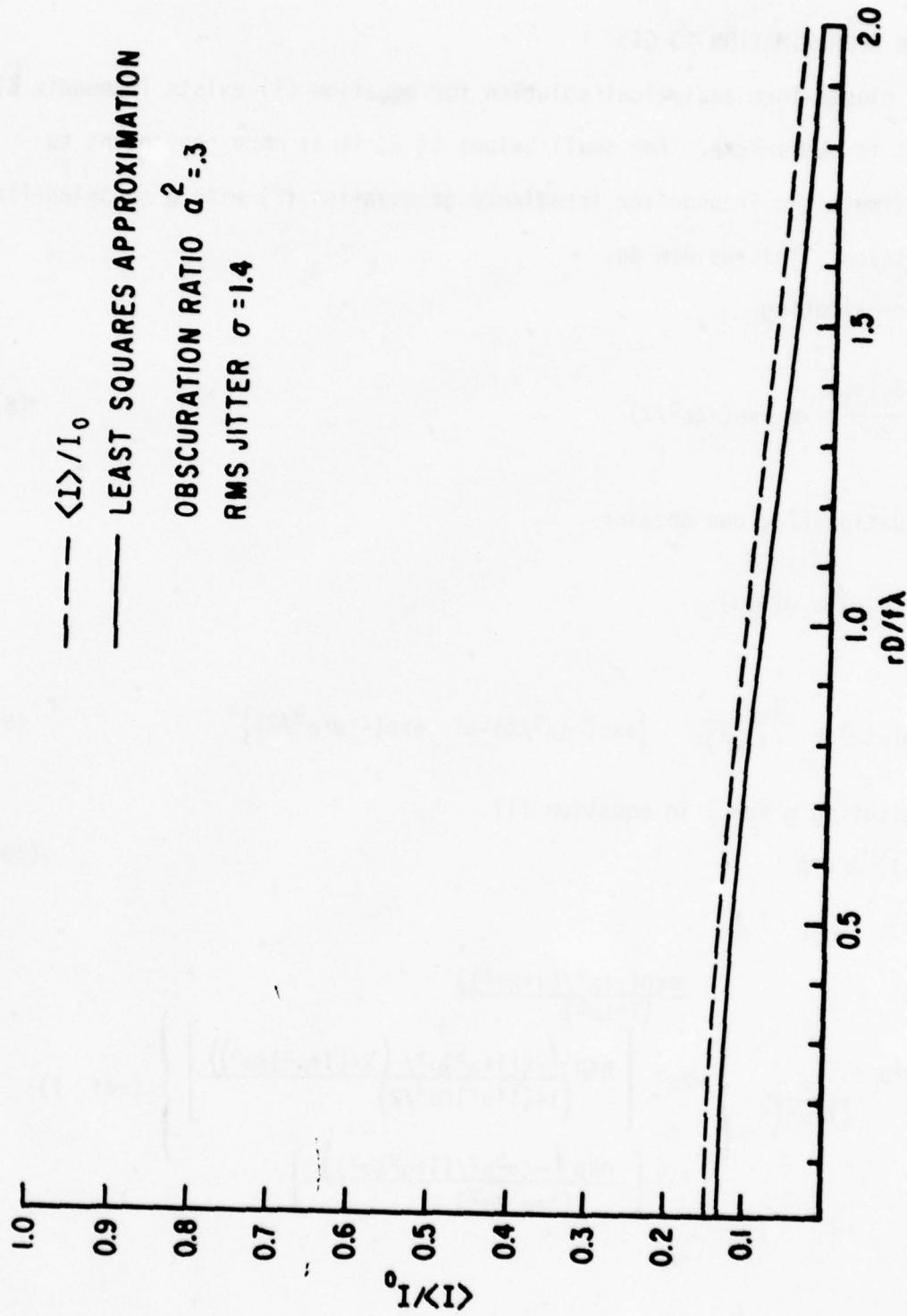


Figure 5(d)-8. Effective Far-Field Irradiance Distribution as a Function of Far-Field Radial Coordinate.

2. AN APPROXIMATION TO  $\langle I \rangle$ 

A closed-form analytical solution for equation (1) exists (Appendix B), but it is cumbersome. For small values of  $\alpha$ , it is more convenient to approximate the Fraunhofer irradiance of equation (2) with a gaussian-like expression. This we now do.

By approximating

$$\frac{2J_1(\pi\rho)}{\pi\rho} \approx \exp(-\xi\rho^2/2) \quad (8)$$

in equation (2), one obtains

$$I(\rho; \alpha) \approx g(\rho; \alpha)$$

where

$$g(\rho; \alpha) = \frac{I_o}{(1-\alpha^2)^2} \left\{ \exp(-\xi\rho^2/2) - \alpha^2 \exp(-\xi\alpha^2\rho^2/2) \right\}^2 \quad (9)$$

Substituting  $g$  for  $I$  in equation (1),

$$\langle I \rangle \approx g^*p \quad (10)$$

$$g^*p = \frac{I_o}{(1-\alpha^2)^2} \left\{ \begin{aligned} & \frac{\exp(-\xi\rho^2/(1+\xi\alpha^2))}{(1+\xi\alpha^2)} \\ & -2\alpha^2 \left[ \frac{\exp(-\xi(1+\alpha^2)\rho^2/(2+(1+\alpha^2)\xi\alpha^2))}{1+(1+\alpha^2)\xi\alpha^2/2} \right] \\ & + \alpha^4 \left[ \frac{\exp(-\xi\alpha^2\rho^2/(1+\alpha^2\xi\alpha^2))}{1+\alpha^2\xi\alpha^2} \right] \end{aligned} \right\} \quad (\text{ref. 1})$$

Physically,  $g$  represents the interference pattern of two focused gaussian beams whose beam waists are  $w$  and  $w/\alpha$ , where

$$w^2 = \left(\frac{f\lambda}{D}\right)^2 2/\xi \quad (11)$$

The parameter  $\xi$  is taken to be a function of  $\alpha$  only. By setting  $\rho = 0$  and least-squares fitting equation (10) to the digital computer solutions of equation (7), one obtains

$$\xi \approx 2.724 - 1.193\alpha^2 + .9353\alpha^4 \quad (12)$$

Equation (9) is a useful approximation for  $\langle I \rangle$  when one wishes to find the effect of jittered, uniformly intense, annular "plane waves" where the annulus has a value of  $\alpha$  other than those considered in this paper.

## SECTION VI

### CONCLUSIONS

Jitter is one major cause of laser beam degradation. In effect it broadens an otherwise diffraction-limited beam, causing a corresponding decrease in peak irradiance and washing out interference fringes.

An increase in jitter usually but not always decreases the time-averaged power transmitted by a circle that is centered at the origin of the Fraunhofer plane. If the rms gaussian jitter is small by comparison to the normalized radius of the circle, i.e., if

$$\frac{aD}{f\lambda} \gg \sigma \quad (12)$$

then the effect of jitter upon the power transmitted by a circle of radius  $a$  is negligible within a tolerance of a few percent.

From such a criterion it is apparent that as the radius of the circle decreases to zero, it is impossible to have the effect of jitter be negligible. The power transmitted by a tiny circle centered at the point of time-averaged peak irradiance is highly sensitive to jitter. Normalizing the transmitted power by the area of the circle, one has, in the limit of vanishing area, the value of the peak irradiance. Consequently the peak irradiance falls off much faster as jitter increases than does the irradiance farther off axis. Measurements with detectors whose diameters are much smaller than the nominal beam width are quite sensitive to jitter.

From a viewpoint of raw energy transfer, large area receivers are

insensitive to jitter. Based on Figure 2, a circle whose radius,  $a$ , is approximately  $f\lambda/D$  can tolerate an rms jitter of about  $.5 f\lambda/D$  with about 10% loss of transmitted energy. A circle of twice the diameter can tolerate an rms jitter of about  $1 f\lambda/D$ . A circle of  $a \approx 3 f\lambda/D$  can tolerate about  $1.5 f\lambda/D$ . Hence, a rule of thumb for large area circles is that if the radius of the circle equals  $a$ , then an rms jitter of  $a/2$  will induce about 10% loss of power transmitted by the circle. This rule of thumb seems to be fairly independent of the annular obscuration ratio  $\alpha^2$  for  $\alpha^2 \lesssim .3$ .

From a standpoint of beam quality a given amount of jitter induces a larger drop in central intensity than it does in encircled irradiance. (See figures 4(a) - 4(c).) To anyone concerned with specifying jitter tolerances of a focusing optical system, a specific limitation on the degradation of the central irradiance, say 10% will place a tighter restriction on jitter than the same (10%) degradation limitation on encircled power.

This paper has considered only a few values of the obscuration ratio  $\alpha^2$ . By using equation (8) it is easy to estimate the effect of jitter upon the Fraunhofer irradiance distribution for other obscuration ratios.

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APPENDIX A  
APPLICABILITY OF MODELING JITTER OF  
A FOCUSED BEAM AS A CONVOLUTION

Although it was assumed that jitter caused the Fraunhofer pattern to undergo simple, unaberrated translation over the entire plane surface, in practice this does not happen. Real jitter causes the focal spot to trace out a path that is not contained in a single plane. Also, no excursion is ever infinitely far from the nominal optical axis. In this article, we ignore all wavefront aberrations and assume our instantaneous irradiance distribution is always given by the Fraunhofer pattern, equation (2) of the main body of this paper. With this restriction in mind, we address the following two questions: First, what is the effect of assuming that jitter causes simple translation over a plane? Secondly, what is the error in assuming that jitter can cause the beam to wander infinitely far from the optical axis?

A simple model for jitter is depicted below. A stationary beam strikes a wiggling mirror.  $\phi_0$  represents the line-of-sight off-axis angle to the instantaneous center of the jittering irradiance distribution.

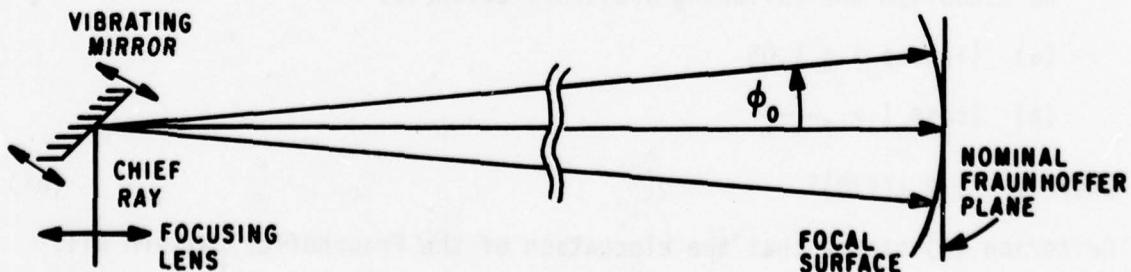


Figure A1. Simple Jitter Model

For large enough angles of jitter, two effects occur. First of all, the beam no longer focuses in the plane it originally focused in. Secondly, even if the depth of focus is sufficient to allow a Fraunhofer pattern to appear in the original focal plane, distortion occurs. The pattern is elongated because the pattern in the nominal Fraunhofer plane is really a projection of the pattern along the true focal surface.

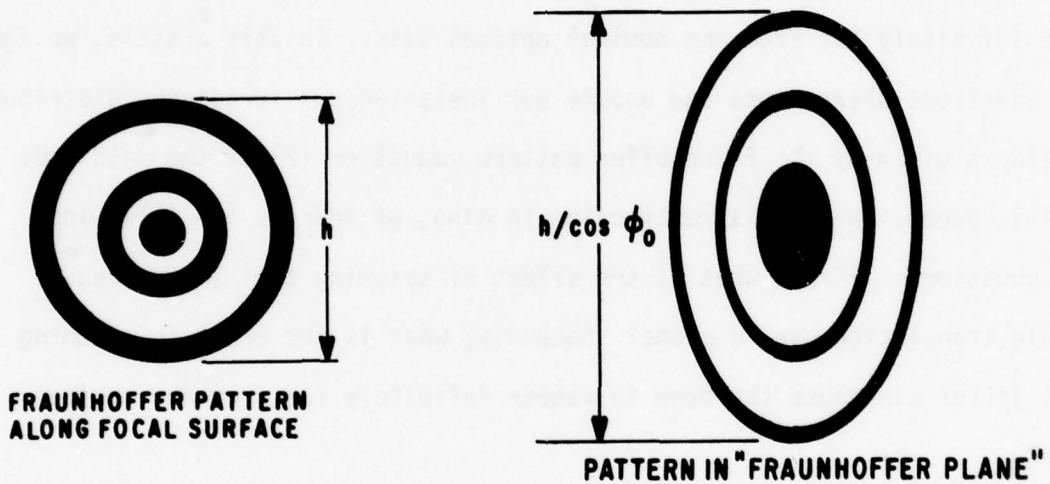


Figure A2. Elongation of Far-Field Pattern due  
Tilt of Beam by an Amount  $\phi_0$

We establish the following arbitrary criteria:

$$(a) \quad |1/\cos\phi_0| \leq 1.05$$

$$(b) \quad |\cos\phi_0| \geq \frac{f}{f + \ell}$$

where  $\ell = 1/2 \lambda(f/D)^2$  (A1)

Criterion (a) states that the elongation of the Fraunhofer pattern will not exceed five percent. Criterion (b) states the true location of peak

far-field irradiance will always be within the depth of focus of the nominal focal plane. The equation for depth of focus is derived strictly from considerations of decrease in on-axis irradiance as one moves in and out of focus. It does not describe the misfocus one can tolerate and still see higher order Fraunhoffer fringes.

Let  $\phi_{\max}$  be the upper limit of angular jitter. Relating (a) to (b)

$$\frac{f + \lambda}{f} \leq 1.05$$

and substituting equation (1) for  $\lambda$ , we obtain equation (A2).

$$\frac{\lambda f}{D^2} \leq 0.1 \quad (A2)$$

Typically, the diameter of a Fraunhoffer pattern is on the order of  $2 f \lambda / D$ . Therefore, whenever one is far enough away such that the diameter of the Fraunhoffer pattern equals or exceeds a fifth of the diameter of the focusing aperture, and for

$$\phi_{\max} \leq \cos^{-1} \left( \frac{1}{1.05} \right) \quad (A3)$$

this paper considers angular-jitter-induced focal spot wander to be confined to a plane, and jitter-induced beam elongation is considered negligible.

In this simplified model of jitter we have considered only angular jitter. We disregard displacements due to pure translational jitter on the basis that the effects of pure translational jitter are limited by the distance which the diffracting aperture can slide back and forth. Usually this distance is much smaller than the diameter of the aperture. Since, by equation (A2), we have already assumed that the size of the focused beam is

at least as large as a fifth of the aperture, pure translational jitter displacements will be small compared to the size of the focused beam whenever they are small compared to the aperture diameter.

We now seek to know the angles of jitter within which our convolution model of jitter is applicable. Let us arbitrarily stipulate that  $\phi_{\max}$  in equation (A3) is related to  $\sigma$  by

$$\phi_{\max} > 2\sigma (\lambda/D) \quad (A4)$$

where  $\sigma(\lambda/D)$  denotes the rms angular jitter. That is, we allow for the possibility of infinite jitter displacements but the time-averaged probability that the center of the jittered Fraunhofer pattern is not more than  $\phi_{\max}$  radians away from the optical axis is almost unity. That is,

$$P(\phi_{\circ} \leq \phi_{\max}) = P(\rho_{\circ} \leq 2\sigma) = 0.98$$

Therefore, effects of jitter are assumed negligible for all angles of jitter greater than  $\phi_{\max}$ . It is considered allowable to integrate (1) to the limit  $\rho_{\circ} \rightarrow \infty$ , although physically there is an upper bound on  $\rho_{\circ}$ .

In this paper the maximum value of  $\sigma$  is 3.00. Substituting this value of  $\sigma$  into inequality (A4) and relating that to inequality (A3), we infer that all calculations in this paper are a good model when

$$\lambda/D < \frac{1}{6} \cos^{-1} \left( \frac{1}{1.05} \right) \quad (A5)$$

This relation along with relation (A2)

$$\frac{\lambda f}{D^2} \geq 0.1 \quad (A2)$$

are the two basic restrictions which limit the applicability of results presented in this paper.

APPENDIX B  
CLOSED-FORM SOLUTION  
TO EQUATION (1)

Closed-form analytical solutions for  $\langle I \rangle$  exist. In the form presented here, they involve infinite sums of Laguerre polynomials. Define  $f$ ,  $g$ ,  $h$ ,  $\rho$ , and  $(\theta - \theta_0)$  as follows:

$$f(\rho) = \frac{4 f_0}{(1-\alpha^2)^2} \left\{ \frac{J_1(\pi\rho)}{\pi\rho} - \alpha^2 \frac{J_1(\pi\alpha\rho)}{\pi\alpha\rho} \right\}^2 \quad (1)$$

$$g(\rho) = \frac{e^{-\rho^2/\sigma^2}}{\pi \sigma^2} \quad (2)$$

$$h = f * g = \frac{1}{\pi\sigma^2} \iint_{-\infty}^{\infty} f(\sqrt{x_0^2+y_0^2}) \exp \left[ -\frac{(x-x_0)^2+(y-y_0)^2}{\sigma^2} \right] dx_0 dy_0 \quad (3)$$

$$= \frac{1}{\pi\sigma^2} e^{-\left(\frac{x^2+y^2}{\sigma^2}\right)} \iint_{-\infty}^{\infty} F(\sqrt{x_0^2+y_0^2}) e^{-\left(\frac{x_0^2+y_0^2}{\sigma^2}\right)} e^{-\frac{2xx_0+2yy_0}{\sigma^2}} dx_0 dy_0$$

$$\rho^2 \equiv x^2 + y^2, \quad \rho_0^2 \equiv x_0^2 + y_0^2 \quad (4)$$

$$\rho\rho_0 \cos(\theta-\theta_0) \equiv xx_0 + yy_0 \quad (5)$$

$$I_0(z) = \frac{1}{2\pi} \int_0^{2\pi} e^{z\sin(\theta-\theta_0)} d\theta = \text{Modified Bessel Function of the First Kind} \quad (6)$$

It follows that

$$\begin{aligned}
 h(\rho) &= \frac{2}{\sigma^2} e^{-\rho^2/\sigma^2} \int_0^\infty f(\rho_0) e^{-\rho_0^2/\sigma^2} I_0\left(\frac{2\rho\rho_0}{\sigma^2}\right) \rho_0 d\rho_0 \quad (7) \\
 &= \frac{2}{\pi^2 \sigma^2} \left( e^{-\rho^2/\sigma^2} \right) \frac{4 f_0}{(1-\alpha^2)^2} \left\{ \int_0^\infty \rho_0^{-1} J_1^2(\pi\rho_0) e^{-\rho_0^2/\sigma^2} I_0\left(\frac{2\rho\rho_0}{\sigma^2}\right) d\rho_0 \right. \\
 &\quad - 2\alpha \int_0^\infty \rho_0^{-1} J_1(\pi\rho_0) J_1(\pi\alpha\rho_0) e^{-\rho_0^2/\sigma^2} I_0\left(\frac{2\rho\rho_0}{\sigma^2}\right) d\rho_0 \\
 &\quad \left. + \alpha^2 \int_0^\infty \rho_0^{-1} J_1^2(\pi\alpha\rho_0) e^{-\rho_0^2/\sigma^2} I_0\left(\frac{2\rho\rho_0}{\sigma^2}\right) d\rho_0 \right\}
 \end{aligned}$$

From equation 8.447(1) in Reference 5,

$$I_0(\xi \rho_0) = \sum_{p=0}^{\infty} \frac{(\xi/2)^{2p}}{(p!)^2} (\rho_0)^{2p} \quad (1) \quad (8)$$

Letting  $\xi = 2\rho/\sigma^2$ , (9)

Eqn (7) becomes

$$\begin{aligned}
 h(\rho) &= \frac{8}{\pi^2 \sigma^2} \frac{f_0}{(1-\alpha^2)^2} e^{-\rho^2/\sigma^2} \sum_{p=0}^{\infty} \frac{(\xi/2)^{2p}}{(p!)^2} \\
 &\quad \times \left\{ \int_0^\infty \rho_0^{2p-1} J_1^2(\pi\rho_0) e^{-\rho_0^2/\sigma^2} d\rho_0 \right. \\
 &\quad - 2\alpha \int_0^\infty \rho_0^{2p-1} J_1(\pi\rho_0) J_1(\pi\alpha\rho_0) e^{-\rho_0^2/\sigma^2} d\rho_0 \\
 &\quad \left. + \alpha^2 \int_0^\infty \rho_0^{2p-1} J_1^2(\pi\alpha\rho_0) e^{-\rho_0^2/\sigma^2} d\rho_0 \right\} \quad (10)
 \end{aligned}$$

We now define  $V(a, b, p; \sigma)$ .

$$V(a, b; p; \sigma) \equiv \int_0^\infty \rho_0^{2p-1} J_1(a\rho_0) J_1(b\rho_0) e^{-\rho_0^2/\sigma^2} d\rho_0 \quad (11)$$

Using equation 6.633(1) in Reference 5, this equation becomes

$$V(a, b; p; \sigma) = \frac{a b \sigma^{2(p+1)}}{8} \sum_{m=0}^{\infty} \frac{(m+p)!}{m!} \frac{(-a^2 \sigma^2/4)^m}{(m+1)!} F(-m, -m-1; 2; \frac{b^2}{a^2}) \quad (12)$$

$$h(\rho) = \frac{8}{\pi^2 \sigma^2} \frac{f_0}{(1-\alpha^2)^2} e^{-\rho^2/\sigma^2} \sum_{p=0}^{\infty} \frac{(\xi/2)^{2p}}{p!^2} \left\{ V(\pi, \pi; p; \sigma) \right. \quad (13)$$

$$\left. - 2\alpha V(\pi, \alpha\pi; p; \sigma) + \alpha^2 V(\alpha\pi, \alpha\pi; p; \sigma) \right\}$$

$$h(\rho) = \frac{8}{\pi^2 \sigma^2} \frac{f_0 e^{-\rho^2/\sigma^2}}{(1-\alpha^2)^2} \sum_{p=0}^{\infty} \frac{(\xi/2)^{2p}}{p!^2} \frac{2p}{\sigma} \frac{2p}{\sigma^2} \frac{\pi^2}{8} \quad (14)$$

$$x \left\{ \sum_{m=0}^{\infty} \frac{(m+p)!}{m!} \frac{(-1)^m}{(m+1)!} \frac{(\pi\sigma/2)^{2m}}{F(-m, -m-1; 2; 1)} \right.$$

$$\left. - 2\alpha^2 \sum_{m=0}^{\infty} \frac{(m+p)!}{m!} \frac{(-1)^m}{(m+1)!} \frac{(\pi\sigma/2)^{2m}}{F(-m, -m-1; 2; \alpha^2)} \right\}$$

$$+ \alpha^4 \sum_{m=0}^{\infty} \frac{(m+p)!}{m!} \frac{(-1)^m}{(m+1)!} \frac{(\pi\alpha\sigma/2)^{2m}}{F(-m, -m-1; 2; 1)} \right\}$$

$$h(\rho) = \frac{f_0 e^{-\rho^2/\sigma^2}}{(1-\alpha^2)^2} \sum_{p=0}^{\infty} \frac{(\sigma\xi/2)^{2p}}{p!^2} \sum_{m=0}^{\infty} \frac{(m+p)!}{m!} \frac{(-1)^m}{(m+1)!} \frac{(\pi\sigma/2)^{2m}}{F(-m, -m-1; 2; 1)} \quad (15)$$

$$x \left\{ F(-m, -m-1; 2; 1) - 2\alpha^2 F(-m, -m-1; 2; \alpha^2) + \alpha^{4+2m} F(-m, -m-1; 2; 1) \right\}$$

Uniform convergence over  $\rho$  has been assumed.

Interchanging the order of summation,

$$h(\rho) = \frac{f_0 e^{-\rho^2/\sigma^2}}{(1-\alpha^2)^2} \sum_{m=0}^{\infty} \sum_{p=0}^{\infty} \frac{(-1)^m \left(\frac{\pi\alpha}{2}\right)^{2m}}{(m+1)!} \left[ \frac{(m+p)!}{m! p!} \frac{[\sigma^2 \xi^2/4]^p}{p!} \right] \quad (16)$$

$$h(\rho) = \frac{f_0 e^{-\rho^2/\sigma^2}}{(1-\alpha^2)^2} \sum_{m=0}^{\infty} \frac{\left(\frac{\pi\alpha}{2}\right)^{2m} (-1)^m}{(m+1)!} \times \left\{ F(-m, -m-1; 2; 1) - 2 \alpha^2 F(-m, -m-1; 2; \alpha^2) + \alpha^{4+2m} F(-m, -m-1; 2; 1) \right\} \\ \times \sum_{p=0}^{\infty} \frac{\left(\frac{\sigma^2 \xi^2}{4}\right)^p}{p!} \left[ \frac{(m+p)!}{m! p!} \right] \quad (17)$$

We pause momentarily to verify

$$\sum_{p=0}^{\infty} \frac{x^p}{p!} \left[ \frac{(m+p)!}{m! p!} \right] = e^x L_m(-x) \quad (18)$$

$$\text{Proof: } \frac{d}{dx} (x^{m+p}) = (m+p) x^{m+p-1}$$

$$\frac{d^2}{dx^2} (x^{m+p}) = (m+p) (m+p-1) x^{m+p-2}$$

$$\frac{d^m}{dx^m} (x^{m+p}) = (m+p) (m+p-1) \cdots (m+p-m+1) x^{m+p-m}$$

$$= \frac{(m+p)!}{p!} x^p$$

$$\frac{d^m}{dx^m} x^m e^x = \frac{d^m}{dx^m} \sum_{p=0}^{\infty} \frac{x^{m+p}}{p!}$$

$$= \sum_{p=0}^{\infty} \frac{x^p}{p!} \frac{(m+p)!}{p!}$$

$$\frac{1}{m!} \frac{d^m}{dx^m} \left( (-x)^m e^x \right) = \frac{(-1)^m}{m!} \frac{d^m}{dx^m} x^m e^x$$

$$= \frac{(-1)^m}{m!} \sum_{p=0}^{\infty} \frac{x^p}{p!} \left[ \frac{(m+p)!}{p!} \right]$$

$$\frac{(-1)^m}{m!} \frac{d^m}{dx^m} \left( (-x)^m e^x \right) = \sum_{p=0}^{\infty} \frac{x^p}{p!} \left[ \frac{(m+p)!}{m! p!} \right]$$

By definition, the  $m$ -th order Laguerre polynomial is

$$L_m(x) \equiv \frac{e^x}{m!} \frac{\partial^m}{\partial x^m} (e^{-x} x^m)$$

$$L_m(-x) = \frac{e^{-x}}{m!} \frac{\partial^m}{\partial(-x)^m} (e^x (-x)^m)$$

$$e^x L_m(-x) = \frac{(-1)^m}{m!} \frac{\partial^m}{\partial x^m} (-x)^m e^x$$

$$= \sum_{p=0}^{\infty} \frac{x^p}{p!} \left[ \frac{(m+p)!}{m! p!} \right]$$

From eqn (18)

$$\sum_{p=0}^{\infty} \frac{\left[ \left( \frac{\sigma \xi}{2} \right)^2 \right]^p}{p!} \left[ \frac{(m+p)!}{m! p!} \right] = e^{\left( \frac{\sigma \xi}{2} \right)^2} L_m \left( - \left( \frac{\sigma \xi}{2} \right)^2 \right) \quad (19)$$

From eqn (17)

$$h(\rho) = \frac{f_0 e^{-\rho^2/\sigma^2}}{(1-\alpha^2)^2} \sum_{m=0}^{\infty} \frac{\left( \frac{\pi \sigma}{2} \right)^{2m} (-1)^m}{(m+1)!} \times \left\{ F(-m, -m-1; 2; 1) - 2 \alpha^2 F(-m, -m-1; 2; \alpha^2) + \alpha^{4+2m} F(-m, -m-1; 2; 1) \right\} \times \left[ L_m \left( - \left( \frac{\sigma \xi}{2} \right)^2 \right) e^{\left( \frac{\sigma \xi}{2} \right)^2} \right] \quad (20)$$

Recall eqn (9)

$$\xi = \frac{2\rho}{\sigma^2}$$

Therefore, from eqns (9) and (20)

$$h(\rho) = \frac{f_0}{(1-\alpha^2)^2} \sum_{m=0}^{\infty} \frac{(\frac{\pi\sigma}{2})(-1)^m}{(m+1)} L_m \left( -(\rho/\sigma)^2 \right) \times \left\{ F(-m, -m-1; 2; 1) - 2\alpha^2 F(-m, -m-1; 2; \alpha^2) + \alpha^{4+2m} F(-m, -m-1; 2; 1) \right\} \quad (21)$$

F = Hypergeometric function

$L_m$  = Laguerre polynomial, order m

By direct comparison of eqns (1), (2) and (3) in this appendix with eqns (2), (3), and (1) in the main body of this report, one sees that

$$h(\rho) = \langle I(\rho) \rangle \quad (22)$$

when  $f_0 = I(0)$ .

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APPENDIX C  
NUMERICAL RESULTS OF JITTER CALCULATIONS

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\*\*\*\*\*  
AVANTAGE INVESTISSEMENT ET CAPITALISATION  
\*\*\*\*\*

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4. TABLE  
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1924 JAN. 1, 1924.

1.  $I(X, Y) = J^{*42}$
2.  $U = (2*J1(P1*2) - 2*ALPHA4*J1(P1*ALPHA4*P1))/P1*(1-ALPHA4*2)$
3.  $R(X, Y) = EXP(-(X**2+Y**2)/( PI*SIGHA4*2))$
4.  $D = \text{OUTER DIAMETER OF ANNULAR APERTURE}$
5.  $ALPHA4 = \text{INNER DIAMETER OF ANNULAR APERTURE}$
6.  $WL = \text{WAVELENGTH OF LASER WAVE}$
7.  $F = \text{DISTANCE FROM APERTURE TO FRAUNHOFER PLANE}$
8.  $S = \text{DISTANCE OF ANY POINT IN THE FRAUNHOFER PLANE FROM THE POINT (0,0) OF THE FRAUNHOFER PLANE}$
9.  $R04 = X**2+Y**2.$
10.  $F = 3*(F*WL)$
11.  $J1 = \text{FIRST ORDER BESSEL FUNCTION OF THE FIRST KIND}$
12.  $ALPHA = .0350$
13.  $T = 1-ALPHA4*2 = 1.0030$

THE JOURNAL OF CLIMATE

CE ANNUAL AVERAGE

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THE FEDERAL BUREAU OF INVESTIGATION, U. S. DEPARTMENT OF JUSTICE, WASHINGTON, D. C.

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\*\*\*\*\* 4V-RING IRRADIANCE OVER A CIRCULAR HOLE \*\*\*\*\*

		HOLE RADIUS	.15	.20	.25	.30	.35	.40	.45	.50
1. 1.20**	*.35									
1. 1.372	*.157*	*.1372	*.1363	*.1359	*.1352	*.1344	*.1334	*.1324	*.1312	
1. 1.360	*.156*	*.1293	*.1295	*.1291	*.1286	*.1280	*.1272	*.1264	*.1255	
1. 1.231	*.125*	*.1223	*.1227	*.1223	*.1213	*.1213	*.1206	*.1199	*.1191	
1. 1.166	*.116*	*.1167	*.1163	*.1160	*.1155	*.1151	*.1145	*.1139	*.1131	
1. 1.153	*.115*	*.1157	*.1152	*.1147	*.1142	*.1139	*.1132	*.1126	*.1120	
1. 1.154	*.115*	*.1153	*.1151	*.1148	*.1145	*.1141	*.1136	*.1130	*.1124	
1. 1.03	*.103*	*.1033	*.1032	*.1003	*.0998	*.0995	*.0991	*.0982	*.0970	
1. 0.956	*.095*	*.0956	*.0955	*.0953	*.0951	*.0948	*.0945	*.0937	*.0926	
1. 0.912	*.091*	*.0912	*.0911	*.0910	*.0905	*.0912	*.0908	*.0904	*.0895	
1. 0.871	*.087*	*.0871	*.0870	*.0863	*.0867	*.0865	*.0862	*.0858	*.0846	
1. 0.832	*.083*	*.0832	*.0831	*.0832	*.0829	*.0827	*.0824	*.0821	*.0810	
1. 0.796	*.079*	*.0796	*.0796	*.0793	*.0793	*.0791	*.0789	*.0786	*.0776	
1. 0.763	*.076*	*.0763	*.0762	*.0761	*.0760	*.0756	*.0753	*.0750	*.0747	
1. 0.731	*.073*	*.0731	*.0730	*.0733	*.0734	*.0726	*.0727	*.0723	*.0717	
1. 0.701	*.070*	*.0701	*.0701	*.0702	*.0702	*.0699	*.0697	*.0693	*.0685	
1. 0.673	*.067*	*.0673	*.0673	*.0672	*.0671	*.0670	*.0669	*.0664	*.0659	
1. 0.647	*.064*	*.0647	*.0647	*.0645	*.0645	*.0644	*.0644	*.0639	*.0633	
1. 0.622	*.062*	*.0622	*.0622	*.0621	*.0620	*.0619	*.0618	*.0614	*.0613	
1. 0.599	*.059*	*.0599	*.0593	*.0594	*.0597	*.0596	*.0595	*.0593	*.0587	
1. 0.577	*.057*	*.0577	*.0577	*.0576	*.0575	*.0574	*.0573	*.0571	*.0568	
1. 0.556	*.055*	*.0556	*.0555	*.0553	*.0554	*.0553	*.0552	*.0551	*.0547	
1. 0.536	*.053*	*.0536	*.0535	*.0533	*.0534	*.0533	*.0532	*.0531	*.0528	
1. 0.517	*.051*	*.0517	*.0517	*.0515	*.0516	*.0515	*.0514	*.0513	*.0508	
1. 0.495	*.049*	*.0495	*.0493	*.0493	*.0493	*.0492	*.0491	*.0490	*.0481	
1. 0.462	*.046*	*.0462	*.0462	*.0461	*.0461	*.0460	*.0459	*.0458	*.0457	
1. 0.406	*.040*	*.0406	*.0406	*.0406	*.0406	*.0406	*.0406	*.0405	*.0404	
1. 0.351	*.035*	*.0351	*.0351	*.0351	*.0351	*.0351	*.0351	*.0351	*.0351	
1. 0.327	*.032*	*.0327	*.0327	*.0326	*.0326	*.0325	*.0325	*.0324	*.0323	
1. 0.292	*.029*	*.0292	*.0292	*.0292	*.0292	*.0292	*.0292	*.0292	*.0292	
1. 0.259	*.025*	*.0259	*.0259	*.0259	*.0259	*.0259	*.0259	*.0259	*.0259	
1. 0.234	*.023*	*.0239	*.0239	*.0239	*.0239	*.0239	*.0239	*.0239	*.0239	

4. GIVEN

1.  $I(x, y) = I_{eff}(x, y)$  = IRRADIANCE DISTRIBUTION CENTERED AT  $(x_0, y_0)$ .  
2.  $\rho(x, y) = \rho(x_0, y_0)$  = PROBABILITY THAT THE IRRADIANCE IS  
CENTERED WITHIN  $(x_0, y_0)$  AT  $(x, y)$ .

3.  $\alpha$  = THE CONVOLUTION OF  $I_{eff}$  AND  $\rho$  OVER THE AREA OF  $A$  IS THE EFFECTIVE  
IRRADIANCE AT  $(x_0, y_0)$ . AT A POINT  $(x, y)$ , THE INTEGRAL OF  $I_{eff}(x, y)$  OVER  
ANY AREA  $A$  DIVIDED BY THE AREA OF  $A$  IS THE AVERAGE  
IRRADIANCE AT  $(x, y)$ .

4.  $I_{eff}(x, y) = I_{eff}(x_0, y_0) + \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} I(x, y) \rho(x, y) dx dy$  IS A SINGLE WHICH IS  
A SUM OF THE AREA OF  $I(x, y)$  AND THE AREA OF  $\rho(x, y)$  DUE TO THE  
POINT  $(x, y)$  OF THE  $I(x, y)$  AND THE POINT  $(x, y)$  OF THE  $\rho(x, y)$ .

5.  $I_{eff}(x, y) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} I(x, y) \rho(x, y) dx dy$  IS A SUM OF THE AREA OF  $I(x, y)$  AND THE AREA OF  $\rho(x, y)$  DUE TO THE  
POINT  $(x, y)$  OF THE  $I(x, y)$  AND THE POINT  $(x, y)$  OF THE  $\rho(x, y)$ .

6.  $I_{eff}(x, y) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} I(x, y) \rho(x, y) dx dy$  IS A SUM OF THE AREA OF  $I(x, y)$  AND THE AREA OF  $\rho(x, y)$  DUE TO THE  
POINT  $(x, y)$  OF THE  $I(x, y)$  AND THE POINT  $(x, y)$  OF THE  $\rho(x, y)$ .

7.  $I_{eff}(x, y) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} I(x, y) \rho(x, y) dx dy$  IS A SUM OF THE AREA OF  $I(x, y)$  AND THE AREA OF  $\rho(x, y)$  DUE TO THE  
POINT  $(x, y)$  OF THE  $I(x, y)$  AND THE POINT  $(x, y)$  OF THE  $\rho(x, y)$ .

8.  $I_{eff}(x, y) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} I(x, y) \rho(x, y) dx dy$  IS A SUM OF THE AREA OF  $I(x, y)$  AND THE AREA OF  $\rho(x, y)$  DUE TO THE  
POINT  $(x, y)$  OF THE  $I(x, y)$  AND THE POINT  $(x, y)$  OF THE  $\rho(x, y)$ .

9.  $I_{eff}(x, y) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} I(x, y) \rho(x, y) dx dy$  IS A SUM OF THE AREA OF  $I(x, y)$  AND THE AREA OF  $\rho(x, y)$  DUE TO THE  
POINT  $(x, y)$  OF THE  $I(x, y)$  AND THE POINT  $(x, y)$  OF THE  $\rho(x, y)$ .

1.  $I(x, y) = U^{**2}$   
2.  $U = (2 * J1(r) * J1(\rho * \alpha * r) * \rho) / (r * (1 - \alpha * \rho)^{**2})$   
3.  $J1(r) = \text{EXP}(-(X^{**2} + Y^{**2}) / (2 * r^2)) / (\text{PI} * \text{SIGMA}^{**2})$

4.  $\alpha = \text{INNER DIAMETER OF ANNULAR APERTURE}$   
5.  $r = \text{WAVELENGTH OF PLANE WAVE}$   
6.  $\rho = \text{DISTANCE FROM APERTURE TO FRAUNHOFER PLANE}$   
7.  $\text{PI} = \text{PI}$   
8.  $\text{SIGMA} = \text{SIGMA}$

9.  $\text{ALPHA} = \text{ALPHA}$

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• GIVE IT  
• TALK  
• PLAY  
• ENTERTAIN

8. THEN THE CONDUCTION OF A CIRCLE IS THE EFFECTIVE CONDUCTION L-ER, AT A PULP (X,Y). THE INTEGRAL OF  $\Sigma \Sigma F(X,Y)$  OVER ANY AREA A DIVIDED BY THE AREA OF A IS THE AVERAGE CONDUCTION OVER AREA A.

3.  $D$  = OUTER DIAMETER OF ANNULAR APERTURE  
 4.  $\alpha$  = ALPHABD IN NED DIAMETER OF ANNULAR APERTURE  
 5.  $\lambda$  = WAVELENGTH OF PLANE WAVE  
 6.  $F$  = DISTANCE FROM APERTURE TO FRAUNHOFFER PLANE  
 7.  $S$  = DISTANCE OF ANY POINT IN THE FRAUNHOFFER PLANE  
 FROM THE POINT  $(0,0)$  OF THE FRAUNHOFFER PLANE  
 8.  $\alpha_1 = \alpha_2 = \alpha^2 + y^2/2$ ,  $\alpha_1$ ,  $\alpha_2$ ,  $x$ , and  $y$  ARE DIMENSIONLESS.  
 9.  $F = \lambda y^2/4\pi^2$   
 10.  $J_1 =$  FIRST ORDER BESSEL FUNCTION OF THE FIRST KIND

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\*\*\*\*\* AVERAGE FREQUENCY OF CIRCULAR HOLE \*

HOLE RADIUS		1.00	0.95	0.90	0.85	0.80	0.75	0.70	0.65	0.60	0.55	0.50	0.45	0.40	0.35	0.30	0.25	0.20	0.15	0.10	0.05
1.50*		1.1312	1.1277	1.1253	1.1228	1.1203	1.1179	1.1255	1.1235	1.1217	1.1193	1.1169	1.1224	1.1185	1.1156	1.1126	1.1095	1.1066	1.1037	1.1008	1.0979
1.55*		1.1244	1.1233	1.1223	1.1213	1.1203	1.1193	1.1207	1.1193	1.1183	1.1179	1.1165	1.1163	1.1147	1.1130	1.1113	1.1095	1.1075	1.1053	1.1033	1.1013
1.60*		1.1161	1.1171	1.1171	1.1163	1.1163	1.1163	1.1163	1.1155	1.1155	1.1155	1.1155	1.1152	1.1147	1.1130	1.1113	1.1095	1.1075	1.1053	1.1033	1.1013
1.65*		1.1123	1.1113	1.1113	1.1103	1.1103	1.1103	1.1103	1.1095	1.1095	1.1095	1.1095	1.1092	1.1087	1.1079	1.1079	1.1066	1.1053	1.1033	1.1013	1.0993
1.70*		1.1095	1.1095	1.1095	1.1095	1.1095	1.1095	1.1095	1.1092	1.1092	1.1092	1.1092	1.1092	1.1087	1.1084	1.1084	1.1071	1.1058	1.1038	1.1017	1.0997
1.75*		1.1071	1.1071	1.1071	1.1071	1.1071	1.1071	1.1071	1.1068	1.1068	1.1068	1.1068	1.1068	1.1063	1.1061	1.1061	1.1058	1.1045	1.1032	1.1017	1.0997
1.80*		1.1054	1.1054	1.1054	1.1054	1.1054	1.1054	1.1054	1.1052	1.1052	1.1052	1.1052	1.1052	1.1047	1.1045	1.1045	1.1042	1.1039	1.1032	1.1017	1.0997
1.85*		1.1043	1.1043	1.1043	1.1043	1.1043	1.1043	1.1043	1.1042	1.1042	1.1042	1.1042	1.1042	1.1038	1.1035	1.1035	1.1032	1.1029	1.1026	1.1017	1.0997
1.90*		1.1036	1.1036	1.1036	1.1036	1.1036	1.1036	1.1036	1.1035	1.1035	1.1035	1.1035	1.1035	1.1032	1.1031	1.1031	1.1028	1.1025	1.1022	1.1017	1.0997
1.95*		1.1031	1.1031	1.1031	1.1031	1.1031	1.1031	1.1031	1.1030	1.1030	1.1030	1.1030	1.1030	1.1027	1.1026	1.1026	1.1023	1.1020	1.1017	1.1017	1.0997
2.00*		1.1026	1.1026	1.1026	1.1026	1.1026	1.1026	1.1026	1.1025	1.1025	1.1025	1.1025	1.1025	1.1022	1.1021	1.1021	1.1018	1.1015	1.1012	1.1012	1.0997
2.05*		1.1023	1.1023	1.1023	1.1023	1.1023	1.1023	1.1023	1.1022	1.1022	1.1022	1.1022	1.1022	1.1019	1.1018	1.1018	1.1015	1.1012	1.1009	1.1009	1.0997
2.10*		1.1020	1.1020	1.1020	1.1020	1.1020	1.1020	1.1020	1.1019	1.1019	1.1019	1.1019	1.1019	1.1017	1.1016	1.1016	1.1013	1.1010	1.1007	1.1007	1.0997
2.15*		1.1017	1.1017	1.1017	1.1017	1.1017	1.1017	1.1017	1.1016	1.1016	1.1016	1.1016	1.1016	1.1014	1.1013	1.1013	1.1010	1.1007	1.1004	1.1004	1.0997
2.20*		1.1014	1.1014	1.1014	1.1014	1.1014	1.1014	1.1014	1.1013	1.1013	1.1013	1.1013	1.1013	1.1011	1.1010	1.1010	1.1007	1.1004	1.1001	1.1001	1.0997
2.25*		1.1012	1.1012	1.1012	1.1012	1.1012	1.1012	1.1012	1.1011	1.1011	1.1011	1.1011	1.1011	1.1009	1.1008	1.1008	1.1005	1.1002	1.1000	1.1000	1.0997
2.30*		1.1010	1.1010	1.1010	1.1010	1.1010	1.1010	1.1010	1.1009	1.1009	1.1009	1.1009	1.1009	1.1007	1.1006	1.1006	1.1003	1.1000	1.0997	1.0997	1.0997
2.35*		1.1008	1.1008	1.1008	1.1008	1.1008	1.1008	1.1008	1.1007	1.1007	1.1007	1.1007	1.1007	1.1005	1.1004	1.1004	1.1001	1.0997	1.0997	1.0997	1.0997
2.40*		1.1007	1.1007	1.1007	1.1007	1.1007	1.1007	1.1007	1.1006	1.1006	1.1006	1.1006	1.1006	1.1004	1.1003	1.1003	1.1000	1.0997	1.0997	1.0997	1.0997
2.45*		1.1006	1.1006	1.1006	1.1006	1.1006	1.1006	1.1006	1.1005	1.1005	1.1005	1.1005	1.1005	1.1003	1.1002	1.1002	1.1000	1.0997	1.0997	1.0997	1.0997
2.50*		1.1005	1.1005	1.1005	1.1005	1.1005	1.1005	1.1005	1.1004	1.1004	1.1004	1.1004	1.1004	1.1002	1.1001	1.1001	1.0998	1.0995	1.0992	1.0992	1.0997
2.55*		1.1004	1.1004	1.1004	1.1004	1.1004	1.1004	1.1004	1.1003	1.1003	1.1003	1.1003	1.1003	1.1001	1.1000	1.1000	1.0997	1.0994	1.0991	1.0991	1.0997
2.60*		1.1003	1.1003	1.1003	1.1003	1.1003	1.1003	1.1003	1.1002	1.1002	1.1002	1.1002	1.1002	1.1000	1.0999	1.0999	1.0996	1.0993	1.0990	1.0990	1.0997
2.65*		1.1002	1.1002	1.1002	1.1002	1.1002	1.1002	1.1002	1.1001	1.1001	1.1001	1.1001	1.1001	1.0999	1.0998	1.0998	1.0995	1.0992	1.0989	1.0989	1.0997
2.70*		1.1001	1.1001	1.1001	1.1001	1.1001	1.1001	1.1001	1.1000	1.1000	1.1000	1.1000	1.1000	1.0998	1.0997	1.0997	1.0994	1.0991	1.0988	1.0988	1.0997
2.75*		1.1000	1.1000	1.1000	1.1000	1.1000	1.1000	1.1000	1.0999	1.0999	1.0999	1.0999	1.0999	1.0997	1.0996	1.0996	1.0993	1.0990	1.0987	1.0987	1.0997
2.80*		1.0999	1.0999	1.0999	1.0999	1.0999	1.0999	1.0999	1.0998	1.0998	1.0998	1.0998	1.0998	1.0996	1.0995	1.0995	1.0992	1.0989	1.0986	1.0986	1.0997
2.85*		1.0998	1.0998	1.0998	1.0998	1.0998	1.0998	1.0998	1.0997	1.0997	1.0997	1.0997	1.0997	1.0995	1.0994	1.0994	1.0991	1.0988	1.0985	1.0985	1.0997
2.90*		1.0997	1.0997	1.0997	1.0997	1.0997	1.0997	1.0997	1.0996	1.0996	1.0996	1.0996	1.0996	1.0994	1.0993	1.0993	1.0990	1.0987	1.0984	1.0984	1.0997
2.95*		1.0996	1.0996	1.0996	1.0996	1.0996	1.0996	1.0996	1.0995	1.0995	1.0995	1.0995	1.0995	1.0993	1.0992	1.0992	1.0989	1.0986	1.0983	1.0983	1.0997
3.00*		1.0995	1.0995	1.0995	1.0995	1.0995	1.0995	1.0995	1.0994	1.0994	1.0994	1.0994	1.0994	1.0992	1.0991	1.0991	1.0988	1.0985	1.0982	1.0982	1.0997

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A. Given  
TIA-AU-YU = 1420 LANGE JAGET AUTION CENTERED AT (X,Y).

1600-Y-200 = INTEGRAL DISTRIBUTION CENTERED AT (160, Y-20)

3.  $D = \text{OUTER DIAMETER OF ANNULAR APERTURE}$   
 4.  $\alpha = \text{INNER DIAMETER OF ANNULAR APERTURE}$

5. ML = WAVELENGTH OF GLENE WAVE  
6. E = DISTANCE FROM APERTURE TO THE GRAUNHOFFER DIANE FROM

17. SEPARATION OF AN ELEMENT IN THE FEINHOFF PLANE  
18. THE POINT (X<sub>0</sub>, Y<sub>0</sub>) OF THE FEINHOFF PLANE  
19. X<sub>0</sub> AND Y<sub>0</sub> DIMENSIONLESS.

$J_1 = \text{FIRST} \{ \text{PROC } \text{RECDL} \text{ FUNCTION } \text{ OF } \text{ THE } \text{ FIRST } \text{ KIN} \}$

$$\text{ALPHA} = -11111 \quad T = 1-\text{ALPHA}^{**2} = 1.0000$$

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#### AVERAGE IRRADIANCE OVER A CIRCULAR HOLE

		Radius									
		Hole					Radius				
		1.00	1.05	1.10	1.15	1.20	1.25	1.30	1.35	1.40	1.45
1.50**	1.116*	.1120	.1135	.1150	.1164	.1182	.1200	.1219	.1238	.1257	.1276
1.55**	1.149*	.1176	.1157	.1133	.1118	.1099	.1073	.1040	.1003	.951	.929
1.60**	1.156*	.1134	.1162	.1192	.1227	.1259	.1291	.1322	.1353	.1384	.1415
1.65**	1.151*	.1151	.1171	.1195	.1216	.1238	.1260	.1281	.1303	.1325	.1347
1.70**	1.155*	.1149	.1159	.1165	.1171	.1178	.1185	.1193	.1201	.1209	.1215
1.75**	1.175*	.1151	.1152	.1153	.1154	.1155	.1156	.1157	.1158	.1159	.1160
1.80**	1.177*	.1152	.1153	.1154	.1155	.1156	.1157	.1158	.1159	.1160	.1161
1.85**	1.179*	.1153	.1154	.1155	.1156	.1157	.1158	.1159	.1160	.1161	.1162
1.90**	1.181*	.1154	.1155	.1156	.1157	.1158	.1159	.1160	.1161	.1162	.1163
1.95**	1.187*	.1155	.1156	.1157	.1158	.1159	.1160	.1161	.1162	.1163	.1164
2.00**	1.191*	.1156	.1157	.1158	.1159	.1160	.1161	.1162	.1163	.1164	.1165
2.05**	1.196*	.1157	.1158	.1159	.1160	.1161	.1162	.1163	.1164	.1165	.1166
2.10**	1.201*	.1158	.1159	.1160	.1161	.1162	.1163	.1164	.1165	.1166	.1167
2.15**	1.207*	.1159	.1160	.1161	.1162	.1163	.1164	.1165	.1166	.1167	.1168
2.20**	1.212*	.1160	.1161	.1162	.1163	.1164	.1165	.1166	.1167	.1168	.1169
2.25**	1.217*	.1161	.1162	.1163	.1164	.1165	.1166	.1167	.1168	.1169	.1170
2.30**	1.221*	.1161	.1162	.1163	.1164	.1165	.1166	.1167	.1168	.1169	.1170
2.35**	1.226*	.1162	.1163	.1164	.1165	.1166	.1167	.1168	.1169	.1170	.1171
2.40**	1.231*	.1163	.1164	.1165	.1166	.1167	.1168	.1169	.1170	.1171	.1172
2.45**	1.235*	.1164	.1165	.1166	.1167	.1168	.1169	.1170	.1171	.1172	.1173
2.50**	1.239*	.1165	.1166	.1167	.1168	.1169	.1170	.1171	.1172	.1173	.1174
2.55**	1.243*	.1166	.1167	.1168	.1169	.1170	.1171	.1172	.1173	.1174	.1175
2.60**	1.247*	.1167	.1168	.1169	.1170	.1171	.1172	.1173	.1174	.1175	.1176
2.65**	1.251*	.1168	.1169	.1170	.1171	.1172	.1173	.1174	.1175	.1176	.1177
2.70**	1.255*	.1169	.1170	.1171	.1172	.1173	.1174	.1175	.1176	.1177	.1178
2.75**	1.259*	.1170	.1171	.1172	.1173	.1174	.1175	.1176	.1177	.1178	.1179
2.80**	1.263*	.1171	.1172	.1173	.1174	.1175	.1176	.1177	.1178	.1179	.1180
2.85**	1.267*	.1172	.1173	.1174	.1175	.1176	.1177	.1178	.1179	.1180	.1181
2.90**	1.271*	.1173	.1174	.1175	.1176	.1177	.1178	.1179	.1180	.1181	.1182
2.95**	1.275*	.1174	.1175	.1176	.1177	.1178	.1179	.1180	.1181	.1182	.1183
3.00**	1.279*	.1175	.1176	.1177	.1178	.1179	.1180	.1181	.1182	.1183	.1184

440 GILVÉN

•  $(X_0, Y_0)$  = IRRADIANCE DISTRIBUTION CENTERED AT  $(X_0, Y_0)$ .  
 •  $P(X_0, Y_0) * D(X_0, Y_0)$  = PROBABILITY THAT THE IRRADIANCE IS  
 CENTERED WITHIN  $(D(X_0, Y_0))$  OF  $(X_0, Y_0)$ .

Fig. 1. A 1-D convolution of  $f(x)$  with  $p$  describes the effective kernel of  $p$  at a point  $(x, y)$ . The integral of  $\text{EFF}(x, y)$  over any area  $A$  divided by the area of  $A$  is the average

FIG. 14. THE ABOVE TABLE OF DATA FOR A CIRCLE WHICH IS CENTERED AT  $(0,0)$  AND HAS A RADIUS  $R$ , AND  $R = 100$ , IS NORMALIZED FRAUNHOFER PATTERN DUE TO  $\pi$  THROUGH  $\pi/4$  IN A HORN ACETYLENE.

22.  $P(X, Y) = \text{EXP}(-(X^2+Y^2)/(\text{S}\sigma^2))$  ( PI\*SIGMA\*\*2 )  
 23. D = OUTER DIAMETER OF ANNULAR APERTURE  
 24. ALPHAD = INNER DIAMETER OF ANNULAR APERTURE  
 25. WL = WAVELENGTH OF PLANE WAVE  
 26. F = DISTANCE FROM APERTURE TO FRAUNHOFER PLANE  
 27. S = DISTANCE OF ANY POINT IN THE FRAUNHOFER PLANE FROM  
 THE POINT (0,0) OF THE FRAUNHOFER PLANE  
 28.  $2\pi/2 = X^2+Y^2/2$ ,  $X$  AND  $Y$  ARE DIMENSIONLESS.

$$\text{ALPHA} = 0.0100 \quad \tau = 1 - \text{ALPHA}^{*2} = 1.0000$$



\*\*\*\*\* AVERAGE IRRADIANCE OVER A CIRCULAR HOLE \*

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•  $z(x=0, t=0) = 1224.6$  LANGMUIR DISTRIBUTION CENTERED AT  $(x_0, y_0)$ .

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21.  $U = (2*\pi*(PI**2) - 2*\pi*LP4*J1(PI*4*ALPHA**C)) / PI / (1 - ALPHA**2)$   
 22.  $P(X, Y) = EXP(-(X**2 + Y**2) / (2 * SIGMA**2)) / (PI * SIGMA**2)$   
 23.  $D = \text{DIAMETER OF ANULAP APERTURE}$   
 24.  $ALPHA*D = \text{IN 1/4 DIAMETER OF ANULAP APERTURE}$   
 25.  $ML = \text{WAVELENGTH OF PLANE WAVE}$   
 26.  $F = \text{DISTANCE FROM 1/4 APERTURE TO FRAUNHOFER PLANE}$   
 27.  $S = \text{DISTANCE OF 1/4 Y POINT IN THE FRAUNHOFER PLANE FROM}$   
 28.  $\text{THE POINT } (0,0) \text{ OF THE FRAUNHOFER PLANE}$   
 29.  $X**2 = X**2 + Y**2$   
 30.  $C = S / (F * ML)$   
 31.  $J1 = \text{FIRST ORDER BESSEL FUNCTION OF THE FIRST KIND}$   
 32.  $ALPHA = 0.2220$   
 33.  $\gamma = 1 - ALPHA**2 = 1.0000$



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\*\*\* AVERAGE IRRADIANCE OVER A CIRCULAR HOLE \*\*\*

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- $\frac{1(X-X_0, Y-Y_0)}{2(X_0, Y_0) + D(X, Y)} =$  INTRADAY DISTRIBUTION CENTERED AT  $(X_0, Y_0)$ .
- $\frac{1(X-X_0, Y-Y_0)}{2(X_0, Y_0) + D(X, Y)} =$  POSSIBILITY THAT THE IRRADIANCE IS

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- $I(x,y) = U^{**2}$
- $U = (2*J1(PI*x) - 2*ALPHA*J1(PI*ALPHA*x))/PI*x/(1-ALPHA**2)$
- $P(x,y) = EXP(-(X**2+Y**2)/(SIGMA**2)/(PI*SIGHA**2))$
- $D = \text{OUTER DIAMETER OF ANNULAR APERTURE}$
- $D - \text{INNER DIAMETER OF ANNULAR APERTURE}$
- $ML = \text{WAVELENGTH OF PLANE WAVE}$
- $F = \text{DISTANCE FROM APERTURE TO FRAUNHOFER PLANE}$
- $S = \text{DISTANCE OF 1/M POINT IN THE FRAUNHOFER PLANE}$
- $\text{THE POINT } (0,0) \text{ OF THE FRAUNHOFER PLANE}$
- $X**2 = X**2+Y**2$
- $Y = S*D/(F*ML)$
- $J1 = \text{FIRST ORDER BESSEL FUNCTION OF THE FIRST KIND}$

$$A_{LP-12} = .12300 \quad T = 1 - A_{LP-12} = .9900$$

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\* AVERAGE IRRADIANCE OVER A CIRCULAR HOLE

4. GIVEN  
 1.  $I(x,y) = I(x,y-y_0)$  = IRRADIANCE DISTRIBUTION CENTERED AT  $(x,y_0)$ .  
 2.  $I(x,y) = I(x-y_0, y-y_0)$  = PROBABILITY THAT THE IRRADIANCE IS

$$\begin{aligned}
I(x, y) &= U^{**2} \\
U &= (2 * J1 * P1 * Q) - 2 * ALPH4 * J1 * (P1 * ALPH4 * P) * P1 / R / ((1 - ALPH4)**2) \\
P(x, y) &= EXP(-(X**2 + Y**2) / ( SIGMA**2) / ( PI * SIGMA**2)) \\
Q &= OUTER DIAMETER OF ANULUS, APERATURE \\
ALPH4 &= INNER DIAMETER OF ANULAR APERATURE \\
WL &= WAVELENGTH OF PLANE WAVE \\
F &= DISTANCE FROM APERATURE TO FRAUNHOFER PLANE \\
S &= DISTANCE OF ANY POINT IN THE FRAUNHOFER PLANE FROM \\
&\quad THE POINT (0,0) OF THE FRAUNHOFER PLANE \\
P**2 &= X**2 + Y**2. P, X, AND Y ARE DIMENSIONLESS. \\
P &= 2 * J1 / (F * WL) \\
J1 &= FIRST ORDER BESSEL FUNCTION OF THE FIRST KIND
\end{aligned}$$

$$4 = 1 - \Delta L B H D * 2 = .3500$$

\* \* \* \* \* AVERAGE IRREGULARITY OVER A CIRCULAR HOLE

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1. KAN-YO-YU = INVESTIGATING INVESTIGATION CENTERED AT KAN-YO-YU.  
2. PEKO-YO = DAD & DYO = PROBABLY THAT THE ASSAULT IS  
PERPETRATED BY HIM (KAN-YO-YU) OF KAN-YO-YU.

3. TECH.  
THE REVOLUTION OF I WITH 2 DETERMINED THE EFFECTIVE  
ROTATION OF I, E. AT A POINT (X,Y), THE ROTATIONAL  
VELOCITY OF I IS  $\omega(x,y)$ .

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Fig. 14. The "bow-tie" of the  $100\text{ cm}^{-1}$  Raman circle which is centered at  $100\text{ cm}^{-1}$  and has a radius of  $\pm 40\text{ cm}^{-1}$ .

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## \* AVERAGE IRRADIANCE OVER A CIRCULAR HOLE \*

4. 514-1  
4. 114-1  
4. 114-1  
4. 114-1

do. TAKEN  
THE CONVOLUTION OF  $\tilde{f}$  WITH A JEWEL AND THE EFFECTIVE  
INTERPOLATION,  $\tilde{I}_{\text{EFF}}$ , AT A POINT  $(x, y)$ . THE INTERVAL OF  $I_{\text{EFF}}(x, y)$

- $I(X,Y) = J^{**2}$
- $U = (2**J1(PI*X)) - 2*ALPHA*(J1(PI*ALPHA*X)) / PI/R / (1-ALPHA**2)$
- $P(X,Y) = EXP(-(X**2+Y**2)/(SIGMA**2)) / PI*SIGMA**2$
- $J = \text{OUTER DIAMETER OF ANNULAR APERTURE}$
- $ALPHA = \text{INNER DIAMETER OF ANNULAR APERTURE}$
- $HL = \text{WAVELENGTH OF PLANE WAVE}$
- $Z = \text{DISTANCE FROM APERTURE TO FRAUNHOFER PLANE}$
- $S = \text{DISTANCE OF ANY POINT IN THE FRAUNHOFER PLANE FROM THE POINT } (C,0) \text{ OF THE FRAUNHOFER PLANE}$
- $X = X**2+Y**2$
- $R = C**2 / (F*WL)$
- $J1 = \text{FIRST ORDER BESSEL FUNCTION OF THE FIRST KTN}$

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\* AVERAGE IRRADIANCE OVER A CIRCULAR HOLE \*

六書

1.  $I(X-X_0, Y-Y_0) = I$  IRRADIANCE DENSITY DISTRIBUTED C-NERED AT  $(X_0, Y_0)$ .  
 2.  $I(X_0, Y_0) = P$ OWER DENSITY THAT THE IRRADIANCE IS  
 DIFFUSED WITHIN  $(X_0 \pm \Delta X, Y_0 \pm \Delta Y)$ .

$$1. \quad I(X) = \frac{1}{2} \sum_{i=1}^2 I(X_i) = \frac{1}{2} \sum_{i=1}^2 (H(X_i) - H(X_i|Y)) = \frac{1}{2} \sum_{i=1}^2 (H(X_i) - H(X_i|X_{-i})) = \frac{1}{2} \sum_{i=1}^2 I(X_i|X_{-i})$$

TABLE 4. THE CONVERGENCE OF THE INTEGRAL OF THE EFFECTIVE  
KERR-ALBRECHT LIEFF AT A POINT (X, Y). THE INTEGRAL OF LIEFF (X, Y)  
IS COMPUTED BY THE INTEGRAL OF THE AVERAGE  
KERR-ALBRECHT LIEFF.

4.4.  $\alpha = \frac{2\pi}{\lambda} \frac{d}{s}$  INNER DIAMETER OF ANULAR APERTURE  
 $\lambda$  = WAVELENGTH OF PLANE WAVE  
 $d$  = DISTANCE FROM APERTURE TO FRAUNHOFER PLANE  
 $s$  = DISTANCE OF ANY POINT IN THE FRAUNHOFER PLANE FROM  
 THE POINT ( $x_0$ ) OF THE FRAUNHOFER PLANE

6. IN THE ABOVE TABLE OF FIG. 10, ADD A CIRCLE WHICH IS  
CENTRED AT (0, 0) AND PASSING THROUGH THE POINTS  
(X, 0) = (40.8425, 0) AND (0, 24.0425).

$R = S^4(F^*ML)$   
 $\therefore J_1 = \text{FIRST ORDER ACCESSEL FUNCTION OF THE FSTST KIN}$

AVERAGE IRRADIANCE OVER A CIRCULAR HOLE

M. D. REED  
CA. 1940

3. FIGURE 34-36 ARE THE CONSTRUCTION OF A MITER JOINT WITH THE EFFECTIVE DIVISIONS OF THE MITER JOINT. THE LATENT OF MITER JOINT IS THE AREA OF A MITER JOINT WHICH IS THE AREA OF A MITER JOINT.

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1.  I(XY) = U**2
2.  U = (2*X1 - 2*ALPHA**2 + J1*(1+ALPHA**2)) / (1-ALPHA**2)
3.  D(XY) = EXP(-(X**2+Y**2)/2) - SIGMA**2 / (1+SIGMA**2)

```

33.  $D = 0.443 \cdot \text{DIAMETER} + 0.2$  OF ANNULUS APERTURE  
34.  $\text{ALPHA} \cdot D = \text{INNER DIAMETER OF ANNULUS APERTURE}$

ALL WAVELENGTHS OF LIGHT HAVE THE SAME SPEED, THEREFORE, THE DISTANCE FROM A POINT IN THE EMISSION-REFLECTION PLANE TO THE POINT (C) OF THE EMISSION-REFLECTION PLANE

$x^2 + y^2 = 2$  AND  $x$  AND  $y$  ARE DIMENSIONLESS  
 $F = \mathbb{C}^2 / \langle F \rangle_{\text{UL}}$

### 3. *J<sub>1</sub>* = FIRST ORDER RECESS FUNCTION OF THE FIRST KING

$$12.8448 = 0.7472$$

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\*\*\*\*\* AVERAGE IRRADIANCE OVER A CIRCULAR HOLE \*\*\*\*\*

		HOLE RADIUS											
		.60					.65						
		.40		.50		.55		.60		.65			
1.000	1.000	.3635	+.7626	+.5667	+.1933	+.2586	+.1837	+.1515	+.1380	+.1253	+.1098	1.000	1.000
.997	.997	.3632	+.7673	+.5373	+.3768	+.2573	+.1876	+.1537	+.1372	+.1234	+.1082	.997	.997
.994	.994	.3626	+.7631	+.4382	+.3334	+.2495	+.1923	+.1571	+.1357	+.1196	+.1047	.994	.994
.991	.991	.3621	+.7611	+.4361	+.3374	+.2600	+.2284	+.1863	+.1333	+.1157	+.1010	.991	.991
.988	.988	.3617	+.7595	+.4251	+.3293	+.2296	+.1995	+.1713	+.1281	+.1115	+.0975	.988	.988
.985	.985	.3612	+.7581	+.4231	+.3283	+.2243	+.1884	+.1726	+.1525	+.1354	+.1056	.985	.985
.982	.982	.3607	+.7567	+.4151	+.3153	+.2151	+.1661	+.1451	+.1331	+.1210	+.0981	.982	.982
.979	.979	.3602	+.7552	+.4137	+.3137	+.2135	+.1475	+.1239	+.1156	+.1023	+.0897	.979	.979
.976	.976	.3604	+.7543	+.4123	+.3123	+.2123	+.1423	+.1233	+.1153	+.1005	+.0878	.976	.976
.973	.973	.3605	+.7535	+.4129	+.3113	+.2115	+.1415	+.1223	+.1123	+.1005	+.0861	.973	.973
.970	.970	.3606	+.7526	+.4124	+.3103	+.2103	+.1405	+.1215	+.1115	+.1005	+.0852	.970	.970
.967	.967	.3607	+.7517	+.4077	+.3095	+.2095	+.1395	+.1191	+.1087	+.0973	+.0832	.967	.967
.964	.964	.3608	+.7507	+.4043	+.3086	+.2080	+.1386	+.1182	+.1074	+.0969	+.0820	.964	.964
.961	.961	.3609	+.7494	+.4026	+.3076	+.2070	+.1376	+.1172	+.1067	+.0959	+.0810	.961	.961
.958	.958	.3610	+.7484	+.4016	+.3065	+.2060	+.1365	+.1162	+.1055	+.0950	+.0800	.958	.958
.955	.955	.3611	+.7475	+.4004	+.3053	+.2050	+.1353	+.1151	+.1046	+.0943	+.0795	.955	.955
.952	.952	.3612	+.7466	+.3994	+.3042	+.2040	+.1342	+.1140	+.1036	+.0936	+.0786	.952	.952
.949	.949	.3613	+.7457	+.3984	+.3031	+.2030	+.1331	+.1130	+.1025	+.0926	+.0778	.949	.949
.946	.946	.3614	+.7448	+.3974	+.3020	+.2020	+.1320	+.1120	+.1015	+.0915	+.0769	.946	.946
.943	.943	.3615	+.7439	+.3964	+.3009	+.2010	+.1310	+.1110	+.1005	+.0905	+.0760	.943	.943
.940	.940	.3616	+.7430	+.3954	+.3008	+.2000	+.1300	+.1100	+.1000	+.0900	+.0750	.940	.940
.937	.937	.3617	+.7421	+.3944	+.3007	+.1990	+.1290	+.1090	+.0990	+.0749	+.0749	.937	.937
.934	.934	.3618	+.7412	+.3934	+.3006	+.1980	+.1280	+.1080	+.0980	+.0748	+.0748	.934	.934
.931	.931	.3619	+.7403	+.3924	+.3005	+.1970	+.1270	+.1070	+.0970	+.0747	+.0747	.931	.931
.928	.928	.3620	+.7394	+.3914	+.3004	+.1960	+.1260	+.1060	+.0960	+.0746	+.0746	.928	.928
.925	.925	.3621	+.7385	+.3904	+.3003	+.1950	+.1250	+.1050	+.0950	+.0745	+.0745	.925	.925
.922	.922	.3622	+.7376	+.3894	+.3002	+.1940	+.1240	+.1040	+.0940	+.0744	+.0744	.922	.922
.919	.919	.3623	+.7367	+.3884	+.3001	+.1930	+.1230	+.1030	+.0930	+.0743	+.0743	.919	.919
.916	.916	.3624	+.7358	+.3874	+.3000	+.1920	+.1220	+.1020	+.0920	+.0742	+.0742	.916	.916
.913	.913	.3625	+.7349	+.3864	+.3000	+.1910	+.1210	+.1010	+.0910	+.0741	+.0741	.913	.913
.910	.910	.3626	+.7340	+.3854	+.3000	+.1900	+.1200	+.1000	+.0900	+.0740	+.0740	.910	.910
.907	.907	.3627	+.7331	+.3844	+.3000	+.1890	+.1190	+.0990	+.0739	+.0739	+.0739	.907	.907
.904	.904	.3628	+.7322	+.3834	+.3000	+.1880	+.1180	+.0980	+.0738	+.0738	+.0738	.904	.904
.901	.901	.3629	+.7313	+.3824	+.3000	+.1870	+.1170	+.0970	+.0737	+.0737	+.0737	.901	.901
.898	.898	.3630	+.7304	+.3814	+.3000	+.1860	+.1160	+.0960	+.0736	+.0736	+.0736	.898	.898
.895	.895	.3631	+.7295	+.3804	+.3000	+.1850	+.1150	+.0950	+.0735	+.0735	+.0735	.895	.895
.892	.892	.3632	+.7286	+.3794	+.3000	+.1840	+.1140	+.0940	+.0734	+.0734	+.0734	.892	.892
.889	.889	.3633	+.7277	+.3784	+.3000	+.1830	+.1130	+.0930	+.0733	+.0733	+.0733	.889	.889
.886	.886	.3634	+.7268	+.3774	+.3000	+.1820	+.1120	+.0920	+.0732	+.0732	+.0732	.886	.886
.883	.883	.3635	+.7259	+.3764	+.3000	+.1810	+.1110	+.0910	+.0731	+.0731	+.0731	.883	.883
.880	.880	.3636	+.7250	+.3754	+.3000	+.1800	+.1100	+.0900	+.0730	+.0730	+.0730	.880	.880
.877	.877	.3637	+.7241	+.3744	+.3000	+.1790	+.1090	+.0890	+.0729	+.0729	+.0729	.877	.877
.874	.874	.3638	+.7232	+.3734	+.3000	+.1780	+.1080	+.0880	+.0728	+.0728	+.0728	.874	.874
.871	.871	.3639	+.7223	+.3724	+.3000	+.1770	+.1070	+.0870	+.0727	+.0727	+.0727	.871	.871
.868	.868	.3640	+.7214	+.3714	+.3000	+.1760	+.1060	+.0860	+.0726	+.0726	+.0726	.868	.868
.865	.865	.3641	+.7205	+.3704	+.3000	+.1750	+.1050	+.0850	+.0725	+.0725	+.0725	.865	.865
.862	.862	.3642	+.7196	+.3694	+.3000	+.1740	+.1040	+.0840	+.0724	+.0724	+.0724	.862	.862
.859	.859	.3643	+.7187	+.3684	+.3000	+.1730	+.1030	+.0830	+.0723	+.0723	+.0723	.859	.859
.856	.856	.3644	+.7178	+.3674	+.3000	+.1720	+.1020	+.0820	+.0722	+.0722	+.0722	.856	.856
.853	.853	.3645	+.7169	+.3664	+.3000	+.1710	+.1010	+.0810	+.0721	+.0721	+.0721	.853	.853
.850	.850	.3646	+.7160	+.3654	+.3000	+.1700	+.1000	+.0800	+.0720	+.0720	+.0720	.850	.850
.847	.847	.3647	+.7151	+.3644	+.3000	+.1690	+.990	+.0790	+.0719	+.0719	+.0719	.847	.847
.844	.844	.3648	+.7142	+.3634	+.3000	+.1680	+.980	+.0780	+.0718	+.0718	+.0718	.844	.844
.841	.841	.3649	+.7133	+.3624	+.3000	+.1670	+.970	+.0770	+.0717	+.0717	+.0717	.841	.841
.838	.838	.3650	+.7124	+.3614	+.3000	+.1660	+.960	+.0760	+.0716	+.0716	+.0716	.838	.838
.835	.835	.3651	+.7115	+.3604	+.3000	+.1650	+.950	+.0750	+.0715	+.0715	+.0715	.835	.835
.832	.832	.3652	+.7106	+.3594	+.3000	+.1640	+.940	+.0740	+.0714	+.0714	+.0714	.832	.832
.829	.829	.3653	+.7097	+.3584	+.3000	+.1630	+.930	+.0730	+.0713	+.0713	+.0713	.829	.829
.826	.826	.3654	+.7088	+.3574	+.3000	+.1620	+.920	+.0720	+.0712	+.0712	+.0712	.826	.826
.823	.823	.3655	+.7079	+.3564	+.3000	+.1610	+.910	+.0710	+.0711	+.0711	+.0711	.823	.823
.820	.820	.3656	+.7070	+.3554	+.3000	+.1600	+.900	+.0700	+.0710	+.0710	+.0710	.820	.820
.817	.817	.3657	+.7061	+.3544	+.3000	+.1590	+.890	+.0690	+.0709	+.0709	+.0709	.817	.817
.814	.814	.3658	+.7052	+.3534	+.3000	+.1580	+.880	+.0680	+.0708	+.0708	+.0708	.814	.814
.811	.811	.3659	+.7043	+.3524	+.3000	+.1570	+.870	+.0670	+.0707	+.0707	+.0707	.811	.811
.808	.808	.3660	+.7034	+.3514	+.3000	+.1560	+.860	+.0660	+.0706	+.0706	+.0706	.808	.808
.805	.805	.3661	+.7025	+.3504	+.3000	+.1550	+.850	+.0650	+.0705	+.0705	+.0705	.805	.805
.802	.802	.3662	+.7016	+.3494	+.3000	+.1540	+.840	+.0640	+.0704	+.0704	+.0704	.802	.802
.799	.799	.3663	+.7007	+.3484	+.3000	+.1530	+.830	+.0630	+.0703	+.0703	+.0703	.799	.799
.796	.796	.3664	+.6998	+.3474	+.3000	+.1520	+.820	+.0620	+.0702	+.0702	+.0702	.796	.796
.793	.793	.3665	+.6989	+.3464	+.3000	+.1510	+.810	+.0610	+.0701	+.0701	+.0701	.793	.793
.790	.790	.3666	+.6980	+.3454	+.3000	+.1500	+.800	+.0600	+.0700	+.0700	+.0700	.790	.790
.787	.787	.3667	+.6971	+.3444	+.3000	+.1490	+.790	+.0590	+.0699	+.0699	+.0699	.787	.787
.784	.784	.3668	+.6962	+.3434	+.3000	+.1480	+.780	+.0580	+.0698	+.0698	+.0698	.784	.784
.781	.781	.3669	+.6953	+.3424	+.3000	+.1470	+.770	+.0570	+.0697	+.0697	+.0697	.781	.781
.778	.778	.3670	+.6944	+.3414	+.3000	+.1460	+.760	+.0560	+.0696	+.0696	+.0696	.778	.778
.775	.775	.3671	+.6935	+.3404	+.3000	+.1450	+.750	+.0550	+.0695	+.0695	+.0695	.775	.775
.772	.772	.3672	+.6926	+.3394	+.3000	+.1440	+.740	+.0540	+.0694	+.0694	+.0694	.772	.772
.769	.769	.3673	+.6917	+.3384	+.3000	+.1430	+.730	+.0530	+.0693	+.0693	+.0693	.769	.769
.766	.766	.3674	+.6908	+.3374	+.3000	+.1420	+.720	+.0520	+.0692	+.0692	+.0692	.766	.766
.763	.763	.3675	+.6899	+.3364	+.3000	+.1410	+.710	+.0510	+.0691	+.0691	+.0691	.763	.763
.760	.760	.3676	+.6890	+.3354	+.3000	+.1400	+.700	+.0500	+.0690	+.0690	+.0690	.760	.760
.757	.757	.3677	+.6881	+.3344	+.3000	+.1390	+.690	+.0490	+.0689	+.0689	+.0689	.757	.757
.754	.754	.3678	+.6872	+.3334	+.3000	+.1380	+.680	+.0480	+.0688	+.0688	+.0688	.754	.754
.751	.751	.3679	+.6863	+.3324	+.3000	+.1370	+.670	+.0470	+.0687	+.0687	+.0687	.751	.751
.748	.748	.3680	+.6854	+.3314	+.3000	+.1360	+.660	+.0460	+.0686	+.0686	+.0686	.748	.748
.745	.745	.3681	+.6845	+.3304	+.3000	+.1350	+.650	+.0450	+.0685</				

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TABLE I  
CALCULATION OF THE EFFECTIVE  
RADIANCE AT A POINT  $(x, y)$ . THE INTERVAL OF  $\text{LEFF}(x, y)$   
IS DIVIDED BY  $n$  AND THE AVERAGE  
RADIANCE IS  $\bar{A}_{\text{eff}}$ .

1.  $U = (2 * J1(PI * S) - 2 * ALPHA * J1(PI * ALPHA * P)) / PI * P / (1 - ALPHA * ALPHA * 2)$

2.  $P(X, Y) = EXP(-(X * 2 * Y * * 2) / (SIGMA * * 2)) / (PI * SIGMA * * 2)$

3.  $D = \text{OUTER DIAMETER OF ANNULAR APERTURE}$

4.  $ALPHA = \text{INNER DIAMETER OF ANNULAR APERTURE}$

5.  $ML = \text{WAVELLENGTH OF PLANE WAVE}$

6.  $F = \text{DISTANCE FROM APERTURE TO FRAUNHOFER PLANE}$

7.  $S = \text{DISTANCE OF ANY POINT IN THE FRAUNHOFER PLANE FROM THE POINT } (C, 0) \text{ OF THE FRAUNHOFER PLANE}$

8.  $P * * 2 = X * 2 * Y * * 2$ ,  $X, Y$  ARE DIMENSIONLESS.

9.  $F = C * D / (F * ML)$

10.  $J1 = \text{FIRST ORDER BESSEL FUNCTION OF THE FIRST KIND}$

11.  $T = 1 - ALPHA * ALPHA * 2 = 1 - 75/36 = -5/12$

\* AVERAGE TRANSLIANCE OVER A CIRCULAR HOLE \*

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THE IRRADIANCE IS DETERMINED BY THE PRODUCT OF THE IRRADIANCE CENTERED AT  $\lambda_{0,1}$  AND THE PRODUCT OF THE IRRADIANCE CENTERED AT  $\lambda_{0,2}$ .

30 T. E. H.

THE CONSTRUCTION OF A MITRA DISTRICTS THE EFFECTIVE  
TERRITORIES OF THE DISTRICTS OF THE INDIAN STATE OF JHARKHAND.

1.  $I(X, Y) = U^{**2}$   
 $U = (2 * J1(D1 * R) - 2 * ALPH1 * J1(D1 * ALPH1 * R)) / PI / P((1 - ALPH1)**2)$   
 $P(X, Y) = EXP(-(X**2 + Y**2)) / ( PI * S1 * CHA**2)$   
 $2. I(X, Y) = U^{**2}$   
 $U = (2 * J1(D1 * R) - 2 * ALPH1 * J1(D1 * ALPH1 * R)) / PI / P((1 - ALPH1)**2)$   
 $P(X, Y) = EXP(-(X**2 + Y**2)) / ( PI * S1 * CHA**2)$

4.  $\alpha$  = INNER DIAMETER OF ANULUS APERTURE  
 5.  $\lambda$  = WAVELENGTH OF PLANE HAVE  
 6.  $F$  = DISTANCE FOCAL APERTURE TO FRAUNHOFER PLANE  
 7.  $S$  = DISTANCE OF ANY POINT IN THE FRAUNHOFER PLANE  
 8.  $\theta$  = POINT (0,0) OF THE FRAUNHOFER PLANE  
 9.  $x$  =  $\lambda F \cos \theta$   
 10.  $y$  =  $\lambda F \sin \theta$   
 11.  $\alpha = 2\pi x / \lambda F$   
 12.  $\alpha = 2\pi y / \lambda F$

$$F = \frac{1}{2} \partial_x^2 (F^4 \partial_x^4) \quad J_1 = \text{FIRST ORDER BESSEL FUNCTION OF THE FIRST KIND}$$

AFWL-TR-76-66

APPENDIX D  
FORTRAN PROGRAM JITTRAN

Appendix D IS BEST QUALITY PRACTICABLE  
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per teletex w/ AFWL

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      SUBROUTINE EXTERNAL (TITLE, G, N, W)
      DIMENSION G(4097),TITLE(31),V(11),N(11),W(11)
      COMMON SIGMA,ALPHA
      DATA TITLE/31*1H /
      TITLE( 3)=1HR
      TITLE( 6)=1HM
      TITLE( 9)=1HS
      TITLE(14)=1HJ
      TITLE(17)=1HI
      TITLE(20)=1HT
      TITLE(23)=1HT
      TITLE(26)=1HE
      TITLE(29)=1HR
      130 FOR T=1 TO 29 BY 1 READ (TITLE,T) V(T)
      140 READ (TITLE,31) N(1),N(2),N(3),N(4),N(5),N(6),N(7),N(8),N(9),N(10)
      150 READ (TITLE,32) W(1),W(2),W(3),W(4),W(5),W(6),W(7),W(8),W(9),W(10)
      160 READ (TITLE,33) SIGMA,ALPHA
      170 READ (TITLE,34) T1,T2
      180 READ (TITLE,35) T3,T4
      190 READ (TITLE,36) T5,T6
      200 READ (TITLE,37) T7,T8
      210 READ (TITLE,38) T9,T10
      220 READ (TITLE,39) T11,T12
      230 READ (TITLE,40) T13,T14
      240 READ (TITLE,41) T15,T16
      250 READ (TITLE,42) T17,T18
      260 READ (TITLE,43) T19,T20
      270 READ (TITLE,44) T21,T22
      280 READ (TITLE,45) T23,T24
      290 READ (TITLE,46) T25,T26
      300 READ (TITLE,47) T27,T28
      310 READ (TITLE,48) T29,T30
      320 READ (TITLE,49) T31,T32
      330 READ (TITLE,50) T33,T34
      340 READ (TITLE,51) T35,T36
      350 READ (TITLE,52) T37,T38
      360 READ (TITLE,53) T39,T40
      370 READ (TITLE,54) T41,T42
      380 READ (TITLE,55) T43,T44
      390 READ (TITLE,56) T45,T46
      400 READ (TITLE,57) T47,T48
      410 READ (TITLE,58) T49,T50
      420 READ (TITLE,59) T51,T52
      430 READ (TITLE,60) T53,T54
      440 READ (TITLE,61) T55,T56
      450 READ (TITLE,62) T57,T58
      460 READ (TITLE,63) T59,T60
      470 READ (TITLE,64) T61,T62
      480 READ (TITLE,65) T63,T64
      490 READ (TITLE,66) T65,T66
      500 READ (TITLE,67) T67,T68
      510 READ (TITLE,68) T69,T70
      520 READ (TITLE,69) T71,T72
      530 READ (TITLE,70) T73,T74
      540 READ (TITLE,71) T75,T76
      550 READ (TITLE,72) T77,T78
      560 READ (TITLE,73) T79,T80
      570 READ (TITLE,74) T81,T82
      580 READ (TITLE,75) T83,T84
      590 READ (TITLE,76) T85,T86
      600 READ (TITLE,77) T87,T88
      610 READ (TITLE,78) T89,T90
      620 READ (TITLE,79) T91,T92
      630 READ (TITLE,80) T93,T94
      640 READ (TITLE,81) T95,T96
      650 READ (TITLE,82) T97,T98
      660 READ (TITLE,83) T99,T100
      670 READ (TITLE,84) T101,T102
      680 READ (TITLE,85) T103,T104
      690 READ (TITLE,86) T105,T106
      700 READ (TITLE,87) T107,T108
      710 READ (TITLE,88) T109,T110
      720 READ (TITLE,89) T111,T112
      730 READ (TITLE,90) T113,T114
      740 READ (TITLE,91) T115,T116
      750 READ (TITLE,92) T117,T118
      760 READ (TITLE,93) T119,T120
      770 READ (TITLE,94) T121,T122
      780 READ (TITLE,95) T123,T124
      790 READ (TITLE,96) T125,T126
      800 READ (TITLE,97) T127,T128
      810 READ (TITLE,98) T129,T130
      820 READ (TITLE,99) T131,T132
      830 READ (TITLE,100) T133,T134
      840 READ (TITLE,101) T135,T136
      850 READ (TITLE,102) T137,T138
      860 READ (TITLE,103) T139,T140
      870 READ (TITLE,104) T141,T142
      880 READ (TITLE,105) T143,T144
      890 READ (TITLE,106) T145,T146
      900 READ (TITLE,107) T147,T148
      910 READ (TITLE,108) T149,T150
      920 READ (TITLE,109) T151,T152
      930 READ (TITLE,110) T153,T154
      940 READ (TITLE,111) T155,T156
      950 READ (TITLE,112) T157,T158
      960 READ (TITLE,113) T159,T160
      970 READ (TITLE,114) T161,T162
      980 READ (TITLE,115) T163,T164
      990 READ (TITLE,116) T165,T166
      1000 READ (TITLE,117) T167,T168
      1010 READ (TITLE,118) T169,T170
      1020 READ (TITLE,119) T171,T172
      1030 READ (TITLE,120) T173,T174
      1040 READ (TITLE,121) T175,T176
      1050 READ (TITLE,122) T177,T178
      1060 READ (TITLE,123) T179,T180
      1070 READ (TITLE,124) T181,T182
      1080 READ (TITLE,125) T183,T184
      1090 READ (TITLE,126) T185,T186
      1100 READ (TITLE,127) T187,T188
      1110 READ (TITLE,128) T189,T190
      1120 READ (TITLE,129) T191,T192
      1130 READ (TITLE,130) T193,T194
      1140 READ (TITLE,131) T195,T196
      1150 READ (TITLE,132) T197,T198
      1160 READ (TITLE,133) T199,T200
      1170 READ (TITLE,134) T201,T202
      1180 READ (TITLE,135) T203,T204
      1190 READ (TITLE,136) T205,T206
      1200 READ (TITLE,137) T207,T208
      1210 READ (TITLE,138) T209,T210
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      1230 READ (TITLE,140) T213,T214
      1240 READ (TITLE,141) T215,T216
      1250 READ (TITLE,142) T217,T218
      1260 READ (TITLE,143) T219,T220
      1270 READ (TITLE,144) T221,T222
      1280 READ (TITLE,145) T223,T224
      1290 READ (TITLE,146) T225,T226
      1300 READ (TITLE,147) T227,T228
      1310 READ (TITLE,148) T229,T230
      1320 READ (TITLE,149) T231,T232
      1330 READ (TITLE,150) T233,T234
      1340 READ (TITLE,151) T235,T236
      1350 READ (TITLE,152) T237,T238
      1360 READ (TITLE,153) T239,T240
      1370 READ (TITLE,154) T241,T242
      1380 READ (TITLE,155) T243,T244
      1390 READ (TITLE,156) T245,T246
      1400 READ (TITLE,157) T247,T248
      1410 READ (TITLE,158) T249,T250
      1420 READ (TITLE,159) T251,T252
      1430 READ (TITLE,160) T253,T254
      1440 READ (TITLE,161) T255,T256
      1450 READ (TITLE,162) T257,T258
      1460 READ (TITLE,163) T259,T260
      1470 READ (TITLE,164) T261,T262
      1480 READ (TITLE,165) T263,T264
      1490 READ (TITLE,166) T265,T266
      1500 READ (TITLE,167) T267,T268
      1510 READ (TITLE,168) T269,T270
      1520 READ (TITLE,169) T271,T272
      1530 READ (TITLE,170) T273,T274
      1540 READ (TITLE,171) T275,T276
      1550 READ (TITLE,172) T277,T278
      1560 READ (TITLE,173) T279,T280
      1570 READ (TITLE,174) T281,T282
      1580 READ (TITLE,175) T283,T284
      1590 READ (TITLE,176) T285,T286
      1600 READ (TITLE,177) T287,T288
      1610 READ (TITLE,178) T289,T290
      1620 READ (TITLE,179) T291,T292
      1630 READ (TITLE,180) T293,T294
      1640 READ (TITLE,181) T295,T296
      1650 READ (TITLE,182) T297,T298
      1660 READ (TITLE,183) T299,T300
      1670 READ (TITLE,184) T301,T302
      1680 READ (TITLE,185) T303,T304
      1690 READ (TITLE,186) T305,T306
      1700 READ (TITLE,187) T307,T308
      1710 READ (TITLE,188) T309,T310
      1720 READ (TITLE,189) T311,T312
      1730 READ (TITLE,190) T313,T314
      1740 READ (TITLE,191) T315,T316
      1750 READ (TITLE,192) T317,T318
      1760 READ (TITLE,193) T319,T320
      1770 READ (TITLE,194) T321,T322
      1780 READ (TITLE,195) T323,T324
      1790 READ (TITLE,196) T325,T326
      1800 READ (TITLE,197) T327,T328
      1810 READ (TITLE,198) T329,T330
      1820 READ (TITLE,199) T331,T332
      1830 READ (TITLE,200) T333,T334
      1840 READ (TITLE,201) T335,T336
      1850 READ (TITLE,202) T337,T338
      1860 READ (TITLE,203) T339,T340
      1870 READ (TITLE,204) T341,T342
      1880 READ (TITLE,205) T343,T344
      1890 READ (TITLE,206) T345,T346
      1900 READ (TITLE,207) T347,T348
      1910 READ (TITLE,208) T349,T350
      1920 READ (TITLE,209) T351,T352
      1930 READ (TITLE,210) T353,T354
      1940 READ (TITLE,211) T355,T356
      1950 READ (TITLE,212) T357,T358
      1960 READ (TITLE,213) T359,T360
      1970 READ (TITLE,214) T361,T362
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      1990 READ (TITLE,216) T365,T366
      2000 READ (TITLE,217) T367,T368
      2010 READ (TITLE,218) T369,T370
      2020 READ (TITLE,219) T371,T372
      2030 READ (TITLE,220) T373,T374
      2040 READ (TITLE,221) T375,T376
      2050 READ (TITLE,222) T377,T378
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      2080 READ (TITLE,225) T383,T384
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      2100 READ (TITLE,227) T387,T388
      2110 READ (TITLE,228) T389,T390
      2120 READ (TITLE,229) T391,T392
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      2140 READ (TITLE,231) T395,T396
      2150 READ (TITLE,232) T397,T398
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      2170 READ (TITLE,234) T401,T402
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      2190 READ (TITLE,236) T405,T406
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      5460 READ (TITLE,563) T1059,T1059
      5470 READ (TITLE,564) T1061,T1062
      5480 READ (TITLE,565) T1063,T1064
      5490 READ (TITLE,566) T1065,T
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C  USER MUST INITIALIZE THE FOLLOWING DATA.  ALPHA, SIGMA, SIGMA1, SIGMA2, and T.
C  DROMAX, and M.
C
C      ALPHA = 1.0
C      T = 1.0+ALPHA*ALPHA
C      SIGMA = 1.5
C      SIGMA1 = 1.0
C      SIGMA2 = 0.0
C      DROMAX = 0.05
C      M = 1
C
C  END OF INITIALIZATION
C  TDEL = NO. OF PTS. ON EACH DROMAX INTERVAL FOR WHICH INTEGRAND, G(TD), IS CALCULATED.
C  TDEL = 2
C  11 COLUMNS OF DATA
C  00 1  I=1*11
C  1  V(I)=DRMAX1+(I-1)*DROMAX
C  2  PRINT 100,ALPHA,T
C  3  PRINT 110
C  4  PRINT 110,(V(I)),(I=1,11)
C  5  PRINT 110
C  6  PRINT 140,ALPHA
C  7  I=11
C  M=M+1
C
C  11 COLUMNS OF DATA
C  SIGMAX = DRMAX1+DRMAX2
C  SIGDEL=SIGMAX/2.0
C  JTOT=INT((1.0+SIGMAX)/DROMAX)
C  JTOT=2
C
C  FILE TO STORE OTHER VALUES OF G(I)
C  00 2  J=JL+JTOT+2
C  3  SIG(J)=0.0
C  4  SIG1=SIG2=0.
C
C  COMMENT 1  INITIALIZE J, SUM1, SUM2
C  J=2
C  SUM1=0
C  SUM2=0
C
C  COMMENT 2  INITIALIZE J, SUM1, SUM2, OF OTHER VALUES OF G(I)
C  J=2
C  SUM1=0
C  SUM2=0
C
C  5  I=I+1
C  6  DR= 1.0/T
C  7  JTOT=INT(DR)
C  8  SUM1=SUM1+SIG1
C  9  SUM2=SUM2+SIG2
C
C  COMMENT 3  CALCULATE SIG1
C  10  SIG1=SIG1+JTOT*JTOT
C  11  PRINT 100
C  12  G(J+1)= INTEGRAL(G(I),DR)
C  13  SIG2=SIG2+G(J+1)
C
C  COMMENT 4  PRINT SUM1, SUM2
C
C  14  V(I)=2.0*SIG1+1.0*SIG2+0.5*(SIG1+SIG2)
C  15  V(I)=V(I)*V(I)+SIG2
C
C  COMMENT 5  IF JLT=1, THEN PRINT SUM1, SUM2, AND G(J+1) FOR A STAT.
C  COMMENT 6  PRINT SUM1, SUM2, AND G(J+1) FOR A STAT.
C  COMMENT 7  IF JLT=1, THEN PRINT SUM1, SUM2, AND G(J+1) FOR A STAT.
C  COMMENT 8  PRINT SUM1, SUM2, AND G(J+1) FOR A STAT.

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AD-A065 978

AIR FORCE WEAPONS LAB KIRTLAND AFB N MEX  
GAUSSIAN JITTER OF A FOCUSED BEAM OF LIGHT. (U)  
APR 76 W T WHITE, J P BAUMGARDNER, D A HOLMES

UNCLASSIFIED

AFWL-TR-76-66

SBIE-AD-E200 250

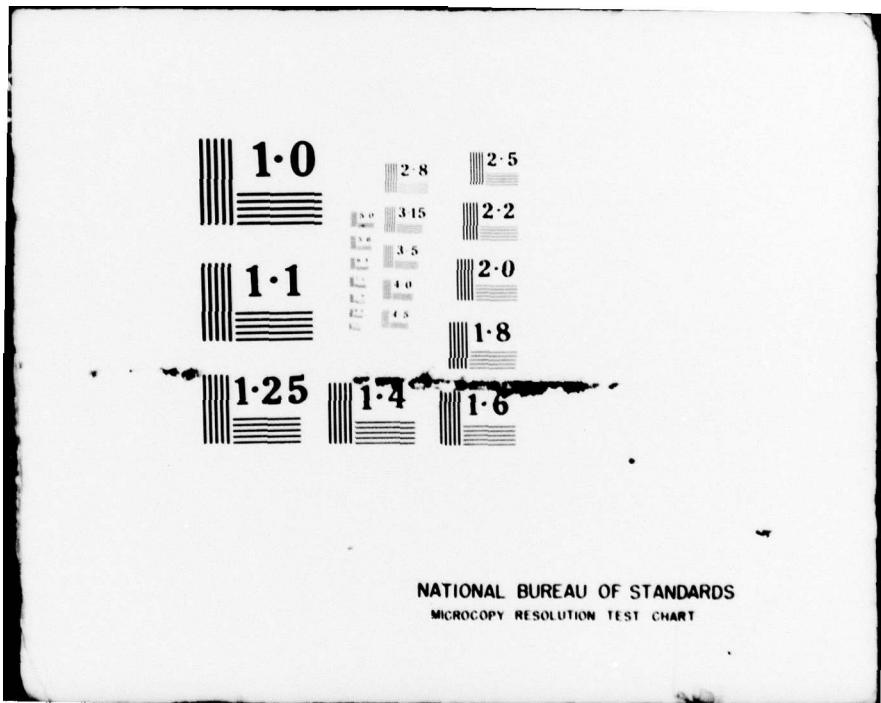
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2 OF 2  
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REAL FUNCTION INTEGRAL(X)

C THIS ALGORITHM ADAPTED FROM PATHOLOGY, A RESTED GAUSSIAN QUADRATURE

C MULTIPLE INTEGRAL SCHEME, AIR FORCE RESEARCH LABORATORY, KIRTLAND AFB, NM

C

COMMENT INTEGRAL(X) INTEGRATES FROM R1 TO THE FAIR TO 4\*8SIGMA AND

C THETA=PI/4. THE INTEGRAL IS A PRODUCT OF (R1-THETA)\*T TIMES A

C PROBABILITY DISTRIBUTION P(X). SINCE THE INTEGRAL IS PERFORMED

C IN SPATIAL COORDINATES, THERE IS AN ADDITIONAL FACTOR OF R1/4

COMMENT THE INTEGRAL.

COMMON SIGMA

DIMENSION A(100),ANSP(1),I(100),R(100),SIG(100)

DIMENSION AT(100),ANSE(100),T(100),RT(100),UT(100)

DATA R1/-,950269356477362,-,74556477413527,-,625538409916329,

1-,183434542493655,10343468249502,-,625538409916329,

2,746666477413527,3602683783562,-,3602683783562,-,313735545877367,

2,222381034453374,-,10126836290376,

DATA RT/-,90269825477362,-,74556477413527,-,625538409916329,

1-,183434542493655,10343468249502,-,625538409916329,

2,746666477413527,3602683783562,-,3602683783562,-,313735545877367,

2,222381034453374,-,10126836290376,

INTEGRAL=0.

IF(SIGMA-.01)10,11,20

C

C FOR SIGMA NEAR ZERO, THE PROBABILITY BECOMES APPROXIMATELY A DELTA

C DELTA FUNCTION. THE INTEGRAL OF F(X)\*P(X) = F(0).

C

10 INTEGRAL=F(0+.01\*X)

RETURN

20 R1=-SIGMA

C

C BREAK INTEGRAL INTO 4 INTEGRALS FROM 0 TO 4\*8SIGMA.

C

30 29 T=1.4

COMMENT

COMMENT BEGIN INTEGRAL=

COMMENT

R1=0.1+8\*1.4

R2=R1+8\*1.4

CHKD=0.

NR=1

AR(1)=0

31 RNR=R

SURP=.

HR=(R2-R1)/NR

DO 29 I=1,NR

R1=1.

32 AR(I+1)=R1+R\*HR

DO 29 I=1,NR

ANSP(I+1)=0.

DO 29 I=1,NR

R(I)=R(I+1)-AR(I+1)\*SURP(I\*HR)+AR(I+1)+AR(I+1)

33 R(I\*HR)=R(I\*HR)/2.

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AFWL-TR-76-66

SINCE THE INTEGRAL IS OF THE NATURE OF THE INTEGRAL FROM ZERO TO PI, AND SINCE THE INTEGRAL IS SYMMETRIC, AND SINCE  $\pi = \pi$ , AND SINCE WE WANT TO INTEGRATE FROM THE  $\pi = 2\pi$  TO THE  $\pi = \pi$ , WE INTEGRATE IN FACTOR OF TWO INTO THE INTEGRAL.

6

**NTEGRAL=2. NTEGRAL  
RETURN  
END**

```

FUNCTION P(x)
C COMPUTE S T-D GAUSSIAN PROBABILITY DENSITY
C COMMON SIGMA
SIGMA2=SIGMA*SIGMA
P=EXP(-X*X/SIGMA2)/3.14159265358979312389
IF P<0.0 P=0.0
END

```

AFWL-TR-76-66

APPENDIX E

$\langle i \rangle$  AS A FUNCTION OF  $\rho$  AND  $\sigma$  FOR THE CASE  $\alpha = 0$

## RADIAL PROFILE OF IEFF(RHO) AT EVENLY-SPACED VALUES OF RHO

IEFF (RHO + .05H)										
RHO	M=0	M=1	M=2	M=3	M=4	M=5	M=6	M=7	M=8	M=9
0.0 + .00000	.99385	.97558	.94575	.90527	.85535	.79745	.73324	.66451	.59711	
.5 + .52086	.44350	.38064	.31568	.25577	.20181	.15440	.11387	.08027	.05119	
1.0 + .03283	.01799	.00814	.00249	.00018	.00037	.00225	.00518	.00846	.01162	
1.5 + .01429	.01622	.01729	.01746	.01630	.01544	.01355	.01133	.00897	.00667	
2.0 + .00457	.00281	.00145	.00055	.00008	.00002	.00029	.00081	.00148	.00220	
2.5 + .00239	.00348	.00390	.00412	.00414	.00396	.00361	.00313	.00257	.00198	
3.0 + .00161	.00090	.00049	.00020	.00004	.00000	.00008	.00026	.00049	.00076	
3.5 + .00103	.00126	.00145	.00156	.00160	.00156	.00145	.00128	.00107	.00074	
4.0 + .00061										
SIGMA=0.0 T=1.00										
IEFF (RHO + .05H)										
RHO	M=0	M=1	M=2	M=3	M=4	M=5	M=6	M=7	M=8	M=9
0.0 + .90889	.90365	.86809	.86266	.82811	.78542	.73579	.69059	.62129	.55043	
.5 + .49652	.43402	.37328	.31549	.26167	.21260	.16886	.13077	.09846	.07183	
1.0 + .05061	.03438	.02258	.01460	.00979	.00745	.00694	.00767	.00910	.01079	
1.5 + .01239	.01366	.01443	.01464	.01427	.01339	.01211	.01054	.00892	.00798	
2.0 + .00545	.00402	.00285	.00199	.00145	.00120	.00121	.00143	.00178	.00220	
2.5 + .00263	.00300	.00329	.00345	.00347	.00336	.00313	.00291	.00241	.00190	
3.0 + .00157	.00119	.00086	.00062	.00046	.00039	.00041	.00049	.00067	.00079	
3.5 + .00095	.00111	.00123	.00131	.00134	.00132	.00124	.00113	.00099	.00082	
4.0 + .00065										
SIGMA= .2 T=1.00										
IEFF (RHO + .05H)										
RHO	M=0	M=1	M=2	M=3	M=4	M=5	M=6	M=7	M=8	M=9
0.0 + .70622	.70286	.69295	.67645	.65408	.62630	.59380	.55736	.51783	.47617	
.5 + .43312	.38372	.34676	.30500	.26511	.22767	.19310	.16174	.13379	.10970	
1.0 + .06626	.07354	.05594	.04416	.03491	.02784	.02259	.01982	.01620	.01447	
1.5 + .01325	.01245	.01186	.01134	.01081	.01022	.00955	.00880	.00799	.00714	
2.0 + .00631	.00551	.00479	.00416	.00364	.00324	.00295	.00275	.00264	.00250	
2.5 + .00258	.00259	.00259	.00258	.00255	.00247	.00237	.00223	.00207	.00189	
3.0 + .00170	.00152	.00136	.00121	.00109	.00100	.00095	.00091	.00091	.00071	
3.5 + .00093	.00095	.00097	.00098	.00098	.00098	.00093	.00098	.00082	.00076	
4.0 + .00069										
SIGMA= .4 T=1.00										
IEFF (RHO + .05H)										
RHO	M=0	M=1	M=2	M=3	M=4	M=5	M=6	M=7	M=8	M=9
0.0 + .60134	.59883	.59137	.57913	.56239	.54153	.51702	.48939	.45924	.42717	
.5 + .33383	.35384	.32581	.29230	.25981	.22878	.19959	.17250	.14774	.12541	
1.0 + .10557	.08319	.07320	.06045	.04978	.04098	.03383	.02811	.02359	.02005	
1.5 + .01731	.01518	.01351	.01218	.01109	.01015	.00931	.00854	.00782	.00712	
2.0 + .00647	.00585	.00528	.00477	.00431	.00392	.00359	.00331	.00309	.00292	
2.5 + .00277	.00266	.00256	.00247	.00238	.00228	.00218	.00207	.00195	.00183	
3.0 + .00170	.00158	.00147	.00136	.00127	.00118	.00112	.00106	.00102	.00100	
3.5 + .00097	.00095	.00093	.00092	.00090	.00087	.00084	.00081	.00077	.00073	
4.0 + .00069										
SIGMA= .5 T=1.00										

## RADIAL PROFILE OF IEFF(RHO) AT EVENLY-SPACED VALUES OF RHO

IEFF(RHO + .05M)										
RHO	M=0	M=1	M=2	M=3	M=4	M=5	M=6	M=7	M=8	M=9
0.0 +	.50722	.50540	.49999	.49110	.47891	.46367	.44566	.42526	.40284	.37882
.5 +	.35361	.32766	.30136	.27512	.24931	.22426	.20024	.17751	.15623	.13656
1.0 +	.11858	.10232	.08780	.07495	.06372	.05400	.04567	.03861	.03268	.02773
1.5 +	.02364	.02127	.01750	.01523	.01336	.01180	.01051	.00941	.00847	.00765
2.0 +	.00694	.00631	.00574	.00524	.00479	.00439	.00404	.00373	.00345	.00322
2.5 +	.00302	.00284	.00267	.00253	.00239	.00227	.00215	.00204	.00193	.00182
3.0 +	.00172	.00162	.00153	.00144	.00136	.00128	.00122	.00116	.00110	.00106
3.5 +	.00101	.00098	.00094	.00091	.00086	.00085	.00081	.00078	.00075	.00072
4.0 +	.00069									

SIGMA= .6 T=1.00

IEFF(RHO + .05M)										
RHO	M=0	M=1	M=2	M=3	M=4	M=5	M=6	M=7	M=8	M=9
0.0 +	.36109	.36016	.35739	.35281	.34650	.33855	.32908	.31824	.30619	.29306
.5 +	.27907	.26439	.24921	.23371	.21807	.20247	.18704	.17195	.15730	.14322
1.0 +	.12980	.11709	.10517	.09406	.08379	.07435	.06574	.05794	.05092	.04464
1.5 +	.03905	.03411	.02976	.02595	.02262	.01974	.01724	.01509	.01324	.01165
2.0 +	.01028	.00911	.00810	.00723	.00649	.00584	.00529	.00481	.00439	.00402
2.5 +	.00369	.00341	.00315	.00293	.00273	.00255	.00238	.00223	.00210	.00197
3.0 +	.00156	.00175	.00166	.00157	.00149	.00141	.00134	.00127	.00121	.00115
3.5 +	.00010	.000105	.000100	.000096	.000091	.000088	.000084	.000080	.000077	.000074
4.0 +	.00071									

SIGMA= .6 T=1.00

IEFF(RHO + .05M)										
RHO	M=0	M=1	M=2	M=3	M=4	M=5	M=6	M=7	M=8	M=9
0.0 +	.26357	.26308	.26161	.25918	.25582	.25156	.24645	.24056	.23393	.22665
.5 +	.21878	.21041	.20162	.19249	.18311	.17356	.16392	.15426	.14465	.13510
1.0 +	.12588	.11682	.10806	.09962	.09154	.08385	.07657	.06971	.06329	.05729
1.5 +	.05172	.04658	.04136	.03753	.03358	.02999	.02675	.02382	.02120	.01885
2.0 +	.01675	.01488	.01323	.01176	.01047	.00933	.00833	.00744	.00667	.00599
2.5 +	.00540	.00487	.00442	.00401	.00366	.00335	.00307	.00293	.00261	.00242
3.0 +	.00225	.00209	.00195	.00183	.00172	.00161	.00152	.00143	.00136	.00128
3.5 +	.00122	.00115	.00110	.00104	.00100	.00095	.00091	.00087	.00083	.00079
4.0 +	.00076									

SIGMA=1.0 T=1.00

IEFF(RHO + .05M)										
RHO	M=0	M=1	M=2	M=3	M=4	M=5	M=6	M=7	M=8	M=9
0.0 +	.19873	.19846	.19763	.19626	.19437	.19196	.18905	.18567	.18185	.17752
.5 +	.17301	.16806	.16280	.15728	.15153	.14560	.13952	.13334	.12709	.12081
1.0 +	.11455	.10932	.10216	.09611	.09019	.08442	.07882	.07341	.06821	.06323
1.5 +	.05848	.05396	.04968	.04564	.04185	.03829	.03497	.03187	.02900	.02535
2.0 +	.02390	.02165	.01959	.01771	.01599	.01442	.01300	.01172	.01056	.00951
2.5 +	.00857	.00773	.00697	.00630	.00569	.00515	.00467	.00424	.00385	.00351
3.0 +	.00321	.00293	.00269	.00248	.00228	.00211	.00196	.00182	.00169	.00159
3.5 +	.00148	.00139	.00130	.00123	.00116	.00109	.00104	.00098	.00093	.00089
4.0 +	.00084									

SIGMA=1.2 T=1.00

## RADIAL PROFILE OF IEFF(RHO) AT EVENLY-SPACED VALUES OF RHO

IEFF(RHO + .05M)										
RHO	M=0	M=1	M=2	M=3	M=4	M=5	M=6	M=7	M=8	M=9
0.0 +	.15443	.15426	.15377	.15296	.15182	.15037	.14862	.14658	.14425	.14168
.5 +	.13854	.13378	.13231	.12904	.12540	.12161	.11769	.11365	.10953	.10534
1.0 +	.10110	.09684	.09256	.08830	.08406	.07986	.07573	.07167	.06769	.06381
1.5 +	.06046	.05639	.05296	.04946	.04619	.04306	.04008	.03724	.03454	.03198
2.0 +	.02957	.02730	.02517	.02317	.02130	.01955	.01793	.01642	.01503	.01374
2.5 +	.01254	.01145	.01044	.00952	.00867	.00790	.00720	.00655	.00597	.00544
3.0 +	.00496	.00452	.00413	.00377	.00345	.00316	.00289	.00266	.00244	.00225
3.5 +	.00207	.00192	.00177	.00165	.00153	.00143	.00133	.00125	.00117	.00110
4.0 +	.00113									

SIGMA=1.4 T=1.00

IEFF(RHO + .05M)										
RHO	M=0	M=1	M=2	M=3	M=4	M=5	M=6	M=7	M=8	M=9
0.0 +	.12313	.12303	.12272	.12221	.12149	.12058	.11947	.11817	.11659	.11504
.5 +	.11322	.11125	.10913	.10686	.10448	.10197	.09937	.09666	.09388	.09103
1.0 +	.08812	.08517	.08218	.07916	.07613	.07311	.07009	.06708	.06411	.06117
1.5 +	.05827	.05543	.05264	.04992	.04726	.04468	.04218	.03976	.03743	.03518
2.0 +	.03302	.03095	.02897	.02708	.02528	.02357	.02194	.02041	.01896	.01759
2.5 +	.01631	.01510	.01397	.01291	.01192	.01099	.01013	.00933	.00859	.00791
3.0 +	.00727	.00668	.00614	.00564	.00518	.00476	.00437	.00402	.00370	.00340
3.5 +	.00313	.00288	.00265	.00244	.00225	.00208	.00193	.00178	.00165	.00153
4.0 +	.00143									

SIGMA=1.6 T=1.00

IEFF(RHO + .05M)										
RHO	M=0	M=1	M=2	M=3	M=4	M=5	M=6	M=7	M=8	M=9
1.0 +	.10033	.10026	.10005	.09972	.09924	.09864	.09791	.09705	.09607	.09497
.5 +	.09375	.09243	.09100	.08948	.08786	.08616	.08437	.08251	.08059	.07860
1.5 +	.07656	.07448	.07235	.07020	.06802	.06582	.06360	.06139	.05917	.05696
2.0 +	.05476	.05257	.05041	.04828	.04618	.04412	.04209	.04011	.03817	.03629
2.5 +	.03445	.03267	.03095	.02928	.02766	.02611	.02462	.02318	.02181	.02050
3.0 +	.01924	.01804	.01690	.01582	.01479	.01381	.01289	.01202	.01120	.01043
3.5 +	.00970	.00901	.00837	.00777	.00721	.00668	.00620	.00574	.00531	.00492
4.0 +	.00455	.00421	.00390	.00360	.00334	.00309	.00285	.00264	.00245	.00227
4.0 +	.00210									

SIGMA=1.8 T=1.00

IEFF(RHO + .05M)										
RHO	M=0	M=1	M=2	M=3	M=4	M=5	M=6	M=7	M=8	M=9
0.0 +	.08324	.08319	.08306	.08283	.08250	.08209	.08159	.08100	.08032	.07957
.5 +	.07873	.07781	.07682	.07576	.07463	.07343	.07218	.07086	.06950	.06888
1.0 +	.06663	.06513	.06359	.06202	.06043	.05881	.05717	.05552	.05386	.05219
1.5 +	.05252	.04885	.04718	.04553	.04388	.04225	.04064	.03904	.03747	.03593
2.0 +	.03442	.03293	.03148	.03006	.02867	.02732	.02501	.02474	.02351	.02232
2.5 +	.02116	.02005	.01898	.01795	.01696	.01601	.01510	.01423	.01340	.01261
3.0 +	.01145	.01113	.01045	.00980	.00918	.00860	.00805	.00753	.00704	.00658
3.5 +	.00614	.00573	.00535	.00498	.00465	.00433	.00403	.00375	.00349	.00325
4.0 +	.00303									

SIGMA=2.0 T=1.00

## RADIAL PROFILE OF IEFF(RHO) AT EVENLY-SPACED VALUES OF RHO

IEFF(RHO + .05M)										
RHO	M=0	M=1	M=2	M=3	M=4	M=5	M=6	M=7	M=8	M=9
0.0 +	.07014	.07010	.07001	.06984	.06962	.06932	.06897	.06855	.06807	.06757
.5 +	.06634	.06529	.06558	.06482	.06401	.06315	.06224	.06129	.06030	.05927
1.0 +	.05821	.05711	.05598	.05482	.05364	.05243	.05120	.04996	.04870	.04743
1.5 +	.04615	.04487	.04358	.04229	.04100	.03972	.03844	.03717	.03591	.03465
2.0 +	.03343	.03221	.03100	.02982	.02866	.02752	.02640	.02510	.02424	.02310
2.5 +	.02217	.02118	.02022	.01928	.01838	.01750	.01665	.01583	.01504	.01427
3.0 +	.01354	.01283	.01215	.01150	.01088	.01028	.00971	.00916	.00854	.00815
3.5 +	.00767	.00722	.00680	.00639	.00601	.00564	.00530	.00497	.00466	.00437
4.0 +	.00410									

SIGMA=2.2 T=1.00

IEFF(RHO + .05M)										
RHO	M=0	M=1	M=2	M=3	M=4	M=5	M=6	M=7	M=8	M=9
0.0 +	.05988	.05985	.05978	.05967	.05950	.05929	.05903	.05873	.05838	.05799
.5 +	.05755	.05708	.05656	.05600	.05541	.05477	.05411	.05340	.05267	.05191
1.0 +	.05111	.05029	.04944	.04857	.04768	.04677	.04583	.04489	.04392	.04295
1.5 +	.04196	.04097	.03996	.03896	.03794	.03693	.03591	.03490	.03389	.03288
2.0 +	.03188	.03089	.02990	.02893	.02796	.02701	.02607	.02514	.02423	.02334
2.5 +	.02246	.02160	.02075	.01993	.01912	.01834	.01757	.01682	.01610	.01539
3.0 +	.01471	.01405	.01340	.01278	.01218	.01160	.01104	.01050	.00998	.00949
3.5 +	.00900	.00854	.00810	.00767	.00727	.00688	.00651	.00615	.00582	.00549
4.0 +	.00519									

SIGMA=2.4 T=1.00

IEFF(RHO + .05M)										
RHO	M=0	M=1	M=2	M=3	M=4	M=5	M=6	M=7	M=8	M=9
0.0 +	.05170	.05168	.05163	.05154	.05142	.05126	.05107	.05085	.05059	.05030
.5 +	.04997	.04962	.04923	.04881	.04837	.04789	.04739	.04686	.04631	.04573
1.0 +	.04513	.04450	.04386	.04319	.04251	.04181	.04109	.04036	.03961	.03885
1.5 +	.03809	.03731	.03652	.03573	.03493	.03412	.03331	.03250	.03169	.03087
2.0 +	.03006	.02925	.02845	.02765	.02685	.02606	.02527	.02450	.02373	.02297
2.5 +	.02222	.02148	.02076	.02004	.01934	.01865	.01797	.01731	.01666	.01603
3.0 +	.01541	.01480	.01421	.01364	.01308	.01254	.01201	.01150	.01100	.01052
3.5 +	.01005	.00960	.00916	.00874	.00834	.00795	.00757	.00720	.00686	.00652
4.0 +	.00620									

SIGMA=2.6 T=1.00

IEFF(RHO + .05M)										
RHO	M=0	M=1	M=2	M=3	M=4	M=5	M=6	M=7	M=8	M=9
0.0 +	.04509	.04507	.04503	.04496	.04487	.04475	.04461	.04444	.04424	.04402
.5 +	.04377	.04350	.04321	.04289	.04255	.04218	.04180	.04140	.04097	.04052
1.0 +	.04006	.03958	.03908	.03856	.03803	.03749	.03693	.03636	.03577	.03519
1.5 +	.03457	.03396	.03333	.03270	.03206	.03142	.03077	.03012	.02946	.02880
2.0 +	.02814	.02748	.02682	.02616	.02551	.02485	.02420	.02355	.02290	.02227
2.5 +	.02163	.02101	.02038	.01977	.01916	.01857	.01798	.01740	.01683	.01627
3.0 +	.01572	.01518	.01465	.01413	.01362	.01312	.01264	.01216	.01170	.01125
3.5 +	.01081	.01038	.00997	.00957	.00917	.00879	.00843	.00807	.00772	.00730
4.0 +	.00706									

SIGMA=2.8 T=1.00

## RADIAL PROFILE OF IEFF(RHO) AT EVENLY-SPACED VALUES OF RHO

RHO	IEFF(RHO + .05M)									
	M=0	M=1	M=2	M=3	M=4	M=5	M=6	M=7	M=8	M=9
1.0 + .03965	.03964	.03961	.03956	.03949	.03940	.03929	.03916	.03903	.03883	
.5 + .03854	.03843	.03820	.03796	.03769	.03741	.03711	.03680	.03646	.03612	
1.0 + .03575	.03538	.03499	.03458	.03416	.03373	.03329	.03284	.03239	.03190	
1.5 + .03142	.03093	.03043	.02993	.02941	.02890	.02837	.02784	.02731	.02677	
2.0 + .02623	.02569	.02515	.02461	.02406	.02352	.02298	.02243	.02189	.02176	
2.5 + .02082	.02029	.01976	.01924	.01872	.01820	.01770	.01719	.01670	.01621	
3.0 + .01572	.01524	.01477	.01431	.01386	.01341	.01297	.01254	.01212	.01171	
3.5 + .01131	.01091	.01053	.01015	.00978	.00942	.00907	.00873	.00840	.00807	
4.0 + .00776										

SIGMA=3.0 T=1.00